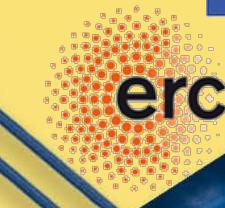


Quantenmechanik mit Schaltkreisen: Photonen und Qubits auf einem supraleitenden Mikrochip

Andreas Wallraff (*ETH Zurich*)

www.qudev.ethz.ch

Team: R. Buijs, M. Collodo, S. Gasparinetti, J. Heinsoo, P. Kurpiers,
M. Mondal, M. Oppliger, M. Pechal, A. Potocnik, Y. Salathe,
M. Stammeier, A. Stockklauser, T. Thiele, T. Walter
(*ETH Zurich*)



Eidgenössische Technische Hochschule Zürich Swiss NATIONAL SCIENCE FOUNDATION
Swiss Federal Institute of Technology Zurich

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Former group members now

Faculty/PostDoc/PhD/Industry

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M. Allan (Leiden)
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J. Basset (U. Paris Sud)
S. Berger (AWK Group)
R. Bianchetti (ABB)
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C. Eichler (Princeton)
A. Fedorov (UQ Brisbane)
A. Fragner (Yale)
S. Filipp (IBM)
J. Fink (Caltech, IST Austria)
T. Frey (Bosch)
M. Goppl (Sensirion)
J. Govenius (Aalto)
L. Huthmacher (Cambridge)

D.-D. Jarausch ([Cambridge](#))

K. Juliusson ([CEA Saclay](#))

C. Lang ([Radionor](#))

P. Leek ([Oxford](#))

P. Maurer ([Stanford](#))

J. Mlynek ([Siemens](#))

G. Puebla ([IBM](#))

A. Safavi-Naeini ([Stanford](#))

L. Steffen ([AWK Group](#))

A. van Loo ([Oxford](#))

S. Zeytinoğlu ([ETH Zurich](#))

J. Faist ([ETH Zurich](#))

J. Gambetta ([IBM](#))

T. Ihn ([ETH Zurich](#))

F. Merkt ([ETH Zurich](#))

L. Novotny ([ETH Zurich](#))

B. Sanders ([Calgary](#))

S. Schmidt ([ETH Zurich](#))

R. Schoelkopf (Yale)

C. Schoenenberger ([Basel](#))

E. Solano ([UPV/EHU](#))

W. Wegscheider ([ETH Zurich](#))

Collaborations with (groups of):

A. Blais ([Sherbrooke](#))

C. Bruder ([Basel](#))

M. da Silva ([Raytheon](#))

L. DiCarlo ([TU Delft](#))

K. Ensslin ([ETH Zurich](#))



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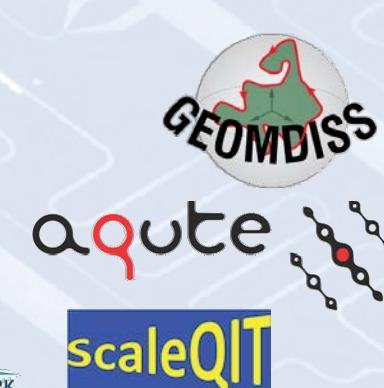
National Centre of Competence in Research



CIRCUIT AND CAVITY
QUANTUM ELECTRODYNAMICS



SEVENTH FRAMEWORK
PROGRAMME



Conventional Electronic Circuits

basic circuit elements:

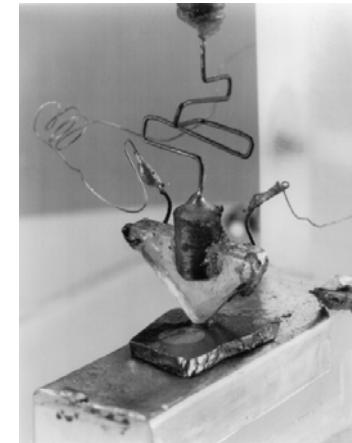


basis of modern
information and
communication
technology

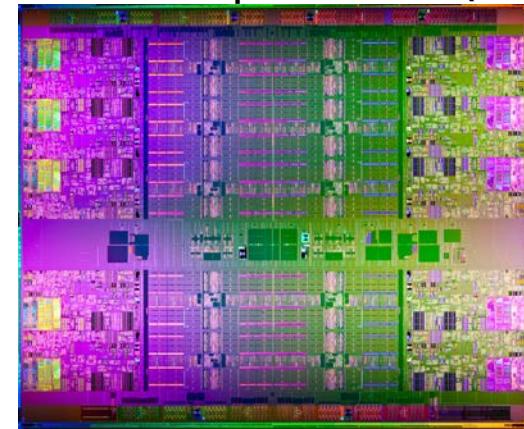
properties :

- classical physics
- no quantum mechanics
- no superposition principle
- no quantization of fields

first transistor at Bell Labs (1947)



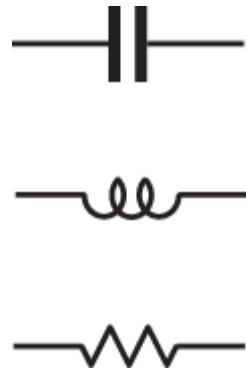
intel xeon processors (2011)



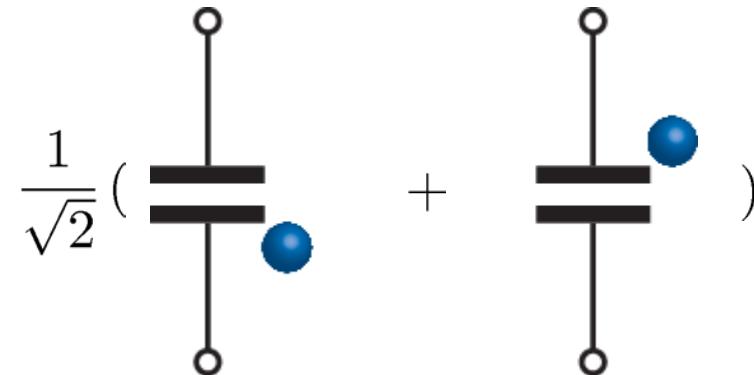
3.000.000.000 transistors
smallest feature size 32 nm
clock speed ~ 3 GHz
power consumption ~ 10 W

Classical and Quantum Electronic Circuit Elements

basic circuit elements:



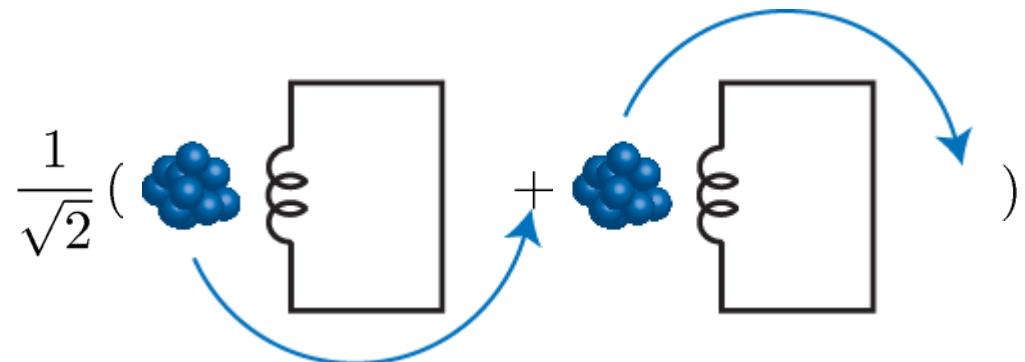
charge on a capacitor:



quantum superposition states:

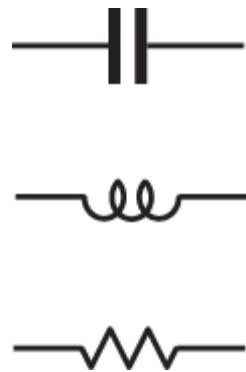
- charge q
- flux ϕ

current or magnetic flux in an inductor:

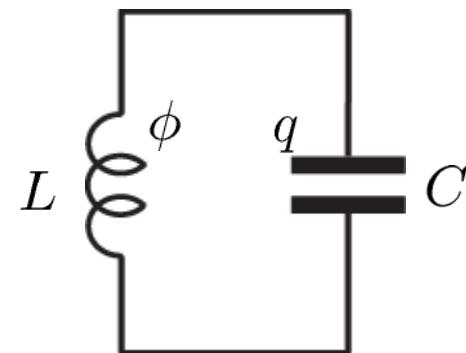


Constructing Linear Quantum Electronic Circuits

basic circuit elements:



harmonic LC oscillator:



$$\omega = \frac{1}{\sqrt{LC}} \sim 5 \text{ GHz}$$

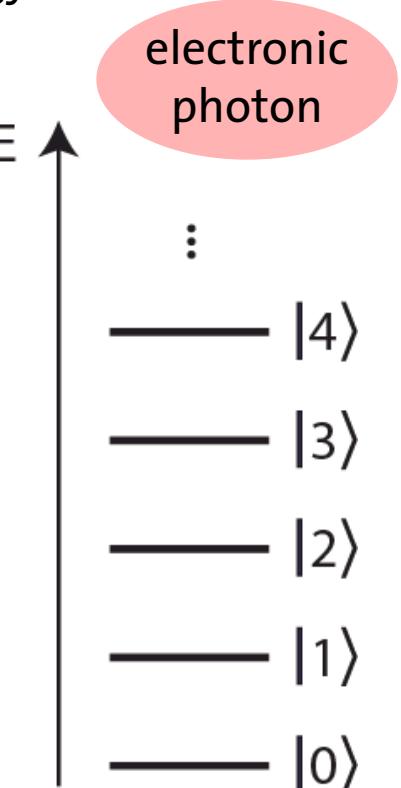
classical physics:

$$H = \frac{\phi^2}{2L} + \frac{q^2}{2C}$$

quantum mechanics:

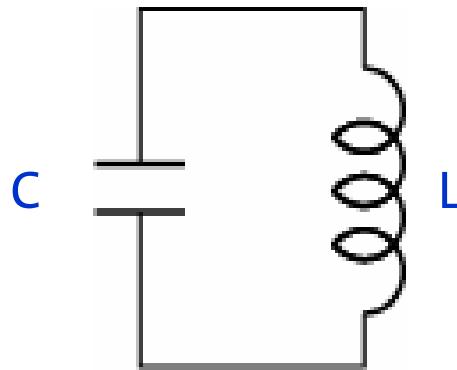
$$\hat{H} = \frac{\hat{\phi}^2}{2L} + \frac{\hat{q}^2}{2C} = \hbar\omega(\hat{a}^\dagger\hat{a} + \frac{1}{2}) \quad [\hat{\phi}, \hat{q}] = i\hbar$$

energy:



Superconducting Harmonic Oscillators

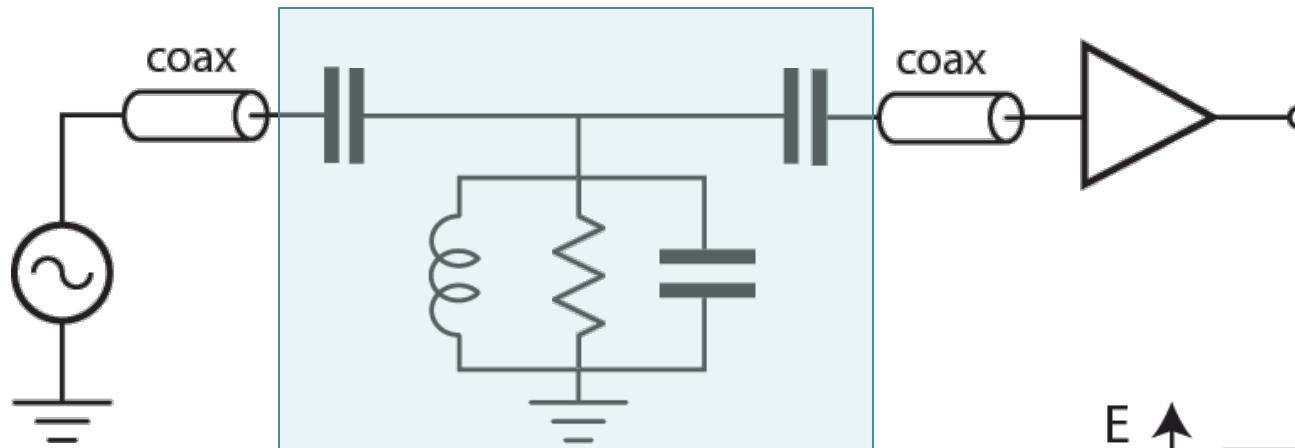
a simple electronic circuit:



- typical inductor: $L = 1 \text{ nH}$
- a wire in vacuum has inductance $\sim 1 \text{ nH/mm}$
- typical capacitor: $C = 1 \text{ pF}$
- a capacitor with plate size $10 \mu\text{m} \times 10 \mu\text{m}$ and dielectric AlO_x ($\epsilon = 10$) of thickness 10 nm has a capacitance $C \sim 1 \text{ pF}$
- resonance frequency

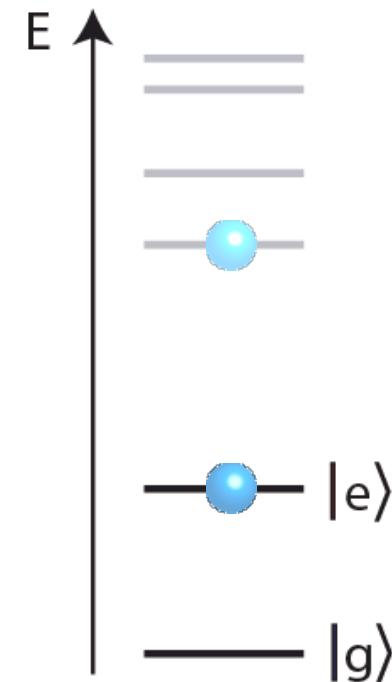
$$\frac{1}{2\pi\sqrt{LC}} \sim 5 \text{ GHz}$$

How to Operate Circuits Quantum Mechanically?



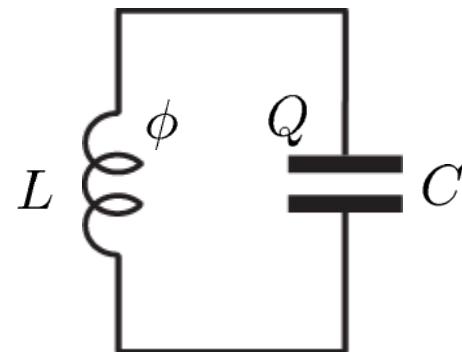
recipe:

- avoid dissipation
- work at low temperatures
- isolate quantum circuit from environment



Quantization of an Electronic Harmonic Oscillator

Harmonic LC oscillator:



$$\begin{aligned}Q &= CV && \text{Charge on capacitor} \\ \phi &= LI && \text{Flux in inductor} \\ V &= -L\dot{I} = -\dot{\phi} && \text{Voltage across inductor}\end{aligned}$$

Classical Hamiltonian:

$$H = \frac{CV^2}{2} + \frac{LI^2}{2} = \frac{Q^2}{2C} + \frac{\phi^2}{2L}$$

Conjugate variables:

$$\begin{aligned}\frac{\partial H}{\partial \phi} &= \frac{\phi}{L} = I = \dot{Q} \\ \frac{\partial H}{\partial Q} &= \frac{Q}{C} = V = -L\dot{I} = -\dot{\phi}\end{aligned}$$

Hamilton operator:

$$\hat{H} = \frac{\hat{\phi}^2}{2L} + \frac{\hat{Q}^2}{2C}$$

Flux and charge operator:

$$\begin{aligned}\hat{\phi} &= \phi \\ \hat{Q} &= -i\hbar \frac{\partial}{\partial \phi}\end{aligned}$$

Commutation relation:

$$[\hat{\phi}, \hat{Q}] = i\hbar$$

Creation and Annihilation Operators for Circuits

Hamilton operator of harmonic oscillator in second quantization:

$$\hat{H} = \frac{\hat{\phi}^2}{2L} + \frac{\hat{Q}^2}{2C} = \hbar\omega(\hat{a}^\dagger\hat{a} + 1/2)$$

$$\hat{a}^\dagger |n\rangle = \sqrt{n+1} |n+1\rangle \quad \text{Creation operator}$$

$$\hat{a} |n\rangle = \sqrt{n} |n-1\rangle \quad \text{Annihilation operator}$$

$$\hat{a}^\dagger\hat{a} |n\rangle = n |n\rangle \quad \text{Number operator}$$

$$\hat{Q} = \sqrt{\frac{\hbar}{2Z_C}}(\hat{a}^\dagger + \hat{a}) \quad \text{Charge/voltage operator}$$

$$\hat{\phi} = i\sqrt{\frac{\hbar Z_C}{2}}(\hat{a}^\dagger - \hat{a}) \quad \text{Flux/current operator}$$

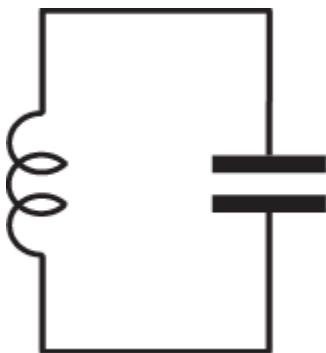
$$\hat{V} = \frac{\hat{Q}}{C} \quad \hat{I} = \frac{\hat{\phi}}{L}$$

With characteristic impedance:

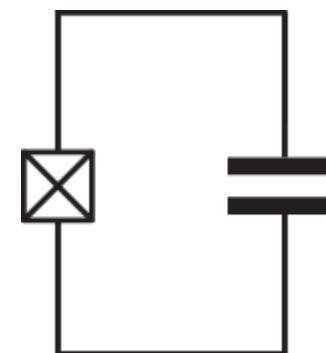
$$Z_C = \sqrt{\frac{L}{C}}$$

Linear vs. Nonlinear Superconducting Oscillators

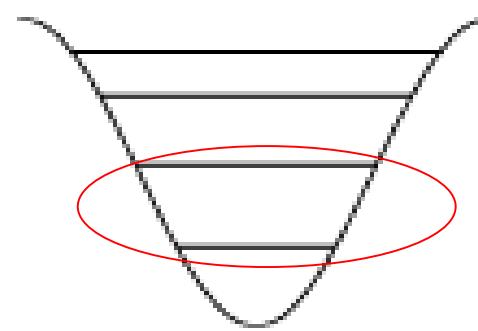
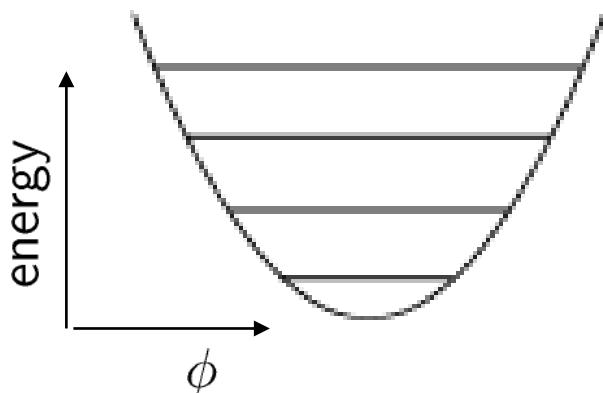
LC resonator:



Josephson junction resonator:
Josephson junction = nonlinear inductor

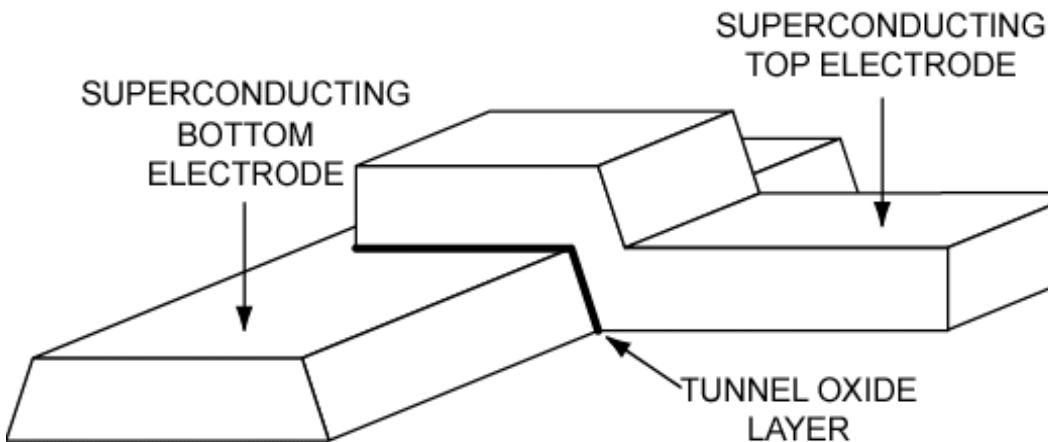


anharmonicity defines effective two-level system



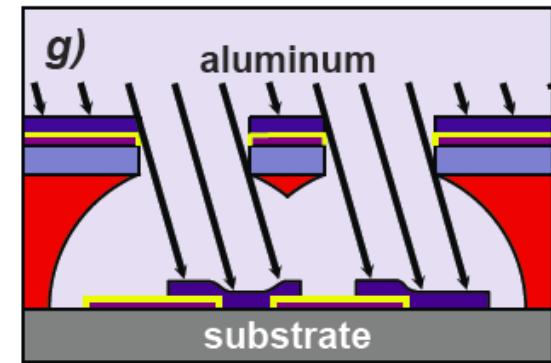
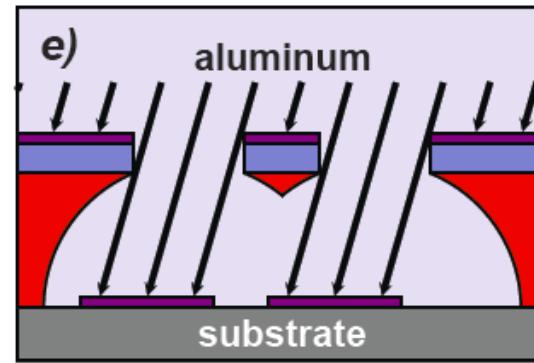
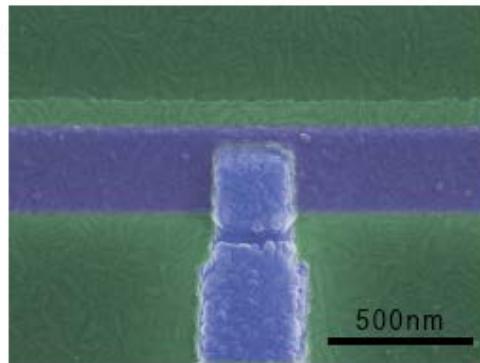
A Low-Loss Nonlinear Element

a (superconducting) Josephson junction:



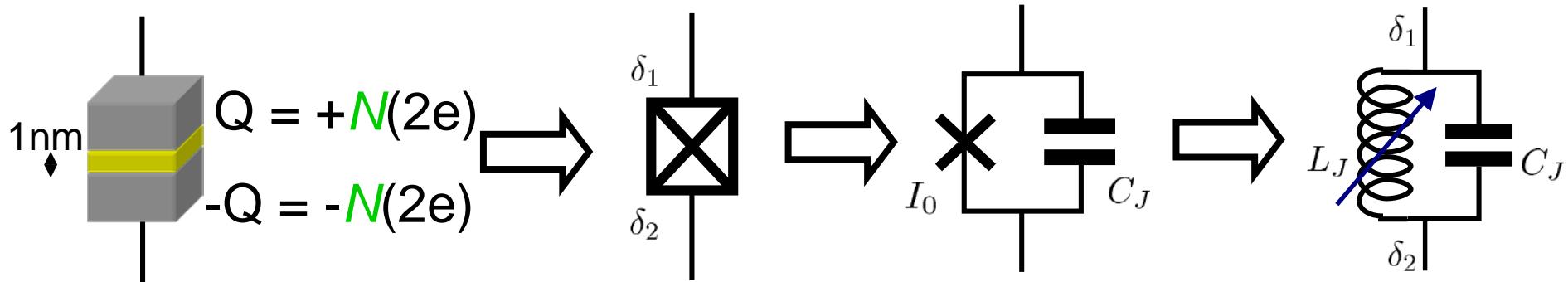
- superconductors: Nb, Al
- tunnel barrier: AlO_x

Josephson junction fabricated by shadow evaporation:



Josephson Tunnel Junction

The only non-linear resonator with no dissipation (BCS, $k_B T < \Delta$)



Tunnel junction parameters:

- Critical current I_0
- Junction capacitance C_J
- Internal resistance R_J

Josephson relations:

$$I = I_0 \sin \delta$$

$$V = \frac{\phi_0}{2\pi} \dot{\delta}$$

Flux quantum:

$$\phi_0 = \frac{h}{2e}$$

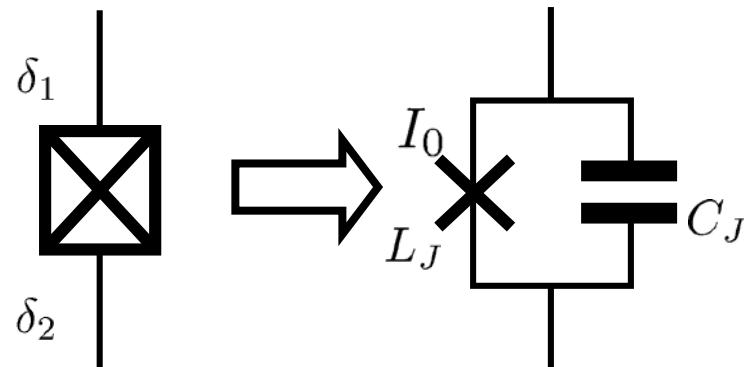
Phase difference:

$$\delta = \delta_2 - \delta_1$$

derivation of Josephson effect, see e.g.: chap. 21 in R. A. Feynman: Quantum mechanics, The Feynman Lectures on Physics. Vol. 3 (Addison-Wesley, 1965)

The Josephson Junction as an ideal Non-Linear Inductor

a nonlinear inductor without dissipation



Josephson relations:

$$I = I_0 \sin \delta = I_0 \sin [2\pi\phi(t)/\phi_0] \quad \text{nonlinear current/phase relation}$$
$$V = \frac{\phi_0}{2\pi} \dot{\delta} = \dot{\phi}$$

gauge inv. phase difference:

$$\delta = \delta_2 - \delta_1 = 2\pi\phi(t)/\phi_0$$

Josephson inductance:

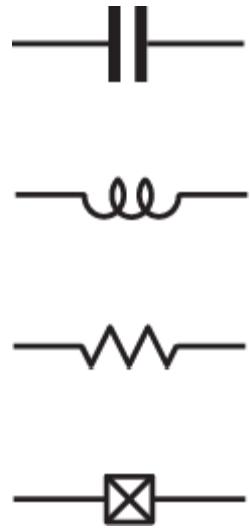
$$V = -L_J \dot{I} = \frac{\phi_0}{2\pi I_0} \frac{1}{\cos \delta} \dot{I}$$

Josephson energy:

$$E_J = \int V I dt = -\frac{I_0 \phi_0}{2\pi} \cos \delta$$

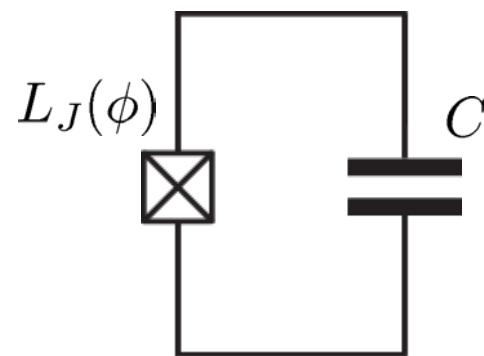
Constructing Non-Linear Quantum Electronic Circuits

circuit elements:



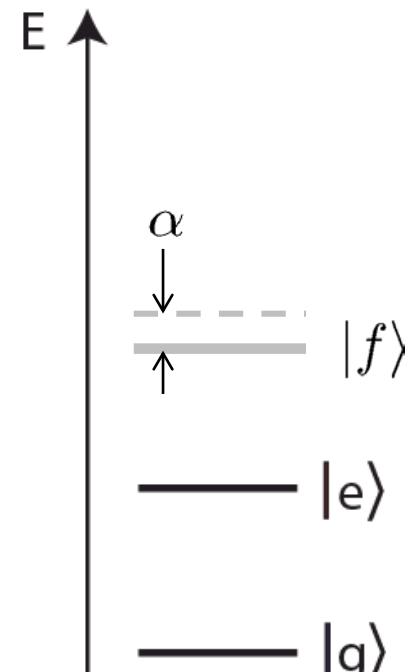
Josephson junction:
a non-dissipative nonlinear
element (inductor)

anharmonic oscillator:



$$H \approx \hbar(\omega_{ge} \hat{b}^\dagger \hat{b} - \frac{\alpha}{2} \hat{b}^\dagger 2 \hat{b}^2)$$

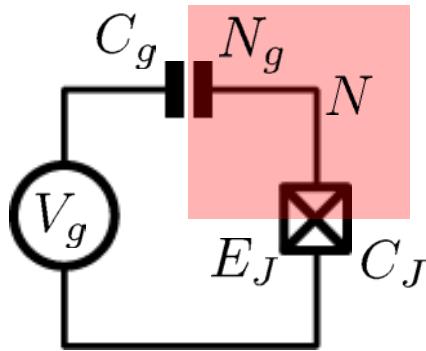
non-linear energy
level spectrum:



electronic
artificial atom

The Cooper Pair Box Qubit

A Charge Qubit: The Cooper Pair Box



discrete charge on island:

$$N = \frac{Q}{2e}$$

continuous gate charge:

$$N_g = \frac{C_g V_g}{2e}$$

total box capacitance

$$C_{\Sigma} = C_g + C_J$$

Hamiltonian: $H = H_{\text{el}} + H_{\text{mag}}$

electrostatic part:

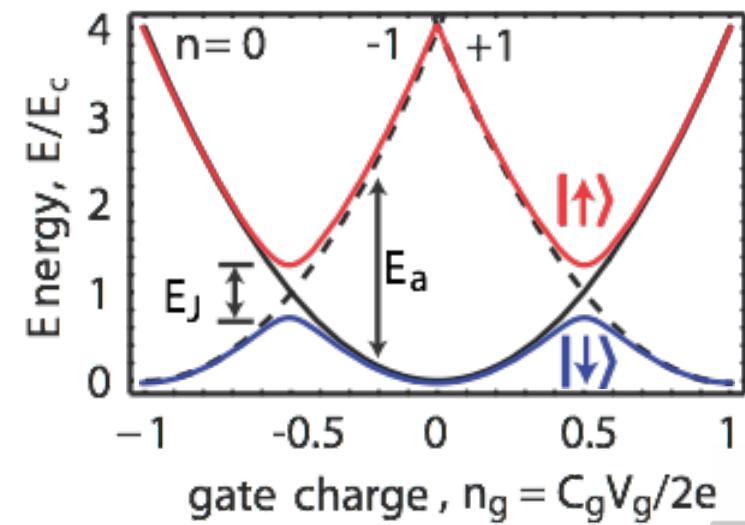
$$H_{\text{el}} = \frac{Q^2}{2C} = \frac{(2e)^2}{2C_{\Sigma}} (N - N_g)^2$$

charging energy E_C

magnetic part:

$$H_{\text{mag}} = -E_J \cos \delta \approx \frac{\phi^2}{2L_{J0}}$$

Josephson energy



Hamilton Operator of the Cooper Pair Box

Hamiltonian: $\hat{H} = \hat{H}_{\text{el}} + \hat{H}_{\text{mag}} = E_C(\hat{N} - N_g)^2 - E_J \cos \hat{\delta}$

commutation relation: $[\hat{\delta}, \hat{N}] = i$ $\cos \hat{\delta} = \frac{1}{2}(e^{i\hat{\delta}} + e^{-i\hat{\delta}})$

charge number operator: $\hat{N}|N\rangle = N|N\rangle$ eigenvalues, eigenfunctions

$$\sum_N |N\rangle\langle N| = 1 \quad \text{completeness}$$

$$\langle N|M\rangle = \delta_{NM} \quad \text{orthogonality}$$

phase basis:

$$|\delta\rangle = \frac{1}{\sqrt{2\pi}} \sum_N e^{iN\delta} |N\rangle \quad \text{basis transformation}$$

$$e^{\pm i\hat{\delta}} |N\rangle = |N \pm 1\rangle$$

Solving the Cooper Pair Box Hamiltonian

Hamilton operator in the charge basis N :

$$\hat{H} = \sum_N \left[E_C(N - N_g)^2 |N\rangle\langle N| - \frac{E_J}{2} (|N\rangle\langle N+1| + |N+1\rangle\langle N|) \right]$$

solutions in the charge basis:

$$\hat{H}|\psi_n(N)\rangle = E_n|\psi_n(N)\rangle$$

Hamilton operator in the phase basis δ :

$$\hat{H} = E_C(\hat{N} - N_g)^2 - E_J \cos \hat{\delta} = E_C(-i \frac{\partial}{\partial \delta} - N_g)^2 - E_J \cos \hat{\delta}$$

transformation of the number operator:

$$\hat{N} = \frac{\hat{Q}}{2e} = -i\hbar \frac{1}{2e} \frac{\partial}{\partial \phi} = -i \frac{\partial}{\partial \delta}$$

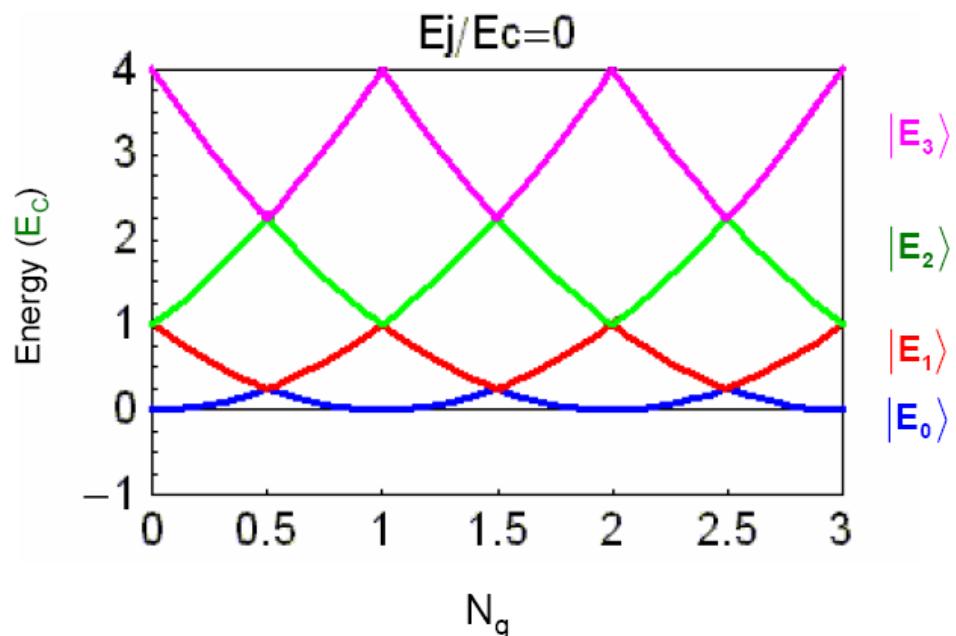
solutions in the phase basis:

$$\hat{H}|\psi_n(\delta)\rangle = E_n|\psi_n(\delta)\rangle$$

Energy Levels

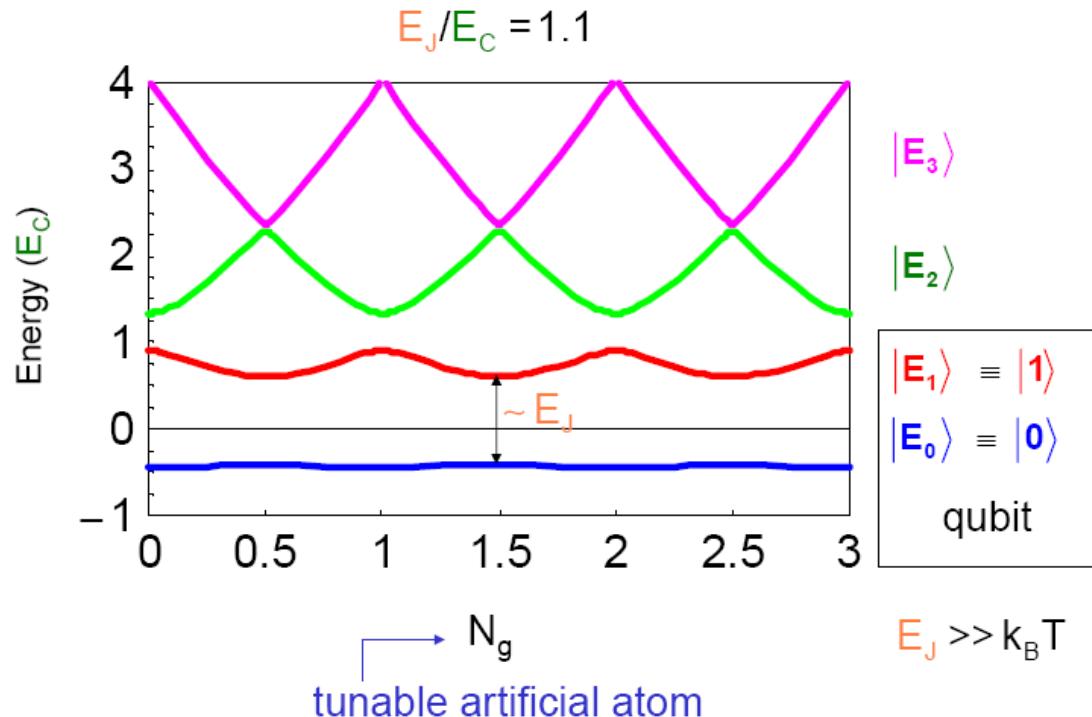
energy level diagram for $E_J=0$:

- energy bands are formed
- bands are periodic in N_g

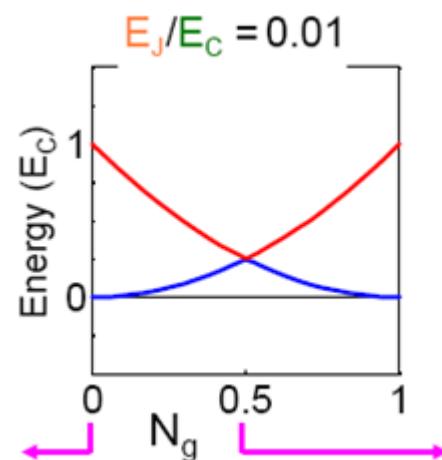
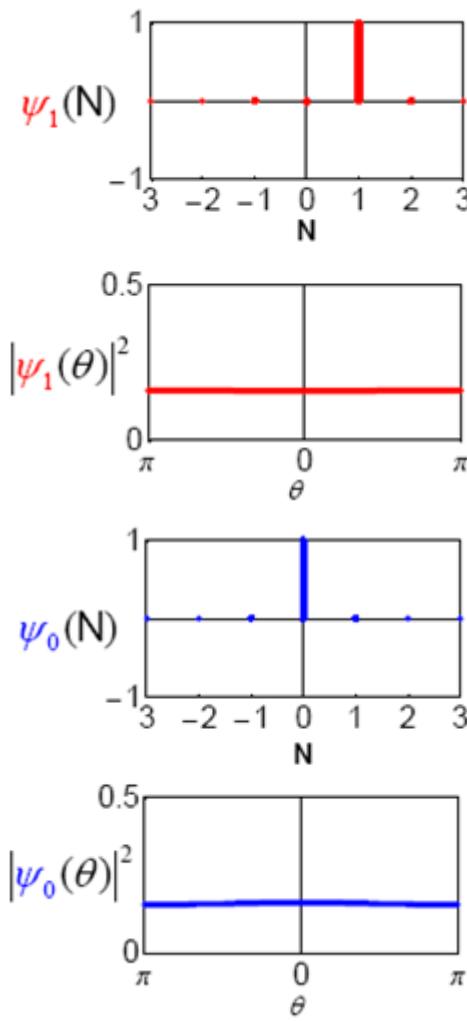


energy bands for finite E_J

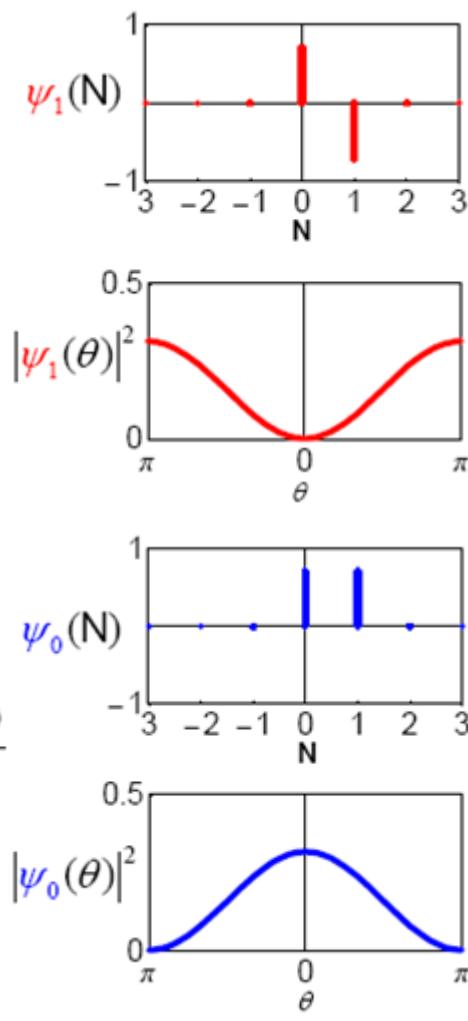
- Josephson coupling lifts degeneracy
- E_J scales level separation at charge degeneracy



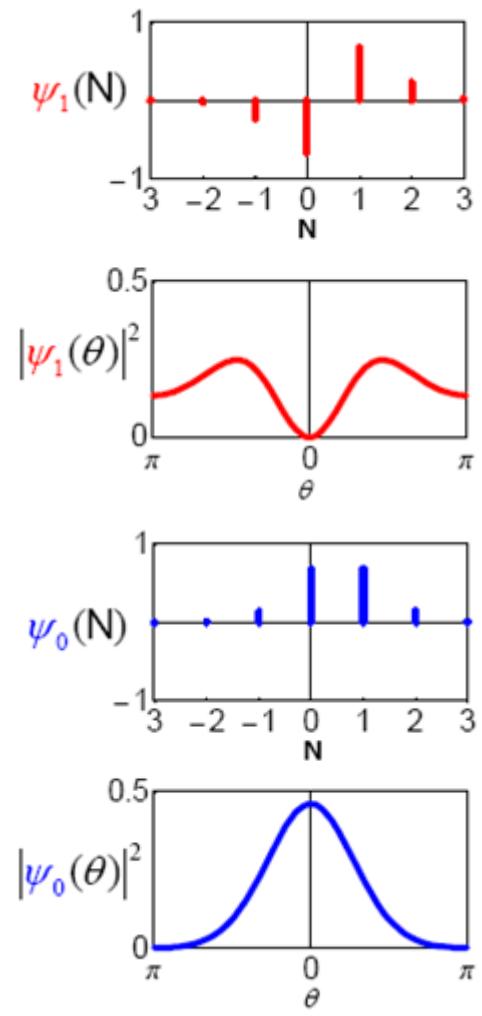
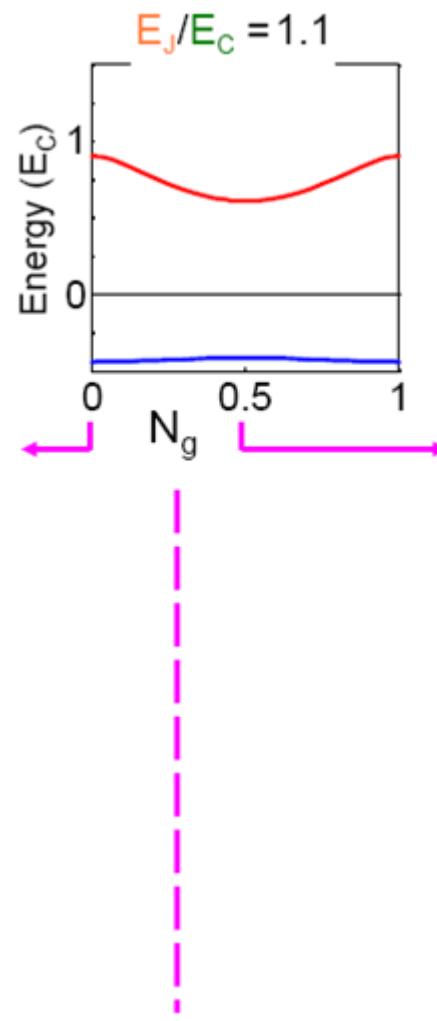
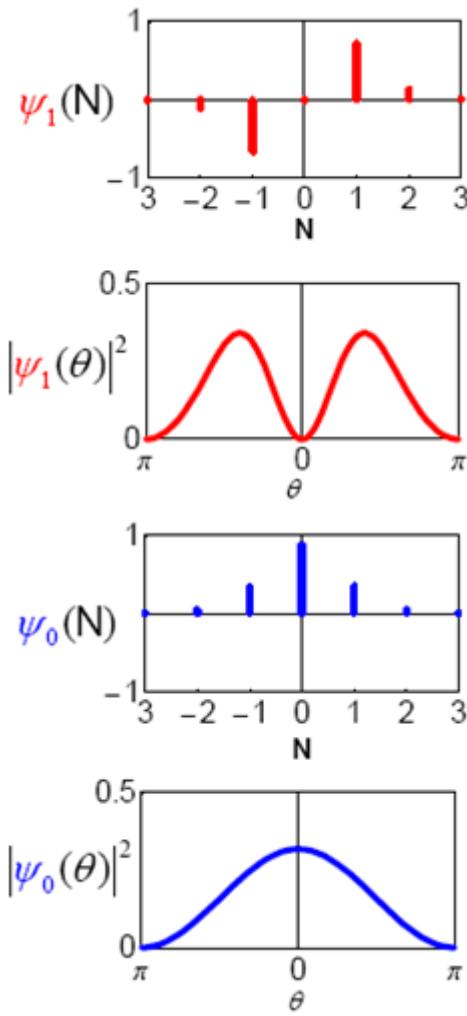
Charge and Phase Wave Functions ($E_J \ll E_C$)



$$\begin{aligned} |\psi_1\rangle &\approx |N=1\rangle & |\psi_1\rangle &\approx \frac{|N=0\rangle - |N=1\rangle}{\sqrt{2}} \\ |\psi_0\rangle &\approx |N=0\rangle & |\psi_0\rangle &\approx \frac{|N=0\rangle + |N=1\rangle}{\sqrt{2}} \end{aligned}$$



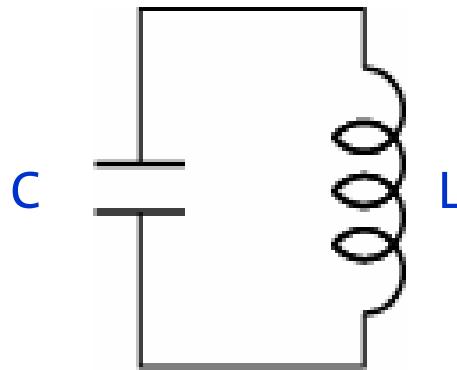
Charge and Phase Wave Functions ($E_J \sim E_C$)



Realizations of Harmonic Oscillators

Superconducting Harmonic Oscillators

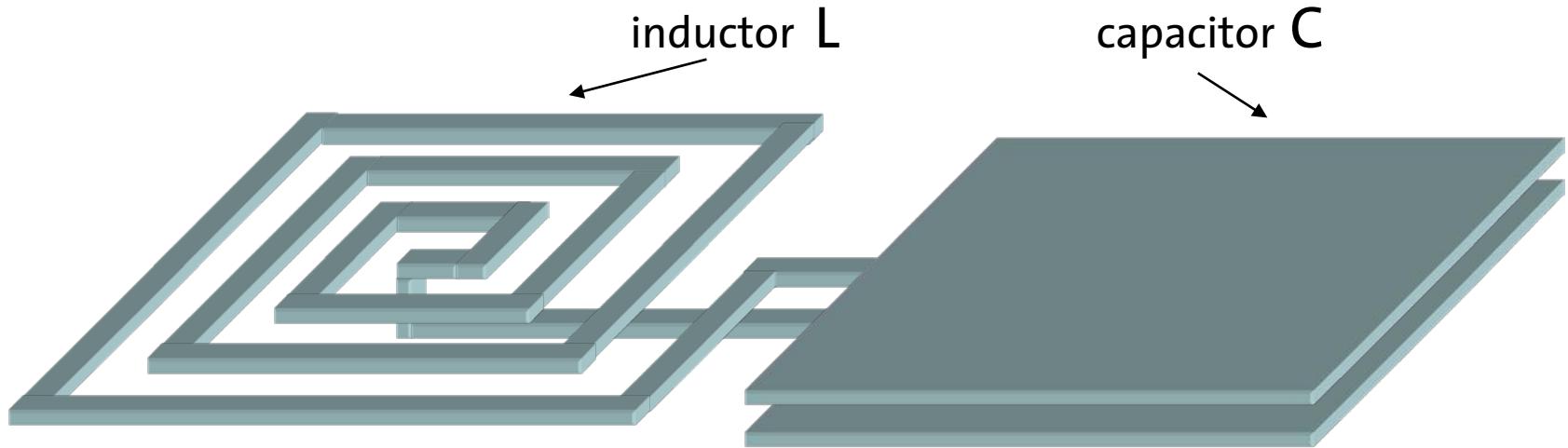
a simple electronic circuit:



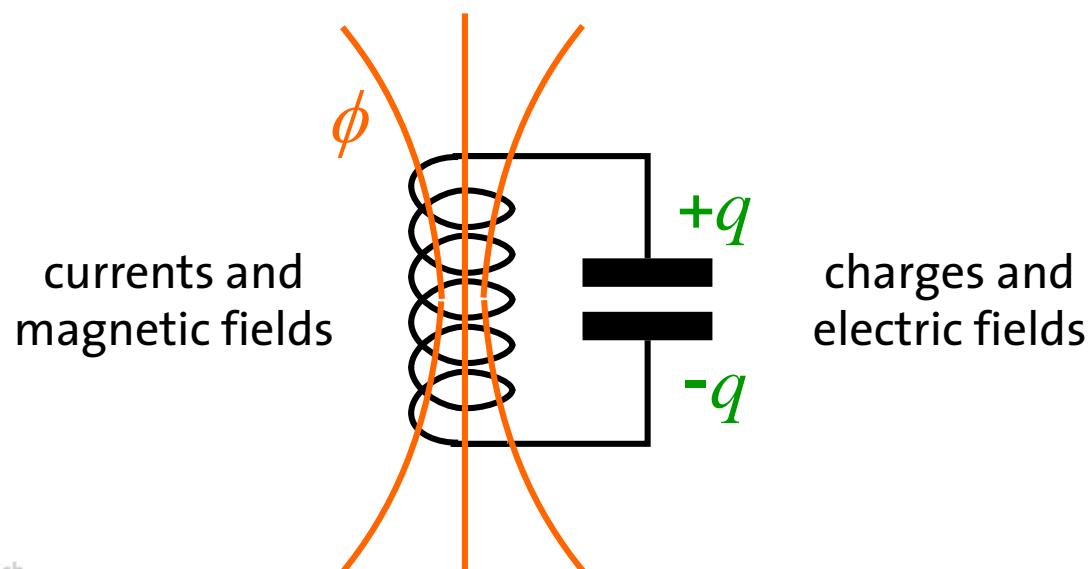
- typical inductor: $L = 1 \text{ nH}$
- a wire in vacuum has inductance $\sim 1 \text{ nH/mm}$
- typical capacitor: $C = 1 \text{ pF}$
- a capacitor with plate size $10 \mu\text{m} \times 10 \mu\text{m}$ and dielectric AlO_x ($\epsilon = 10$) of thickness 10 nm has a capacitance $C \sim 1 \text{ pF}$
- resonance frequency

$$\frac{1}{2\pi\sqrt{LC}} \sim 5 \text{ GHz}$$

Realization of H.O.: Lumped Element Resonator

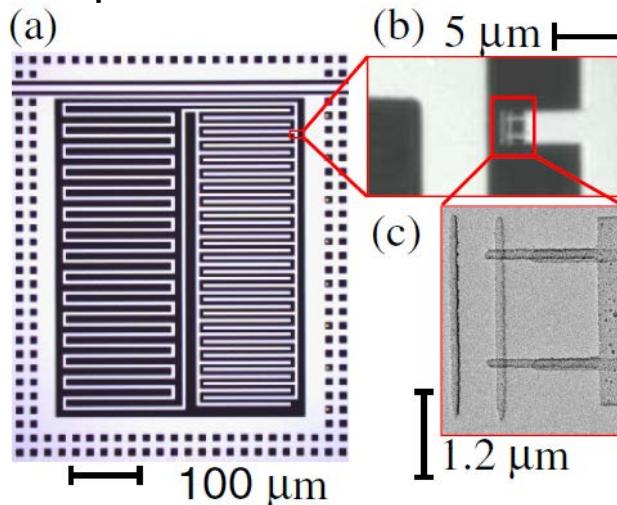


a harmonic oscillator



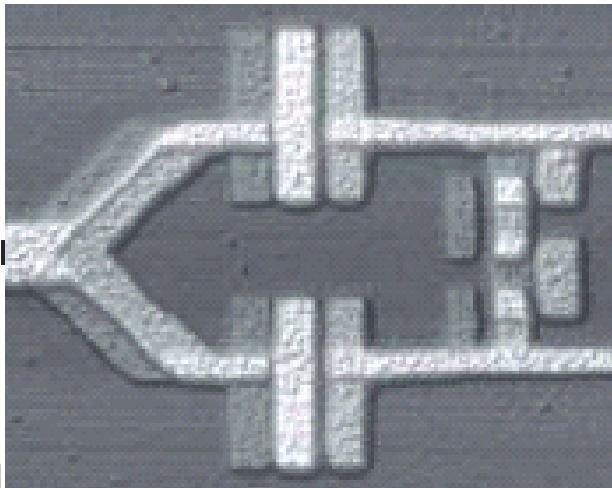
Types of Superconducting Harmonic Oscillators

lumped element resonator:



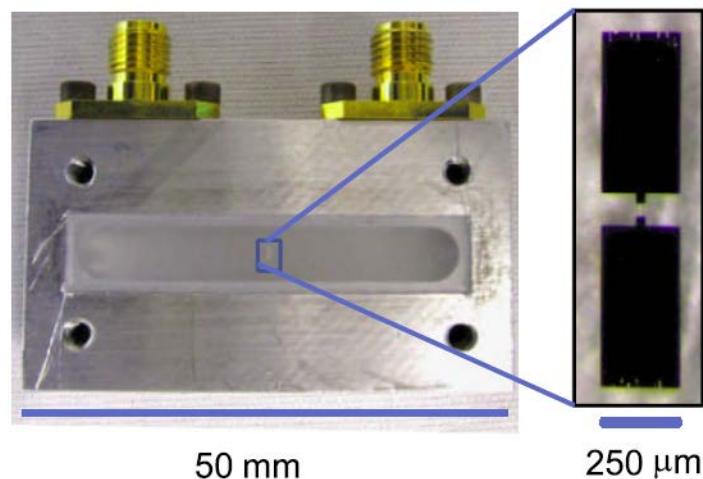
Z. Kim *et al.*, PRL 106, 120501 (2011)

weakly nonlinear junction:



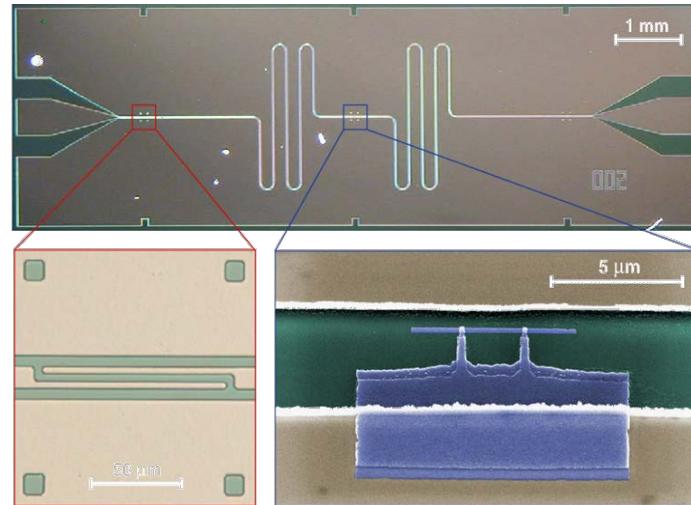
I. Chiorescu *et al.*, Nature 431, 159 (2004)

3D cavity:



H. Paik *et al.*, PRL 107, 240501 (2011)

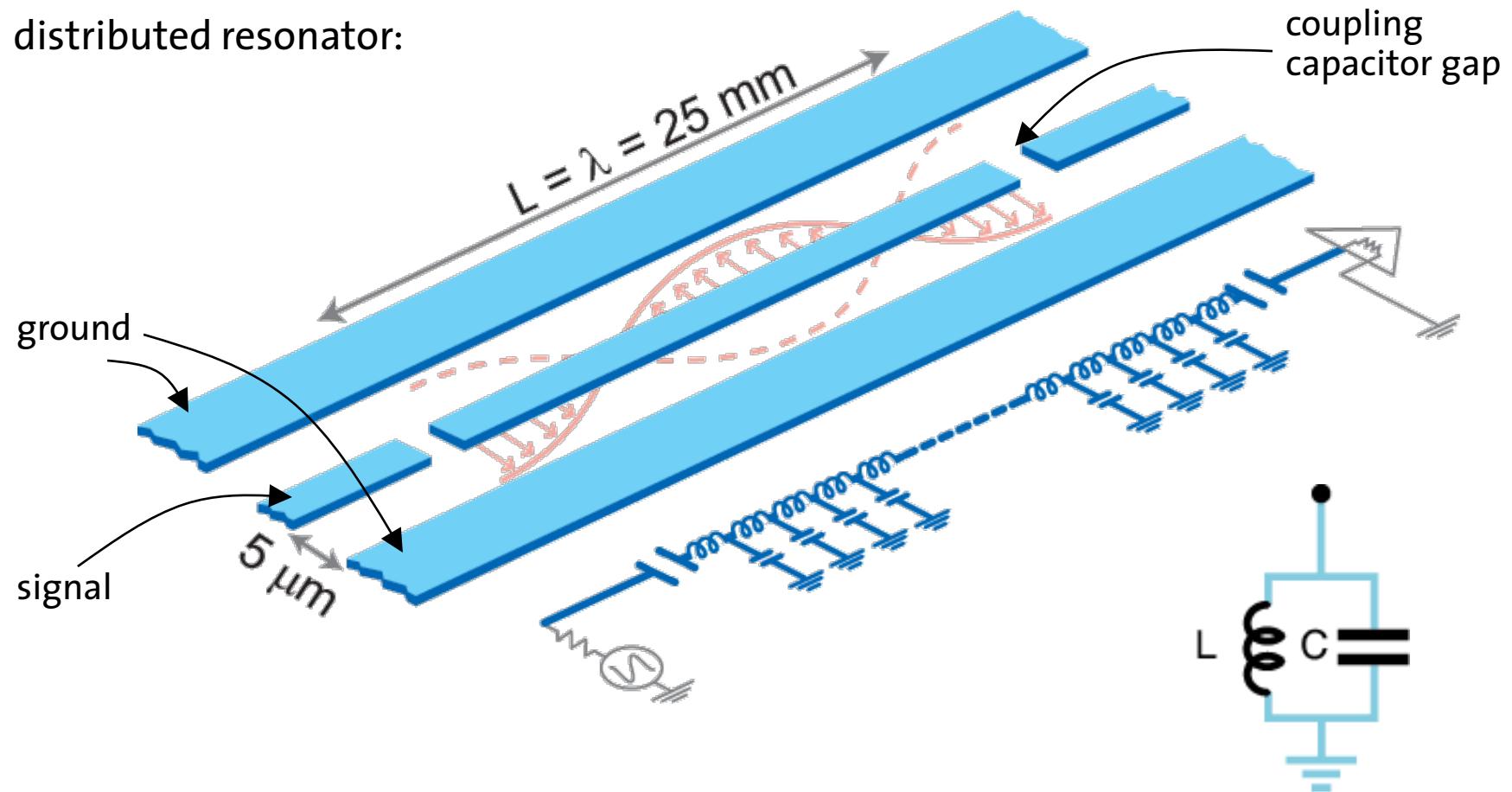
planar transmission line resonator:



A. Wallraff *et al.*, Nature 431, 162 (2004)

Realization of H.O.: Transmission Line Resonator

distributed resonator:

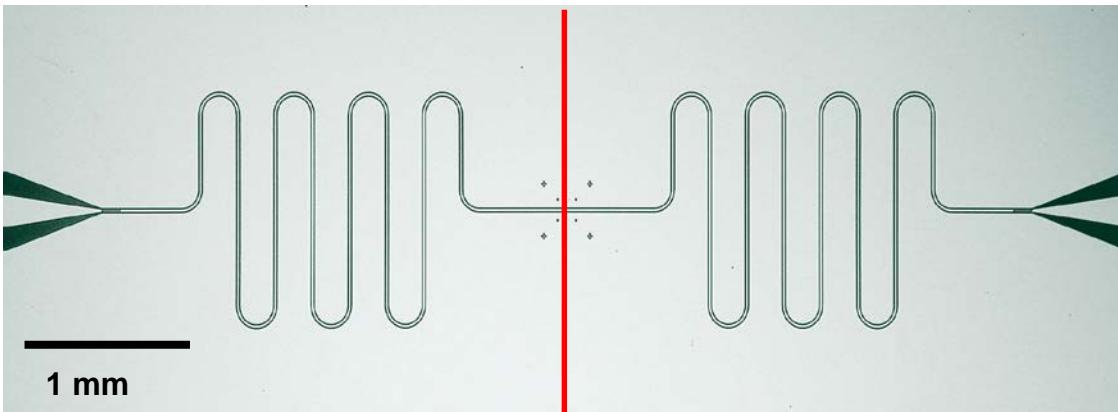


- coplanar waveguide resonator
- close to resonance: equivalent to lumped element LC resonator

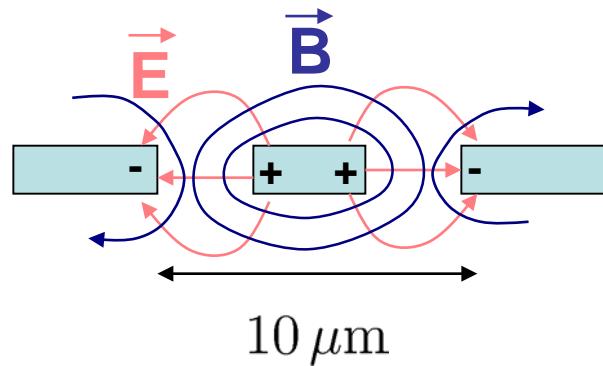
M. Goeppel *et al.*, Coplanar Waveguide Resonators
for Circuit QED, *Journal of Applied Physics* 104, 113904 (2008)

Realization of Transmission Line Resonator

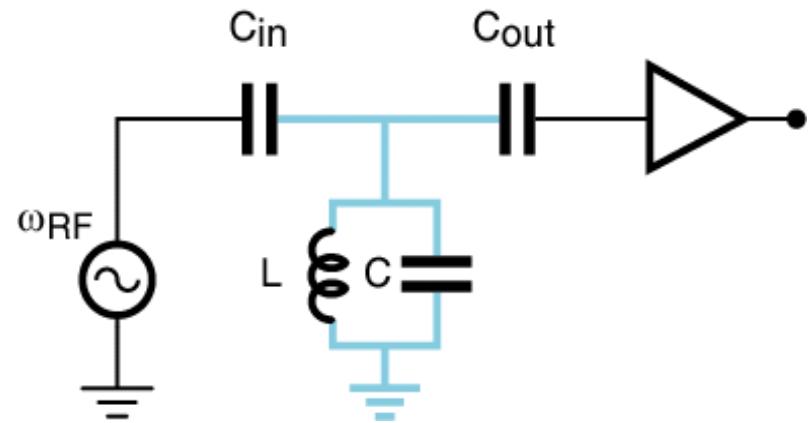
coplanar waveguide:



cross-section of transm. line
(TEM mode):

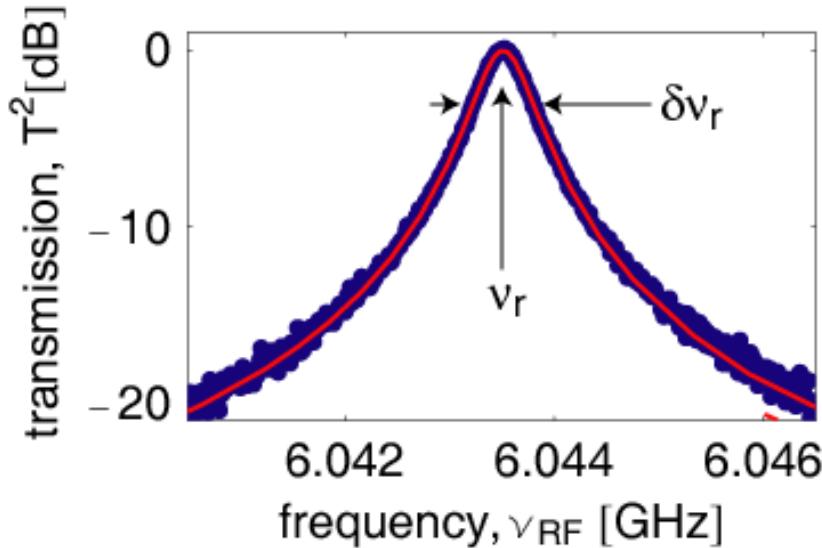


measuring the resonator:



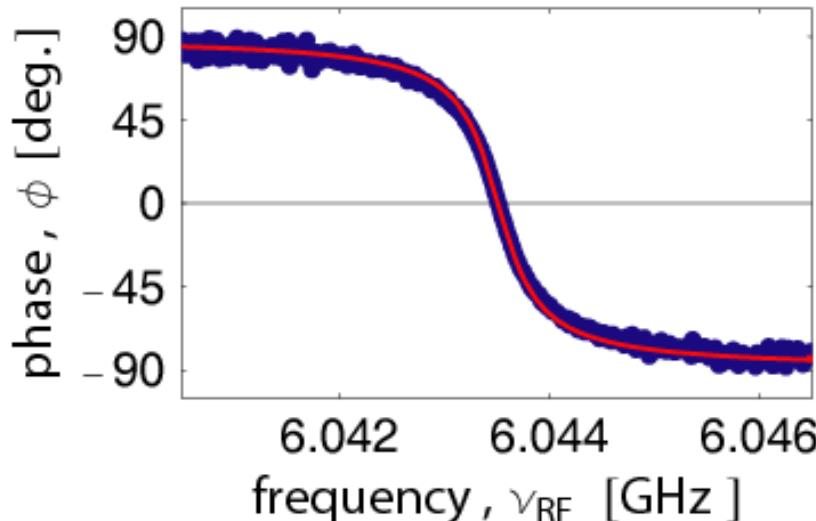
photon lifetime (quality factor) controlled
by coupling capacitors $C_{in/out}$

Resonator Quality Factor and Photon Lifetime



resonance frequency:

$$\nu_r = 6.04 \text{ GHz}$$



quality factor:

$$Q = \frac{\nu_r}{\delta\nu_r} \approx 10^4$$

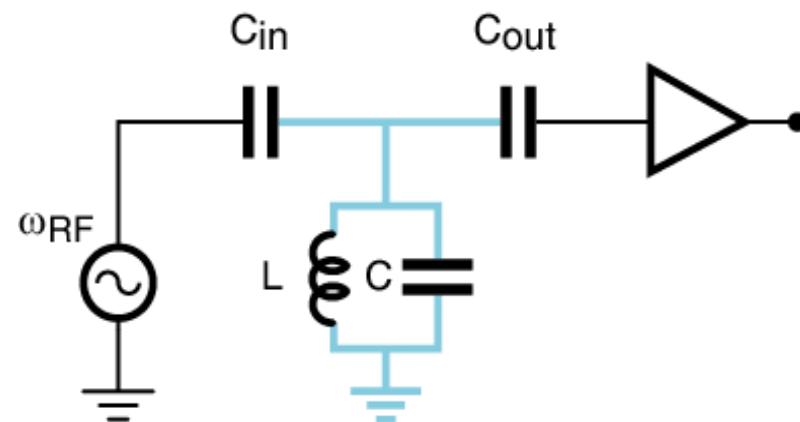
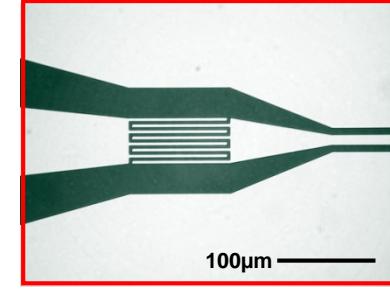
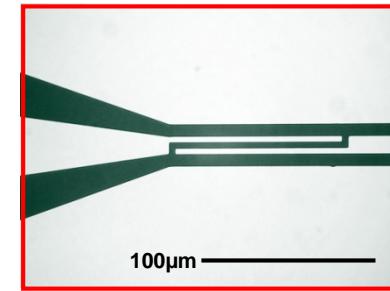
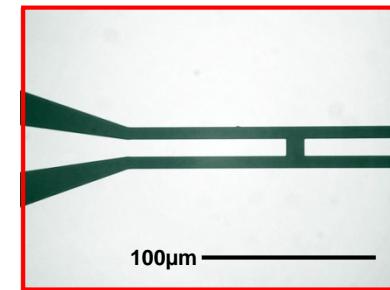
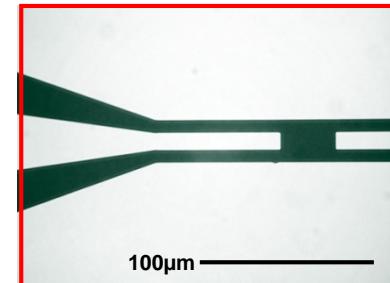
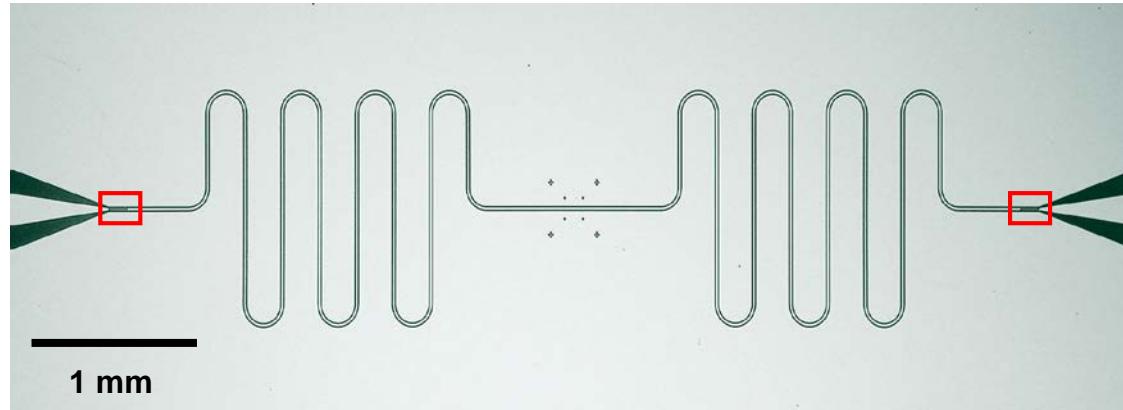
photon decay rate:

$$\frac{\kappa}{2\pi} = \frac{\nu_r}{Q} \approx 0.8 \text{ MHz}$$

photon lifetime:

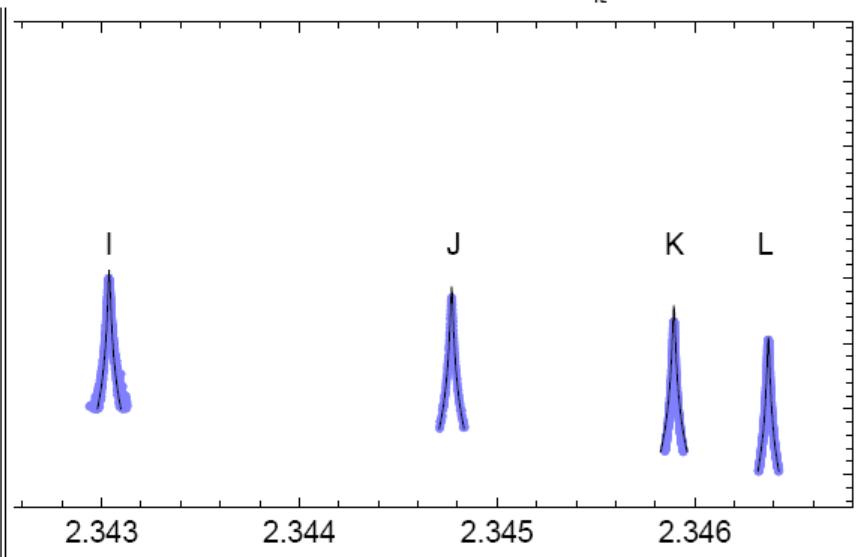
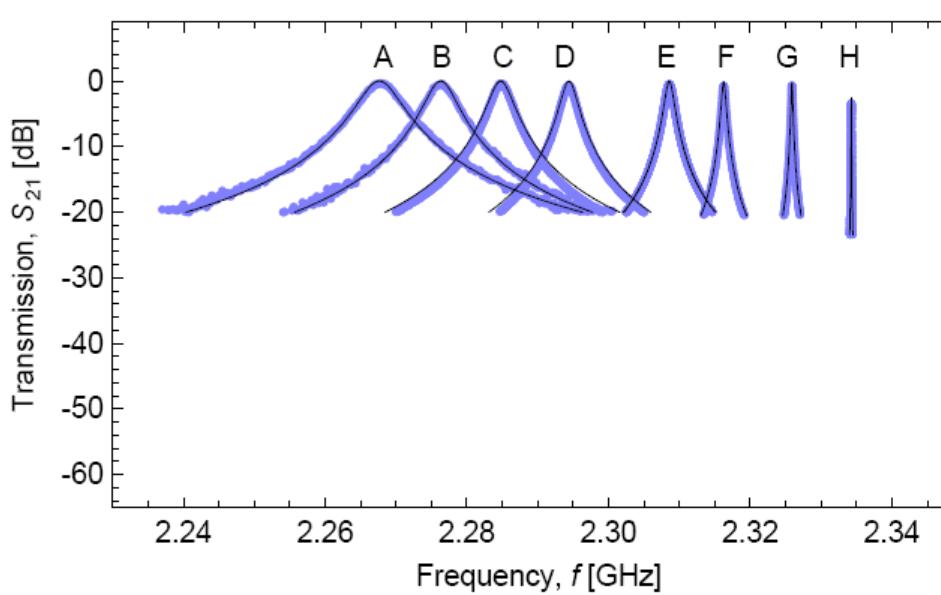
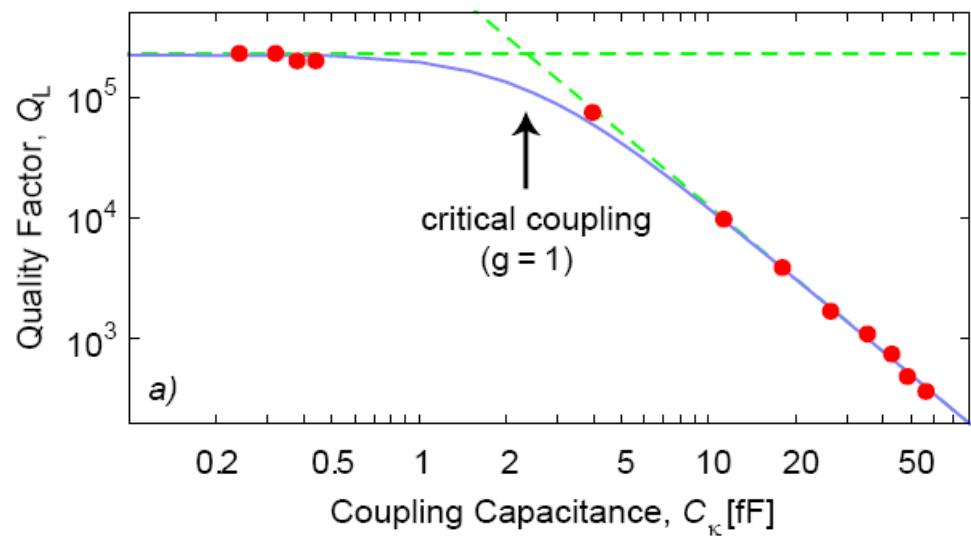
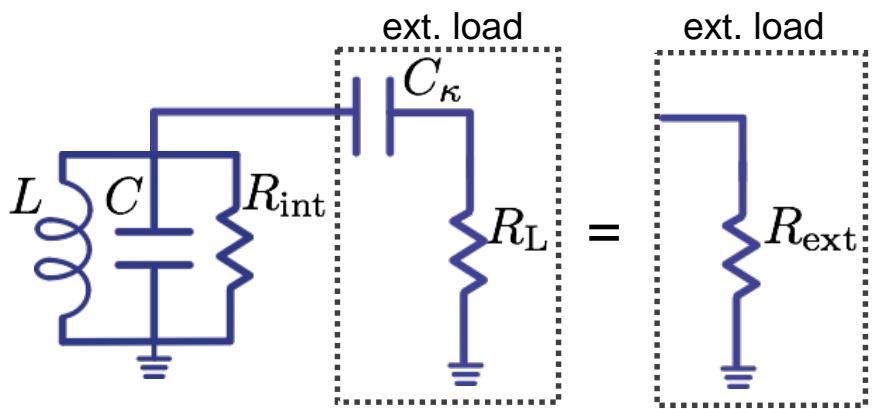
$$T_\kappa = 1/\kappa \approx 200 \text{ ns}$$

Controlling the Photon Life Time



photon lifetime (quality factor)
controlled by coupling capacitor $C_{in/out}$

Quality Factor Measurement

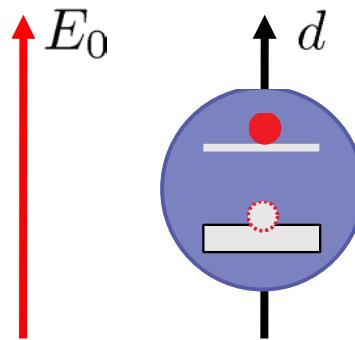


M. Goeppel *et al.*, *J. Appl. Phys.* 104, 113904 (2008)

Cavity Quantum Electrodynamics (QED): Coupling a Harmonic Oscillator to a Qubit

Investigating the Interaction of Light and Matter

challenging on the level of single (artificial) atoms and single photons



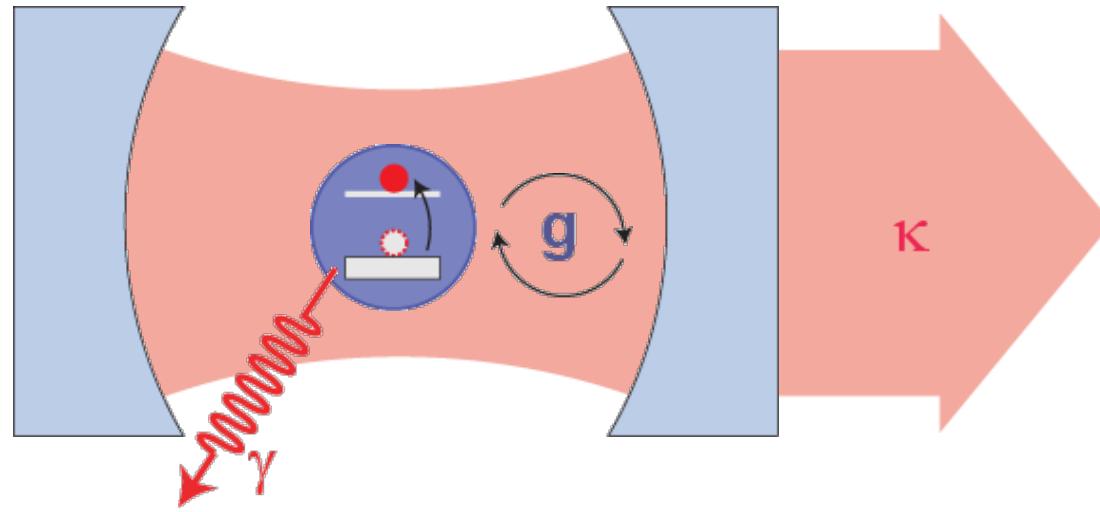
- mode-matching (controlling the absorption probability)
- single photon fields E_o (small in 3D)
- dipole moment d (usually small $\sim ea_o$)
- photon/dipole interaction $\hbar g \sim dE_0$ (usually small)

What to do?

- confine atom and photon in a cavity (cavity QED)
- engineer matter/light interactions, e.g. in solid state circuits

Cavity Quantum Electrodynamics

coupling photons to qubits:



Jaynes-Cummings Hamiltonian

$$H = \hbar\omega_r \left(a^\dagger a + \frac{1}{2} \right) + \frac{\hbar\omega_a}{2} \sigma^z + \hbar g(a^\dagger \sigma^- + a \sigma^+) + H_\kappa + H_\gamma$$

strong coupling limit ($g = dE_0/\hbar > \gamma, \kappa, 1/t_{\text{transit}}$)

Dressed States Energy Level Diagram

$$H = \hbar\omega_r \left(a^\dagger a + \frac{1}{2} \right) + \frac{\hbar\omega_a}{2} \sigma^z + \hbar g(a^\dagger \sigma^- + a \sigma^+)$$

:

:

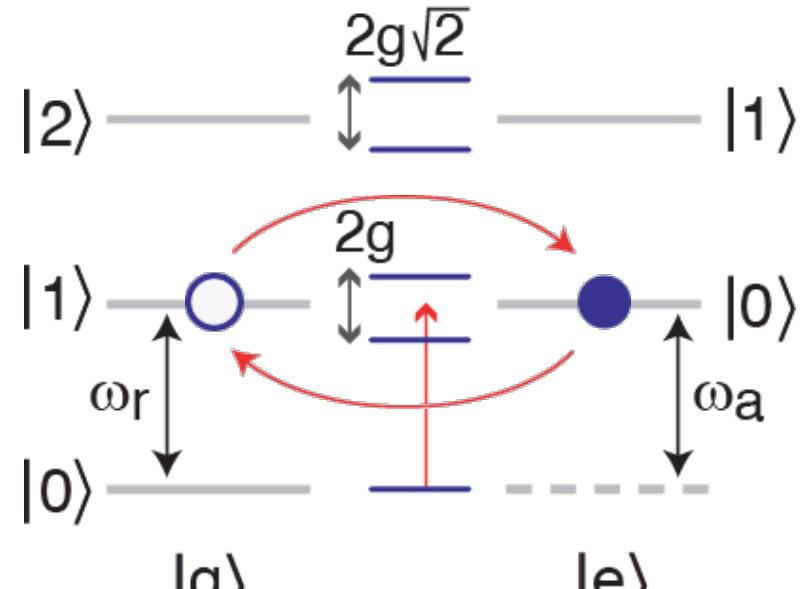
:

in resonance:

$$\omega_a - \omega_r = \Delta = 0$$

strong coupling limit:

$$g = \frac{dE_0}{\hbar} > \gamma, \kappa$$



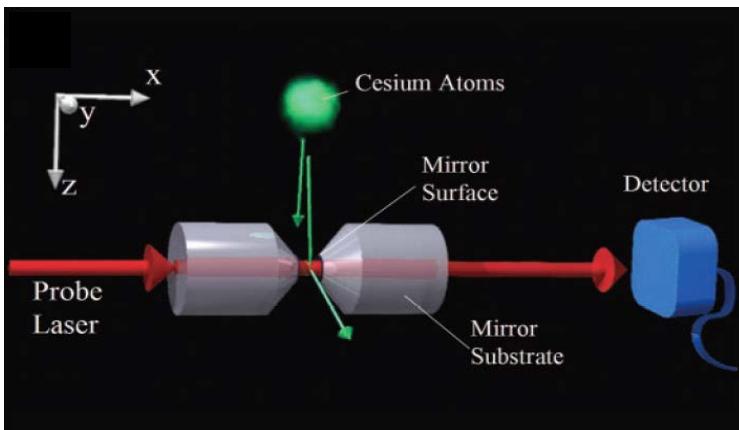
Jaynes-Cummings Ladder

Atomic cavity quantum electrodynamics reviews:

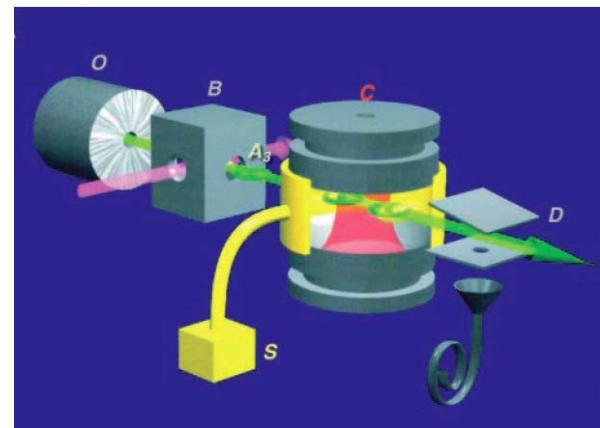
J. Ye., H. J. Kimble, H. Katori, *Science* **320**, 1734 (2008)

S. Haroche & J. Raimond, *Exploring the Quantum*, OUP Oxford (2006)

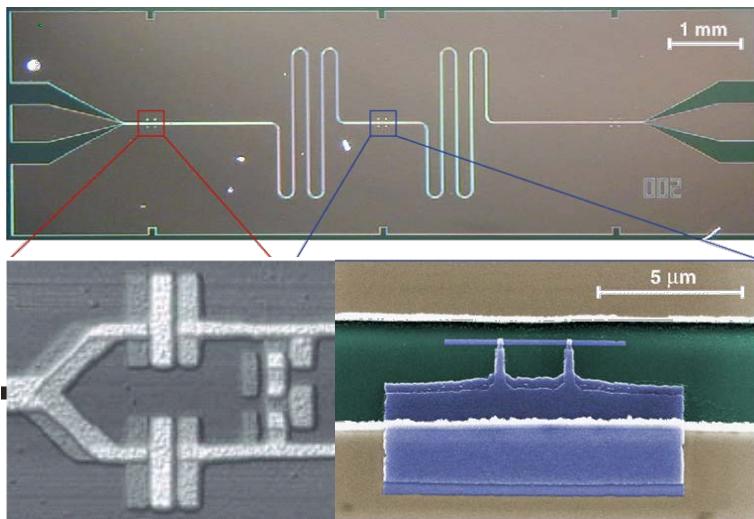
Systems for Exploring Cavity QED



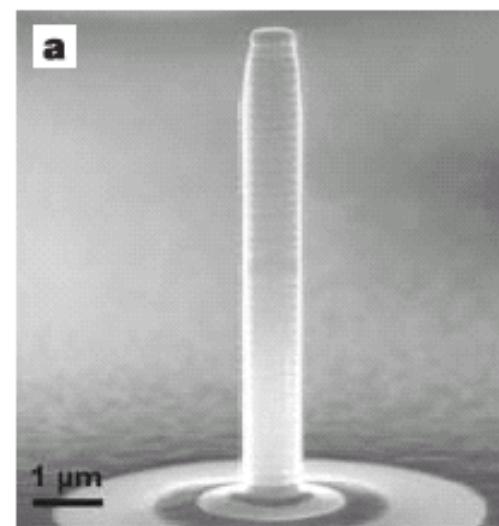
alkali atoms
MPQ, Caltech, ...



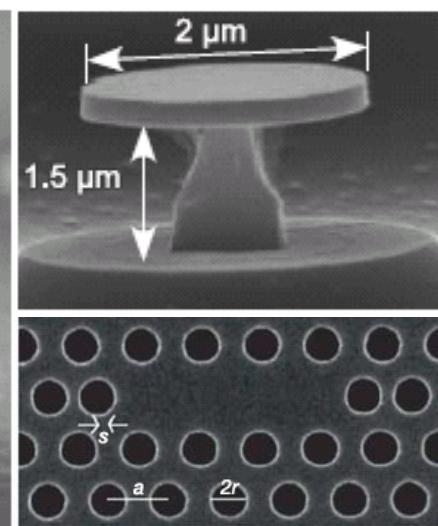
Rydberg atoms
ENS, MPQ, ...



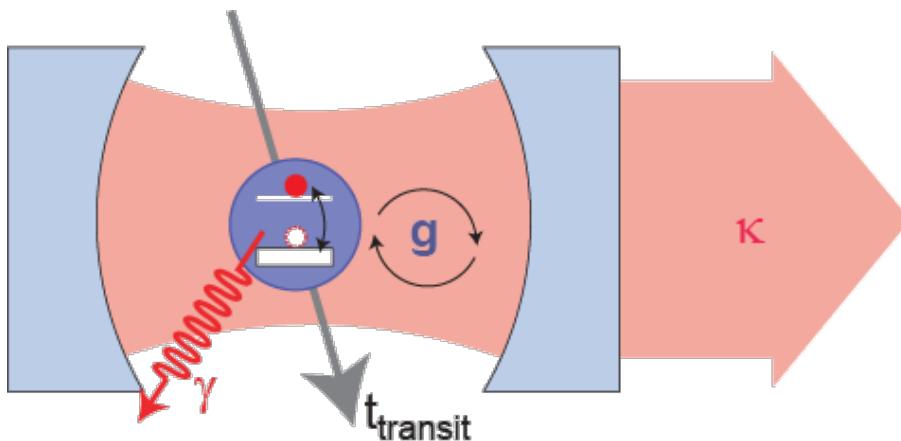
superconductor circuits
Yale, Delft, NTT, ETHZ, NIST, ...



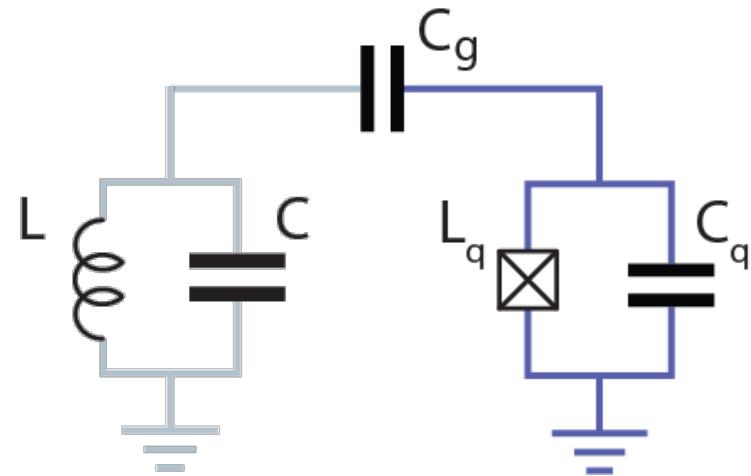
semiconductor quantum dots
Wurzburg, ETHZ, Stanford ...



Cavity QED with Superconducting Circuits



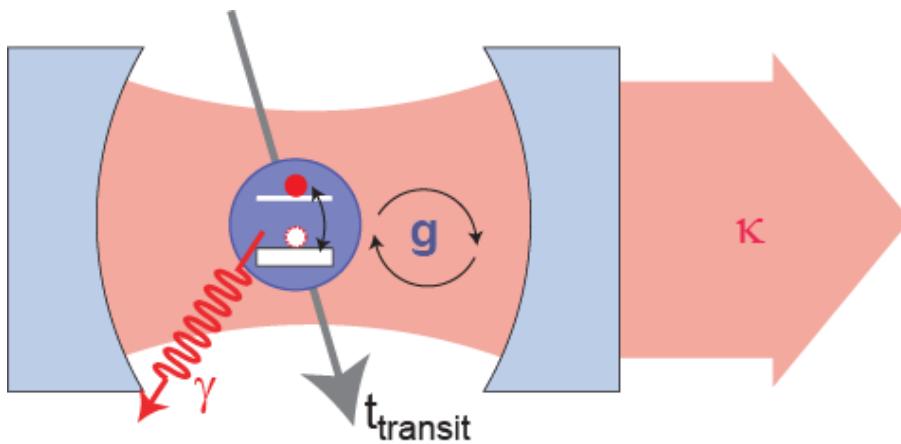
... basic approach:



What is this good for?

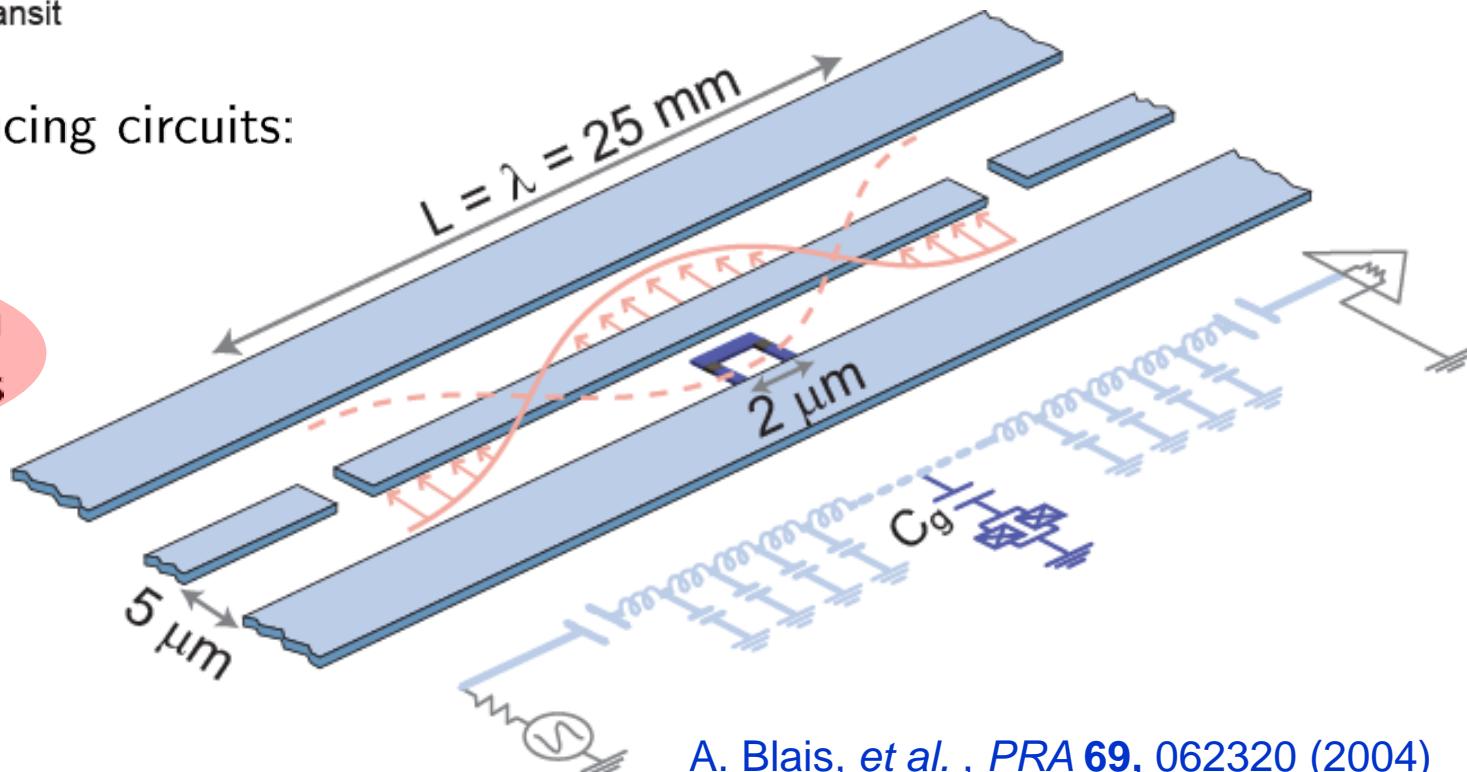
- Study matter light interaction
- Convert qubit states to photon states
- Use concepts to ...
 - ... build single photon sources and detectors
 - ... build quantum computers

Cavity QED with Superconducting Circuits



... in superconducting circuits:

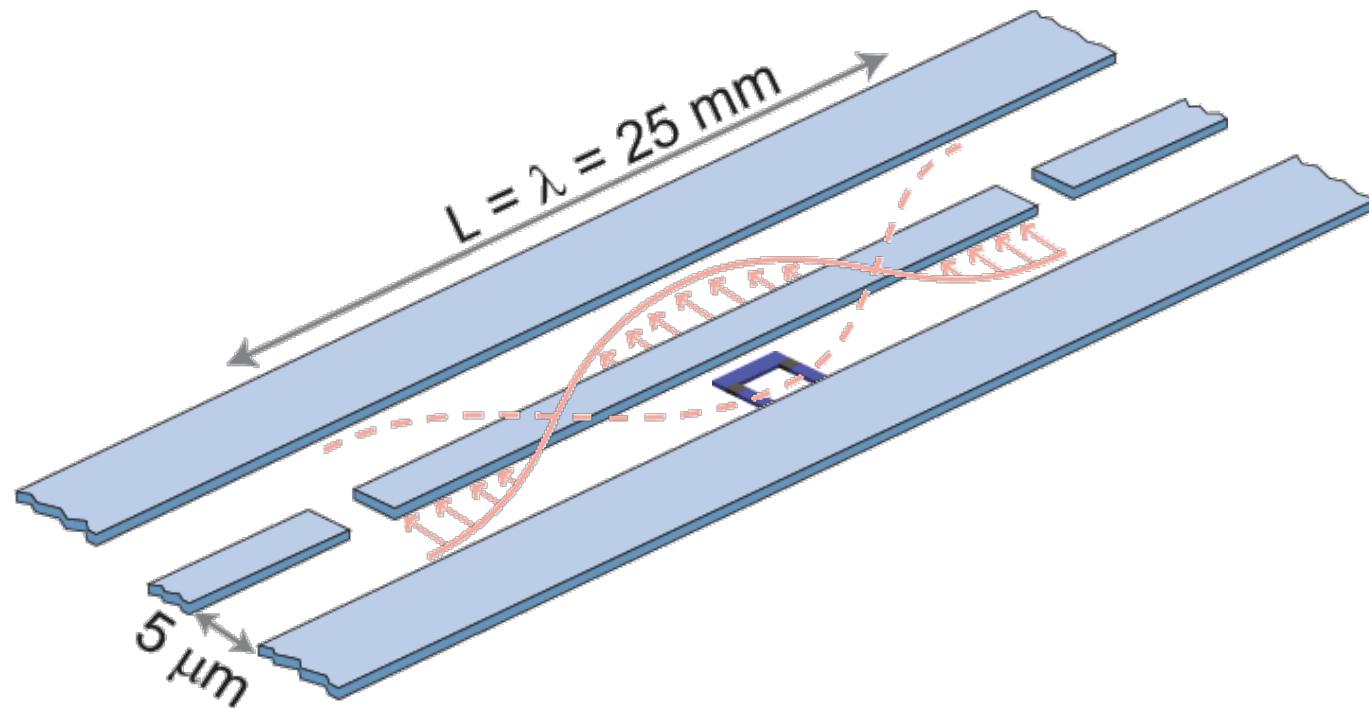
circuit quantum
electrodynamics



A. Blais, et al., PRA **69**, 062320 (2004)

A. Wallraff et al., Nature (London) **431**, 162 (2004)

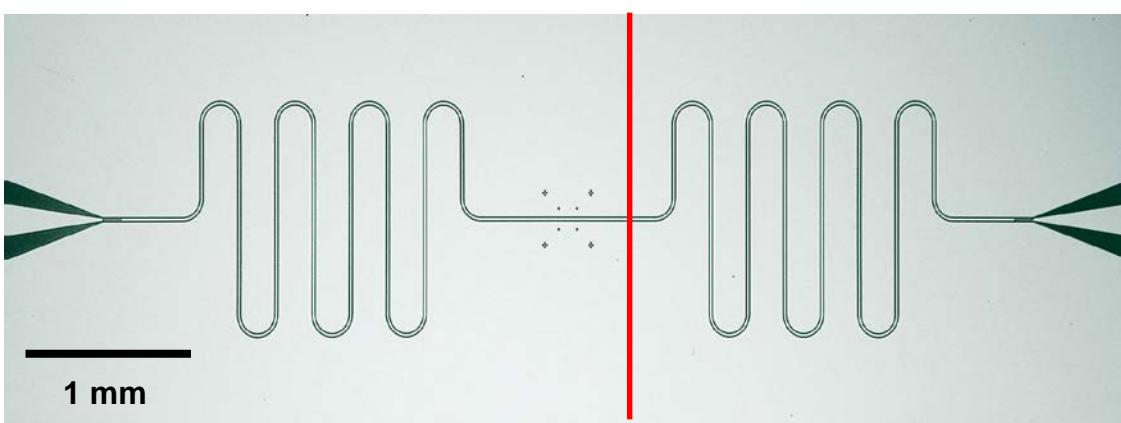
Circuit Quantum Electrodynamics



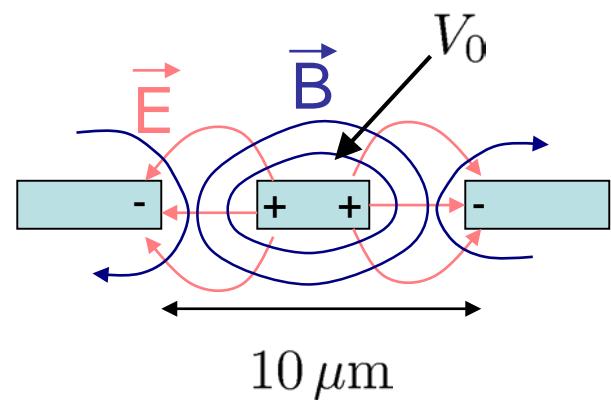
elements

- the cavity: a superconducting 1D transmission line resonator with **large vacuum field E_o** and **long photon life time $1/\kappa$**
- the artificial atom: a Cooper pair box with large E_J/E_C with **large dipole moment d** and **long coherence time $1/\gamma$**

Vacuum Field in 1D Cavity



cross-section
of transm. line (TEM mode):



voltage across resonator in vacuum state ($n = 0$)

$$V_{0,\text{rms}} = \sqrt{\frac{\hbar\omega_r}{2C}} \approx 1 \mu\text{V}$$

harmonic oscillator

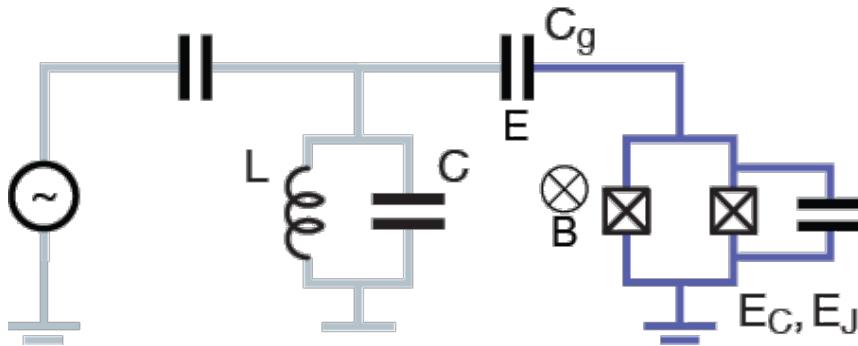
$$H_r = \hbar\omega_r \left(a^\dagger a + \frac{1}{2} \right)$$

$$E_0 = \frac{V_{0,\text{rms}}}{b} \approx 0.2 \text{ V/m}$$

for $\omega_r/2\pi \approx 6 \text{ GHz}$ ($C \sim 1 \text{ pF}$), $b \approx 5 \mu\text{m}$

$\times 10^6$ larger than E_0
in 3D microwave cavity

Qubit/Photon Coupling



Hamilton operator of qubit (2-level approx.) coupled to resonator:

$$\hat{H} = \frac{\hat{Q}^2}{2C} + \frac{\hat{\phi}^2}{2L} + \frac{E_C}{2}(1 - 2(N_g + \hat{N}_g))\hat{\sigma}_z - \frac{E_J}{2}\hat{\sigma}_x$$

quantum part of gate voltage due to resonator

$$\hat{N}_g = \frac{C_g}{2e}\hat{V}_g = \frac{C_g}{2e}\sqrt{\frac{\hbar\omega_r}{2C}}(\hat{a}^\dagger + \hat{a})$$

Jaynes-Cummings Hamiltonian

Consider bias at charge degeneracy $N_g = 1/2$ and change of qubit basis (z to x , x to $-z$)

$$\hat{H} = \hbar\omega_r(\hat{a}^\dagger\hat{a} + 1/2) + \frac{E_J}{2}\hat{\sigma}_z + \frac{E_C}{2}\frac{C_g}{2e}\sqrt{\frac{\hbar\omega_r}{2C}}(\hat{a}^\dagger + \hat{a})\hat{\sigma}_x$$

Use qubit raising and lowering operators $\hat{\sigma}_x = \hat{\sigma}^+ + \hat{\sigma}^-$

Coupling term in the rotating wave approximation (RWA)

$$\hat{H}_g = \frac{E_C}{2}\frac{C_g}{2e}\sqrt{\frac{\hbar\omega_r}{2C}}(\hat{a}^\dagger\hat{\sigma}^- + \hat{a}\cancel{\hat{\sigma}^-} + \hat{a}^\dagger\cancel{\hat{\sigma}^+} + \hat{a}\hat{\sigma}^+) \approx \hbar g(\hat{a}^\dagger\hat{\sigma}^- + \hat{a}\hat{\sigma}^+)$$

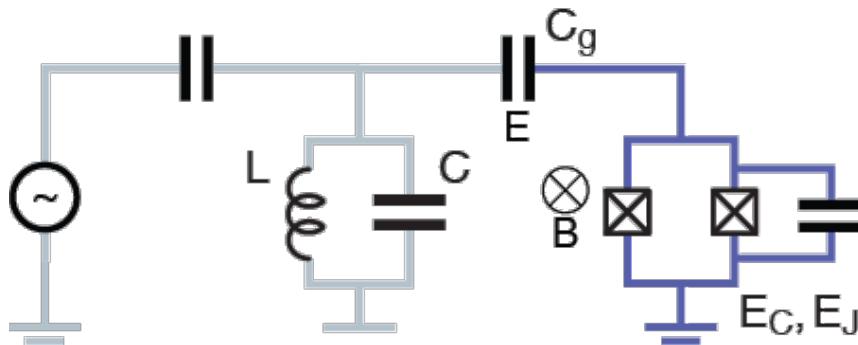
Coupling strength of the Jaynes Cummings Hamiltonian

$$\hbar g = \frac{C_g}{C_\Sigma}2e\sqrt{\frac{\hbar\omega_r}{2C}}$$

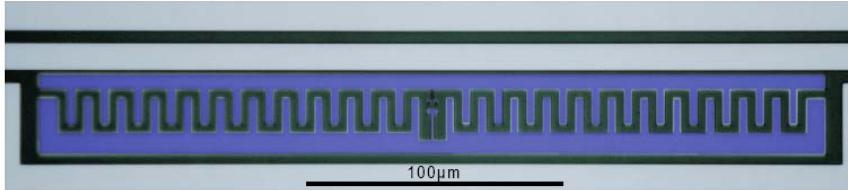
Vacuum-Rabi frequency $\nu_R = \frac{2g}{2\pi} \approx 1 \dots 300 \text{ MHz}$

$g \gg [\kappa, \gamma]$ possible!

Qubit/Photon Coupling in a Circuit



qubit coupled to resonator



coupling strength:

$$\hbar g = eV_{0,\text{rms}} \frac{C_g}{C_\Sigma}$$

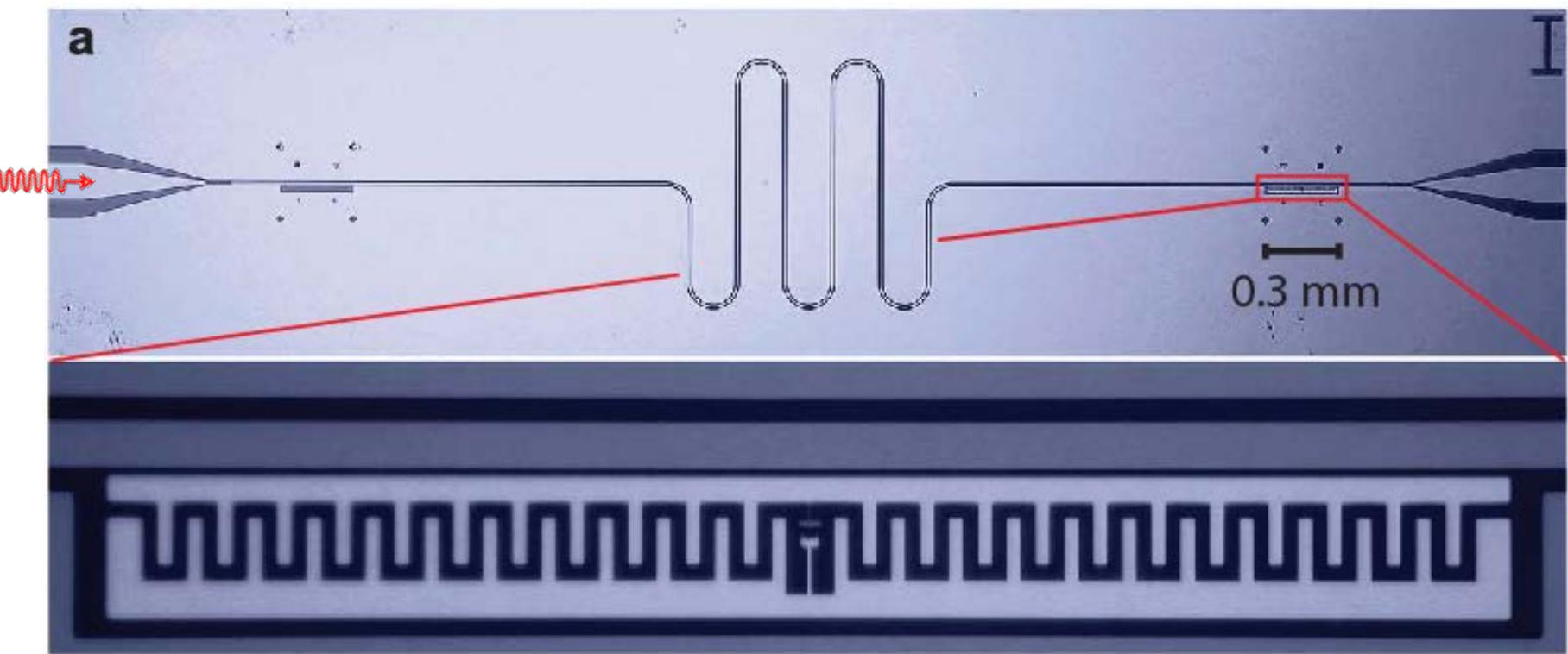
$$\Rightarrow \nu_{\text{vac}} = \frac{g}{\pi} \approx 1 \dots 300 \text{ MHz}$$

$g \gg [\kappa, \gamma]$ possible!

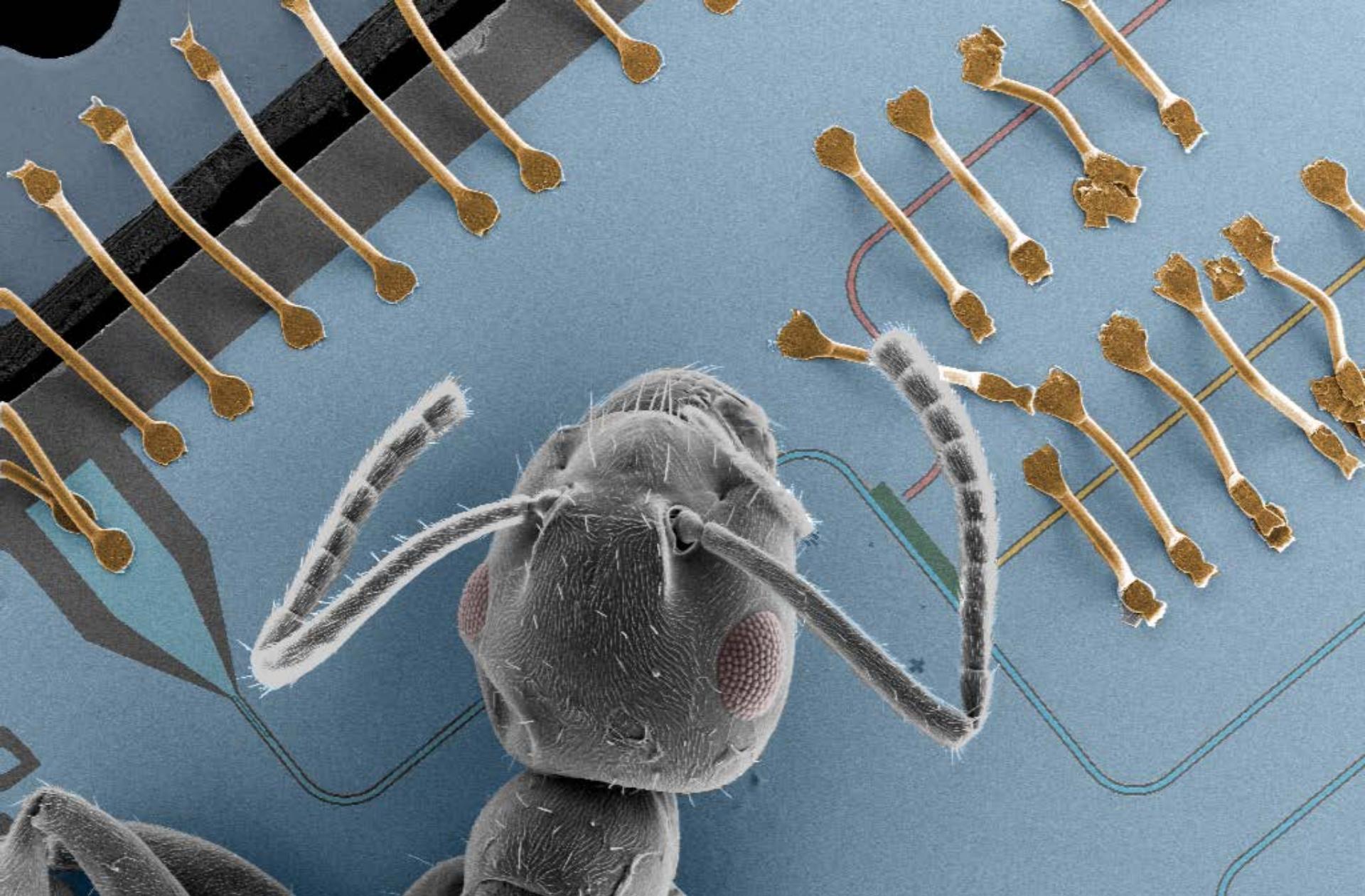
large effective dipole moment

$$d = \frac{\hbar g}{E_0} \sim 10^2 \dots 10^4 ea_0$$

Circuit QED with One Photon



superconducting cavity QED circuit



ETH

Eidgenössische Technische Hochschule Zürich
Swiss Federal Institute of Technology Zurich

J. Mlynek *et al.*, Quantum Device Lab, ETH Zurich (2012)

Sample Mount

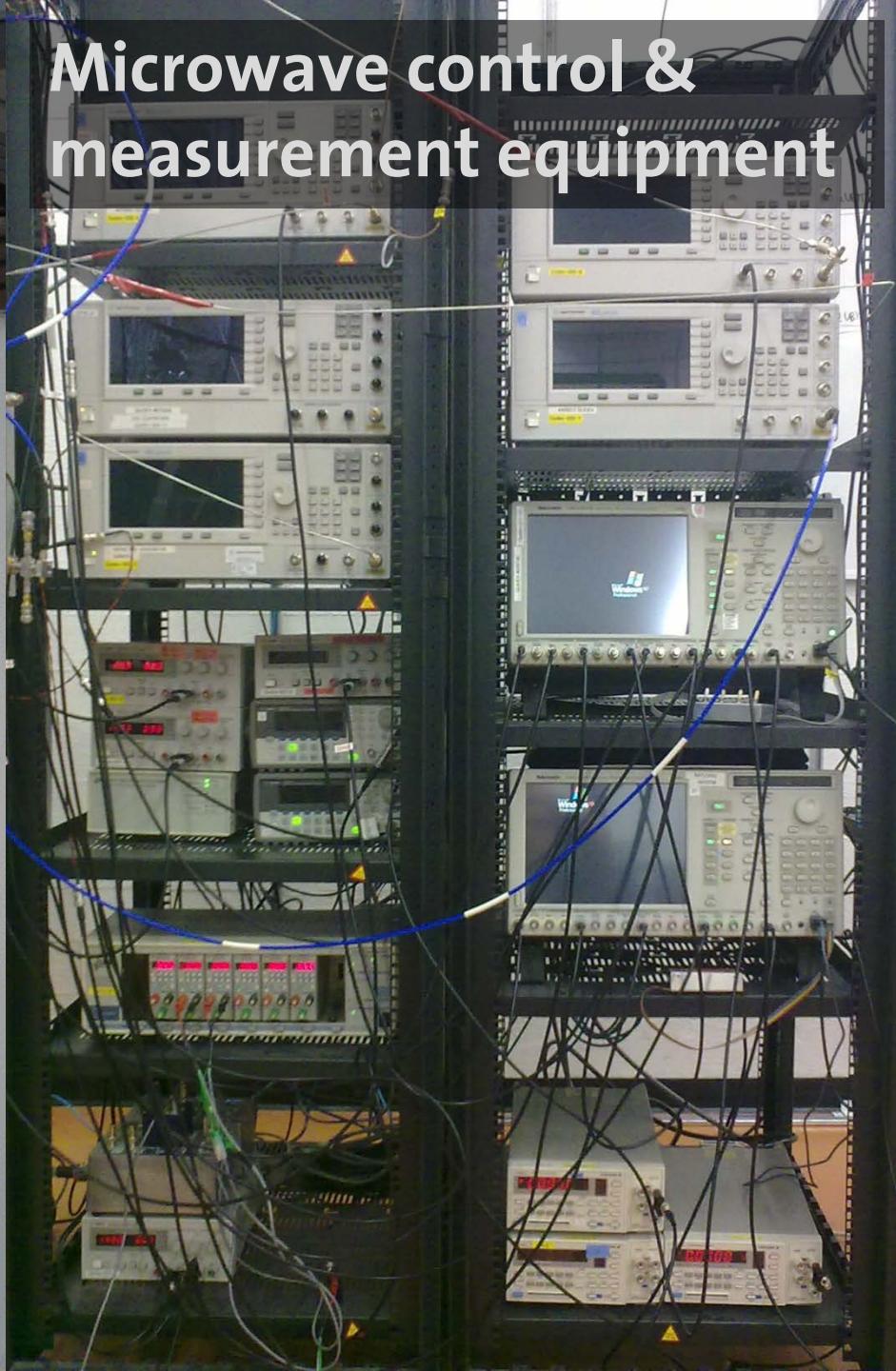


~ 2 cm

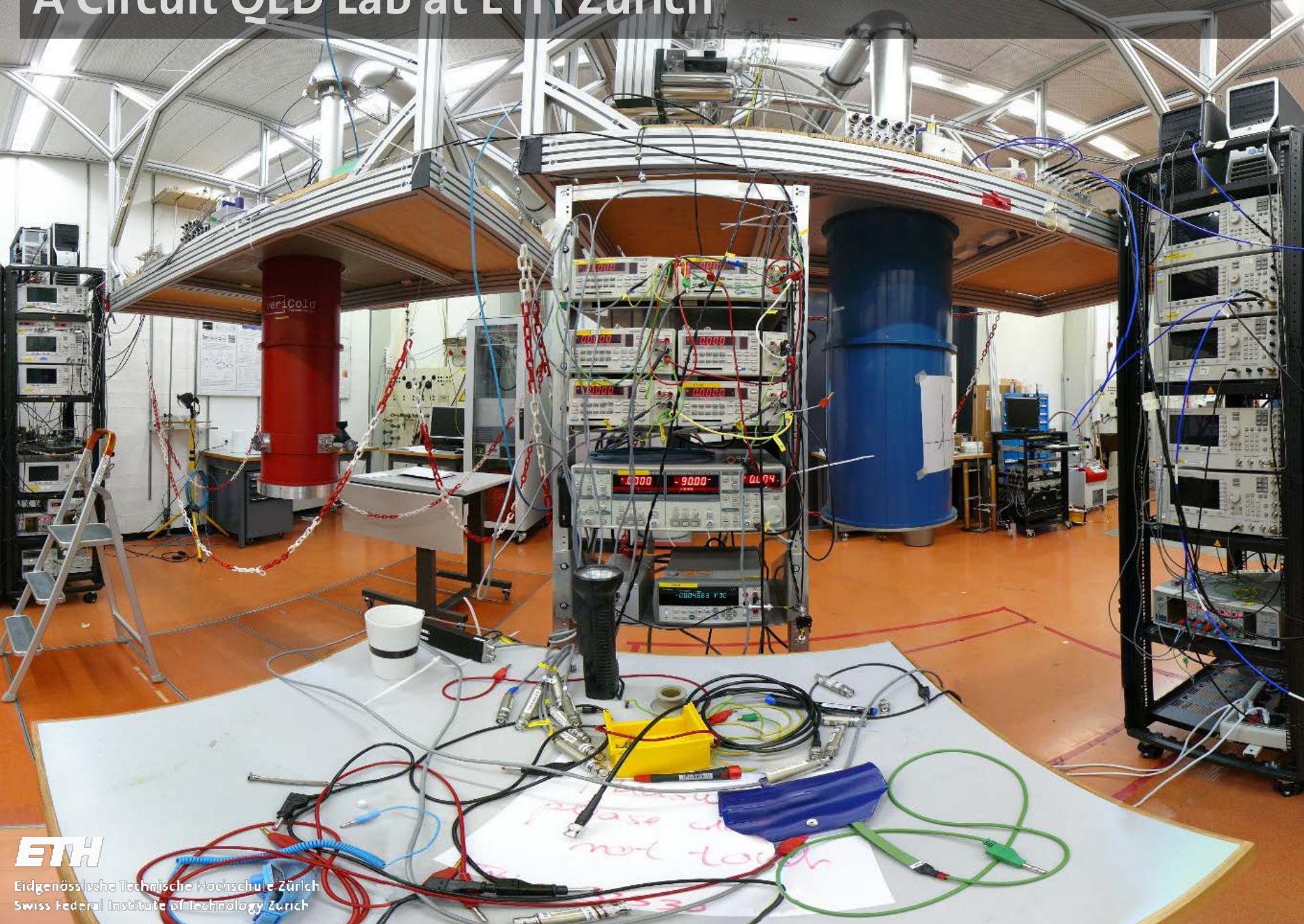
Cryostate for temperatures down to 0.02 K



Microwave control & measurement equipment

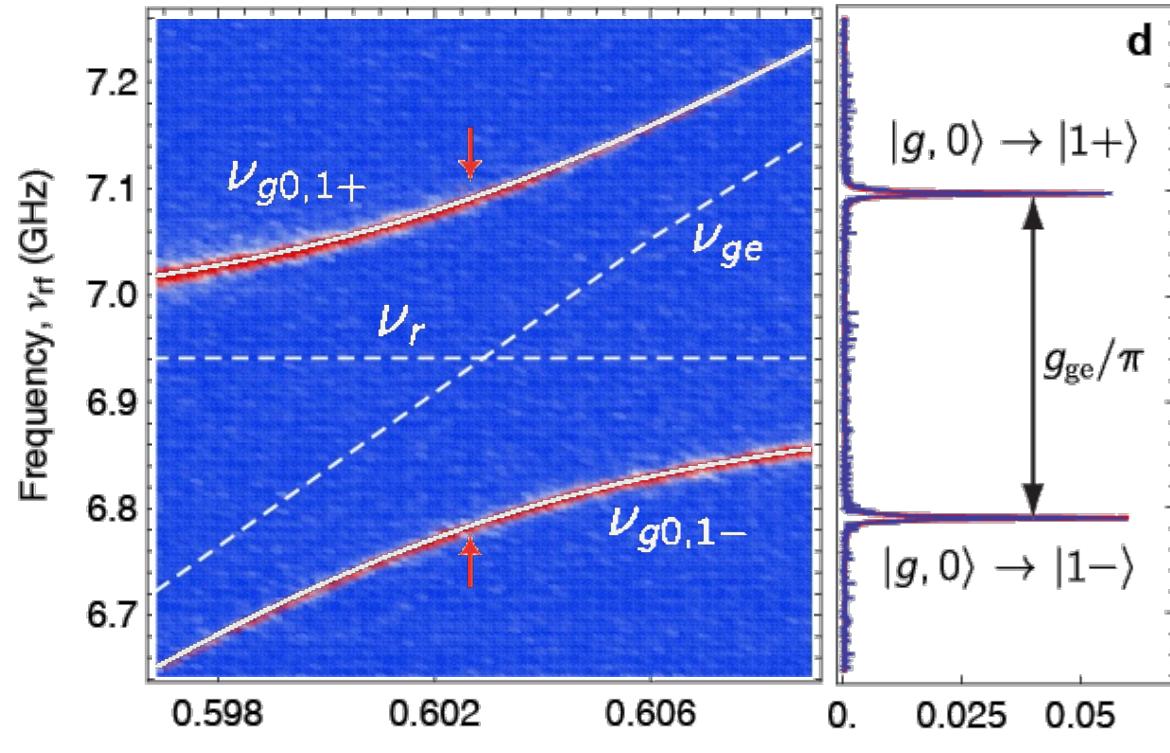


A Circuit QED Lab at ETH Zurich



Resonant Vacuum Rabi Mode Splitting ...

... with one photon ($n=1$):

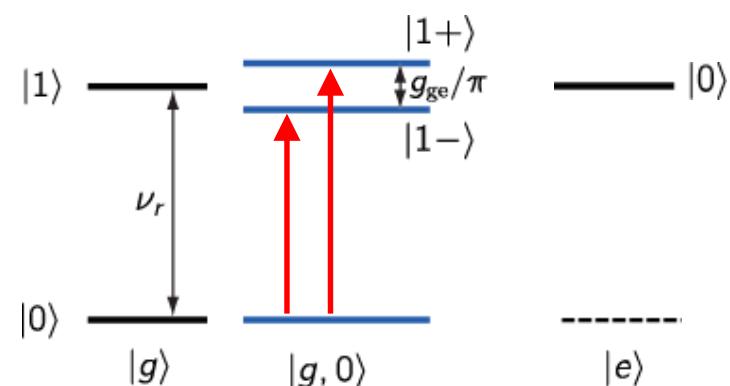


very strong coupling:

$$g_{ge}/\pi = 308 \text{ MHz}$$

$$\kappa, \gamma < 1 \text{ MHz}$$

$$g_{ge} \gg \kappa, \gamma$$



forming a 'molecule' of a qubit and a photon

first demonstration in a solid: A. Wallraff et al., *Nature (London)* 431, 162 (2004)

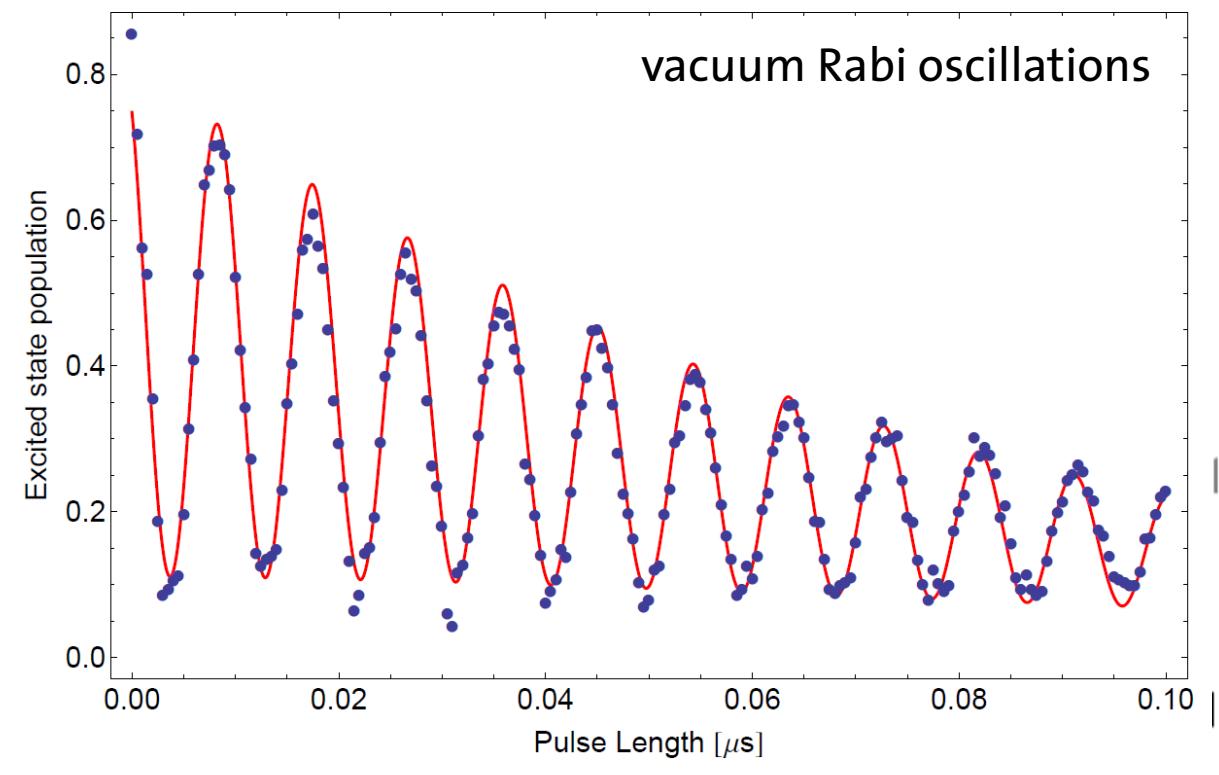
this data: J. Fink et al., *Nature (London)* 454, 315 (2008)

R. J. Schoelkopf, S. M. Girvin, *Nature (London)* 451, 664 (2008)

Resonant Vacuum Rabi Mode Splitting ...

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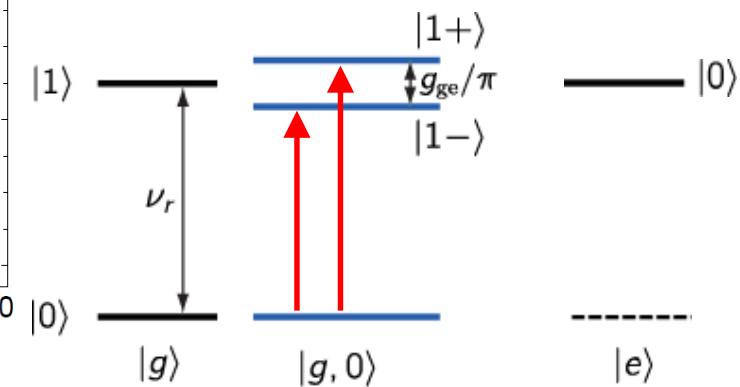
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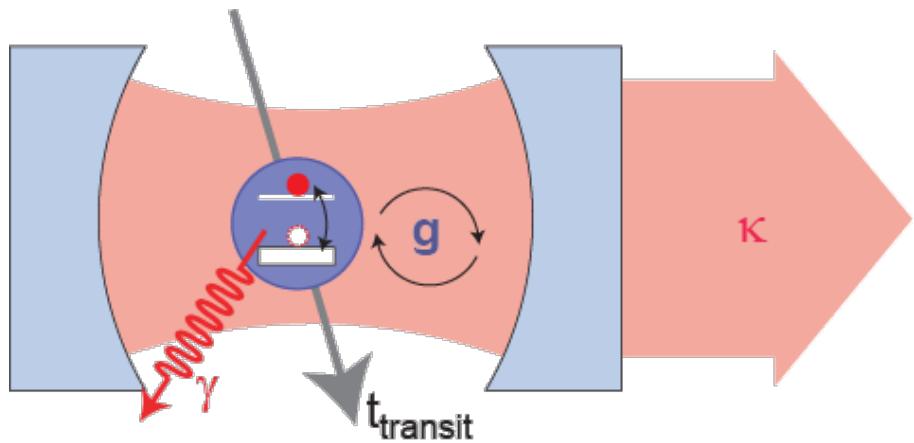
forming a 'molecule' of a qubit and a photon

first demonstration in a solid: A. Wallraff *et al.*, *Nature (London)* 431, 162 (2004)

this data: J. Fink *et al.*, *Nature (London)* 454, 315 (2008)

R. J. Schoelkopf, S. M. Girvin, *Nature (London)* 451, 664 (2008)

Cavity QED with Superconducting Circuits



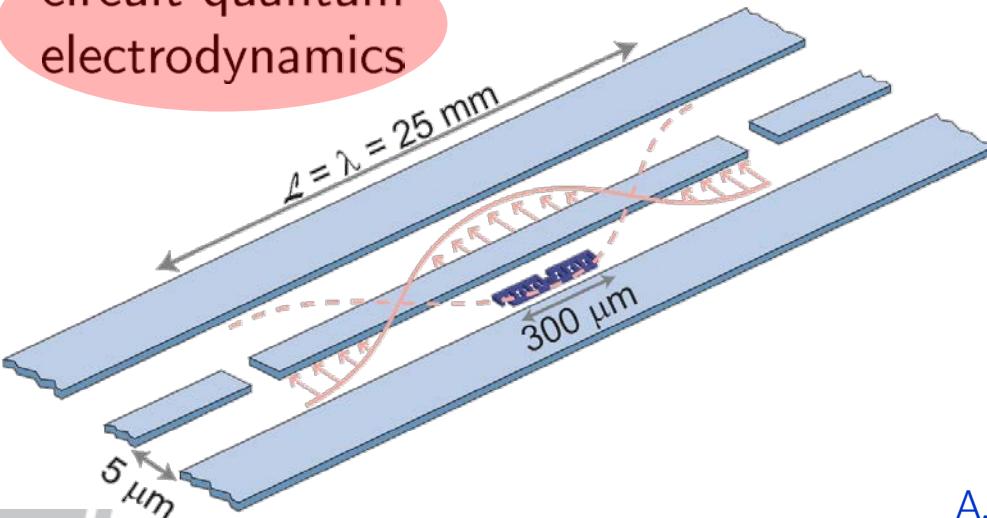
coherent interaction of photons with quantum two-level systems ...

J. M. Raimond *et al.*, *Rev. Mod. Phys.* 73, 565 (2001)

S. Haroche & J. Raimond, *OUP Oxford* (2006)

J. Ye., H. J. Kimble, H. Katori, *Science* 320, 1734 (2008)

circuit quantum
electrodynamics



Properties:

- strong coupling in solid state sys.
- ‘easy’ to fabricate and integrate

Research directions:

- quantum optics
- hybrid quantum systems
- quantum information

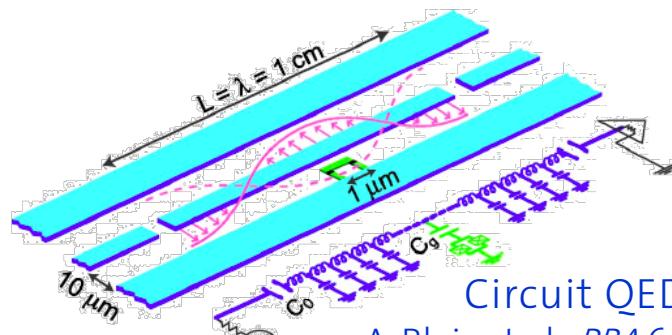
A. Blais, *et al.*, *PRA* 69, 062320 (2004)

A. Wallraff *et al.*, *Nature (London)* 431, 162 (2004)

R. J. Schoelkopf, S. M. Girvin, *Nature (London)* 451, 664 (2008)

Research Directions & Applications

Quantum Computing with Superconducting Circuits

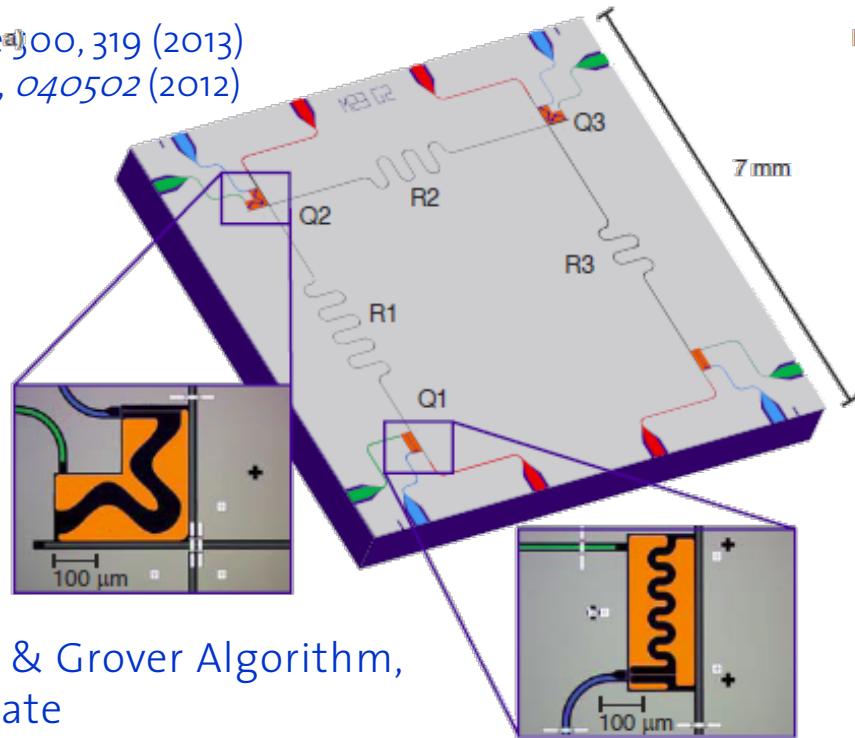


Teleportation

L. Steffen *et al.*, *Nature* 500, 319 (2013)

M. Baur *et al.*, *PRL* 108, 040502 (2012)

- Circuit QED Architecture
- A. Blais *et al.*, *PRA* 69, 062320 (2004)
 - A. Wallraff *et al.*, *Nature* 431, 162 (2004)
 - M. Sillanpaa *et al.*, *Nature* 449, 438 (2007)
 - H. Majer *et al.*, *Nature* 449, 443 (2007)
 - M. Mariantoni *et al.*, *Science* 334, 61 (2011)
 - R. Barends *et al.*, *Nature* 508, 500 (2014)

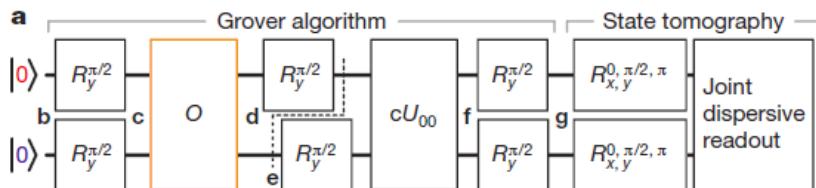


Deutsch & Grover Algorithm, Toffoli Gate

L. DiCarlo *et al.*, *Nature* 460, 240 (2009)

L. DiCarlo *et al.*, *Nature* 467, 574 (2010)

M. Reed *et al.*, *Nature* 481, 382 (2012)



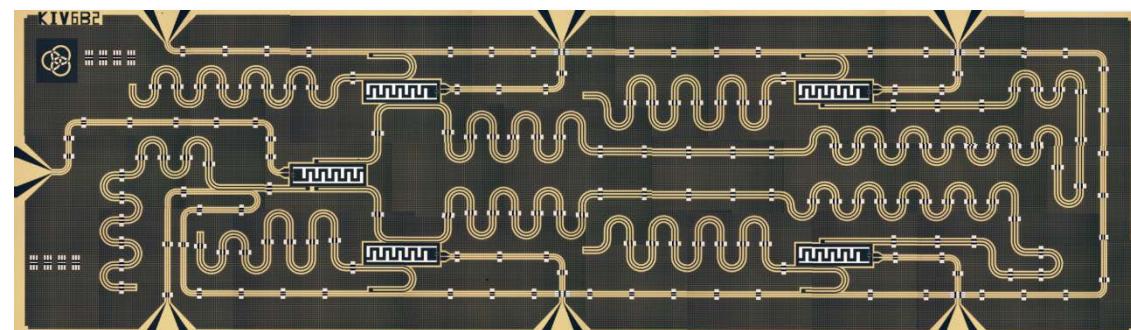
Error Correction

M. Reed *et al.*, *Nature* 481, 382 (2012)

Corcoles *et al.*, *Nat. Com.* 6, 6979 (2015)

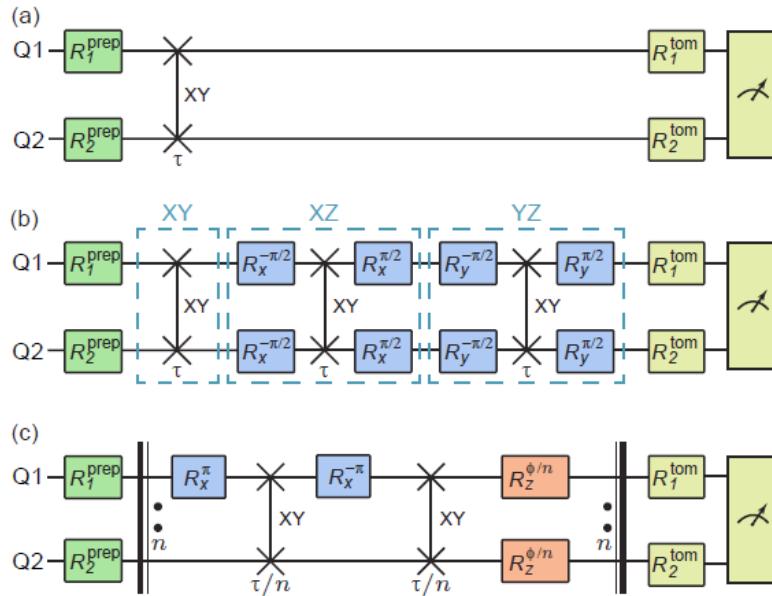
Ristè *et al.*, *Nat. Com.* 6, 6983 (2015)

Kelly *et al.*, *Nature* 519, 66-69 (2015)



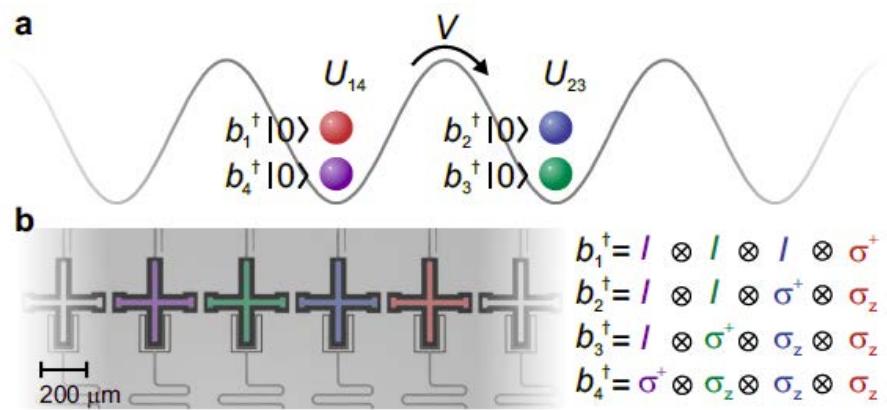
Quantum Simulation with Superconducting Circuits

Digital simulation of exchange,
Heisenberg, Ising spin models

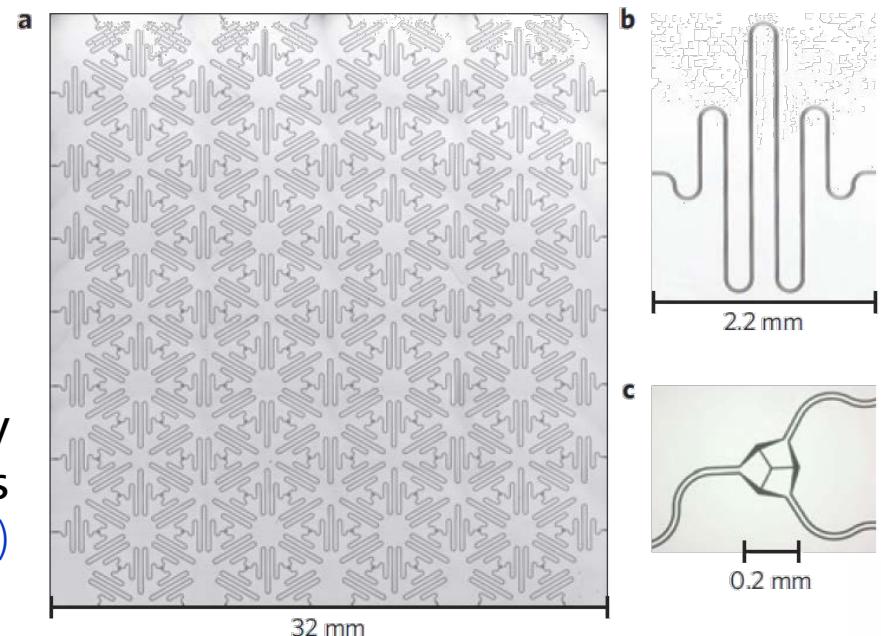


Salathe *et al.*, PRX 5, 021027 (2015)
[arXiv:1502.06778](https://arxiv.org/abs/1502.06778)

... two-mode fermionic Hubbard models

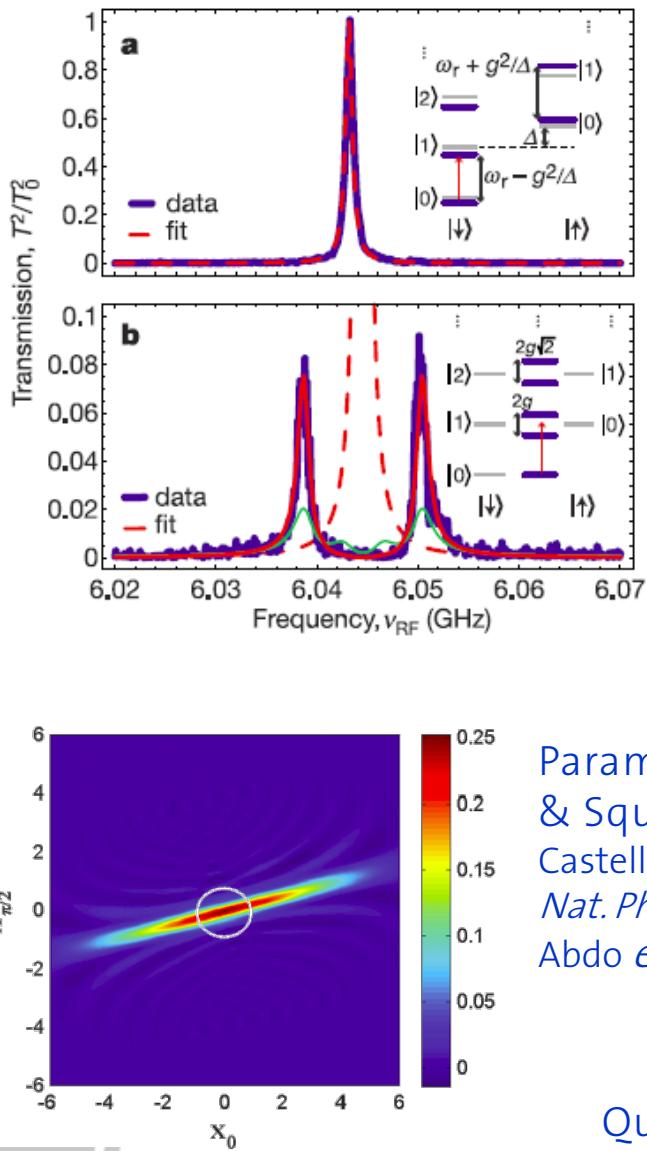


Barends *et al.*, arXiv:1501.07703, (2015)



Analog simulations with cavity
and/or qubit arrays
Houck *et al.*, Nat Phys. 8, 292 (2012)

Quantum Optics with Supercond. Circuits

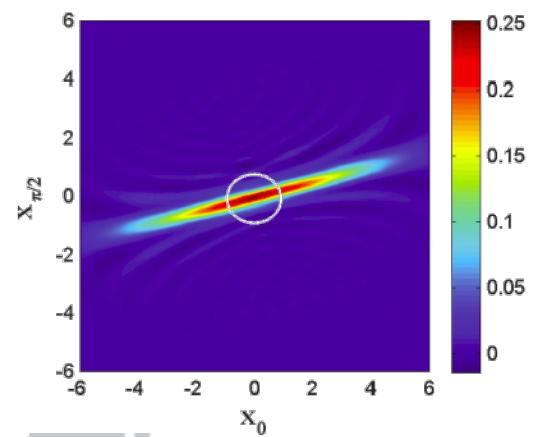


Strong Coherent Coupling
Chiorescu *et al.*, *Nature* 431, 159 (2004)
Wallraff *et al.*, *Nature* 431, 162 (2004)
Schuster *et al.*, *Nature* 445, 515 (2007)

Root n Nonlinearities
Fink *et al.*, *Nature* 454, 315 (2008)
Deppe *et al.*, *Nat. Phys.* 4, 686 (2008)
Bishop *et al.*, *Nat. Phys.* 5, 105 (2009)

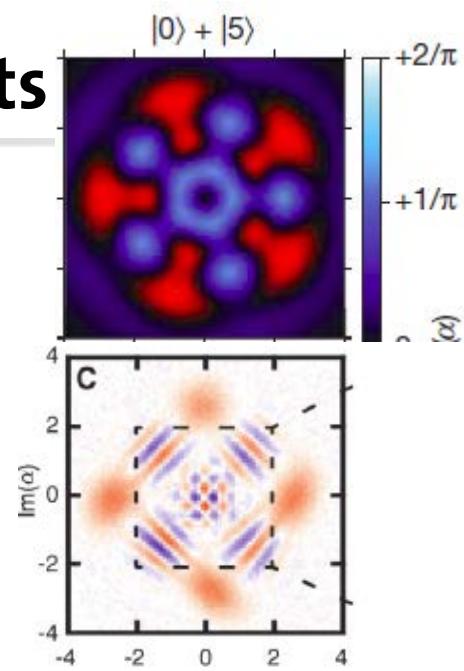
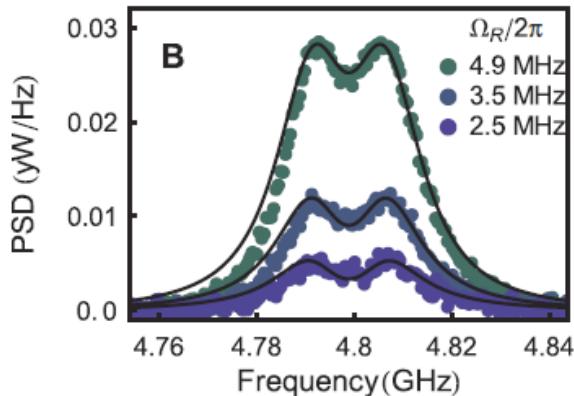


Microwave Fock and Cat States
Hofheinz *et al.*, *Nature* 454, 310 (2008)
Hofheinz *et al.*, *Nature* 459, 546 (2009)
Kirchmair *et al.*, *Nature* 495, 205 (2013)
Vlastakis *et al.*, *Science* 342, 607 (2013)



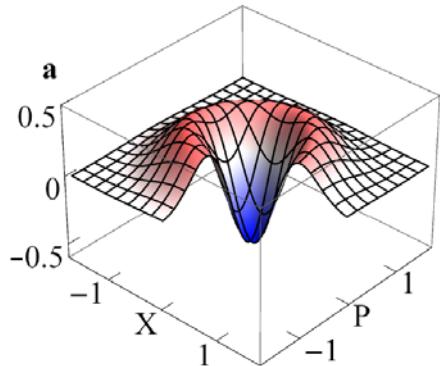
Parametric Amplification & Squeezing
Castellanos-Beltran *et al.*,
Nat. Phys. 4, 928 (2008)
Abdo *et al.*, *PRX* 3, 031001 (2013)

Waveguide QED –
Qubit Interactions in Free Space
Astafiev *et al.*, *Science* 327, 840 (2010)
van Loo *et al.*, *Science* 342, 1494 (2013)



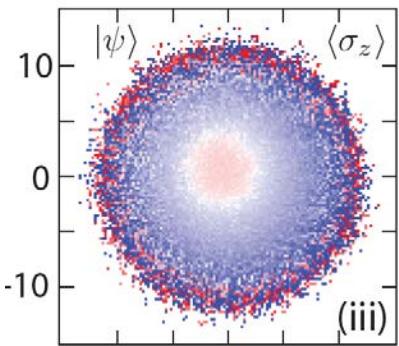
Experiments with Propagating Microwaves in 1D

Full state tomography and Wigner functions of propagating photons



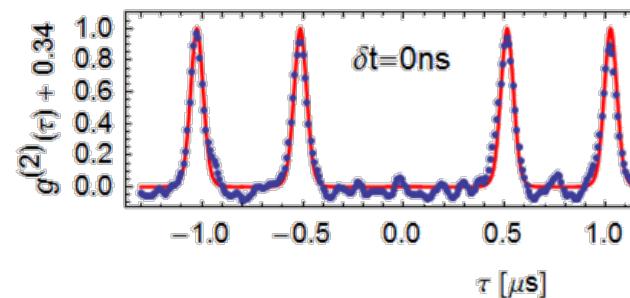
Eichler *et al.*, PRL 106, 220503 (2011)

Preparation and characterization of qubit-propagating photon entanglement

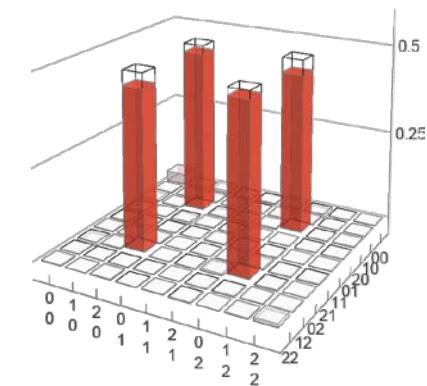


Eichler *et al.*, PRL 109, 240501 (2012)
Eichler *et al.*, PRA 86, 032106 (2012)

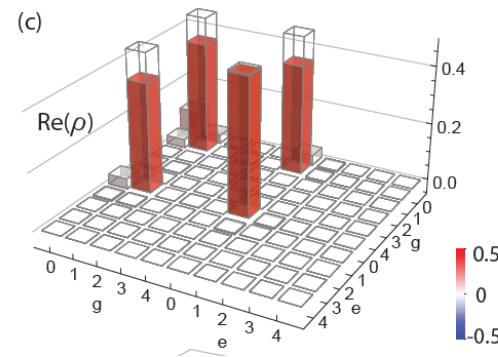
Hong-Ou-Mandel: Two-photon interference incl. msrmnt of coherences at microwave freq.



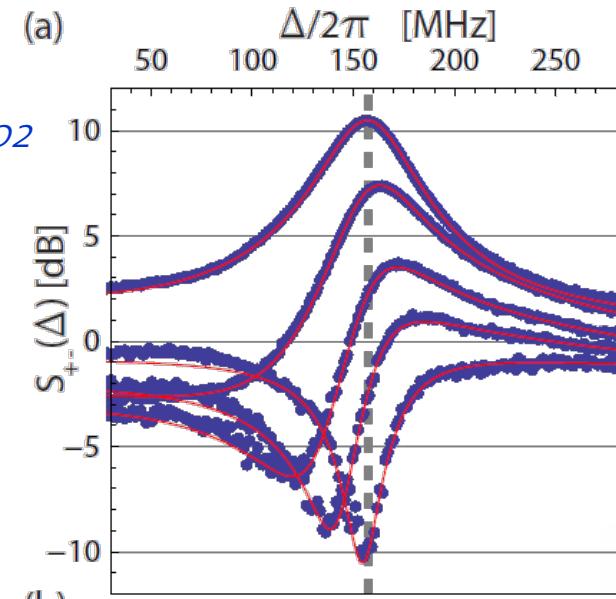
Lang *et al.*, Nat. Phys. 9, 345 (2013)



Squeezing in a Josephson parametric dimer



Eichler *et al.*,
PRL 113, 110502
(2014)

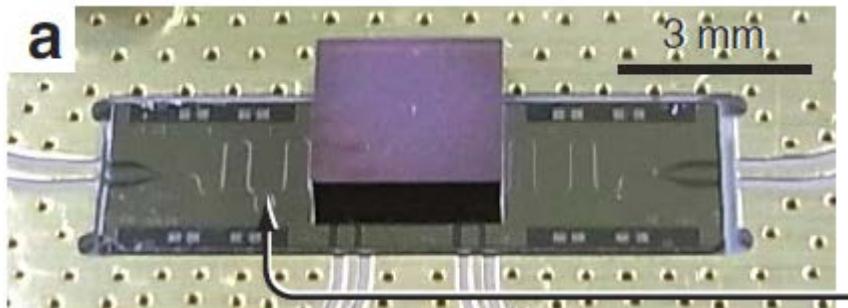


Hybrid Systems with Superconducting Circuits

Spin Ensembles: e.g. NV centers

D. Schuster *et al.*, PRL 105, 140501 (2010)

Y. Kubo *et al.*, PRL 105, 140502 (2010)

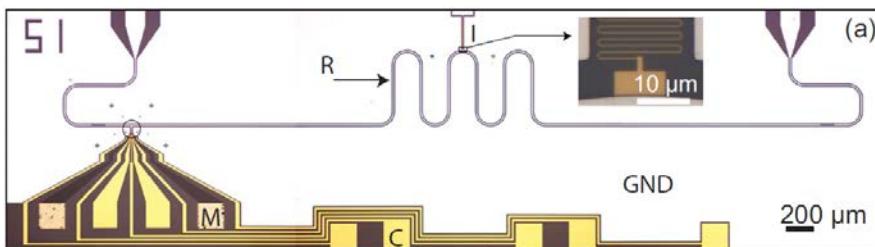


CNT, Gate Defined 2DEG, or nanowire Quantum Dots

M. Delbecq *et al.*, PRL 107, 256804 (2011)

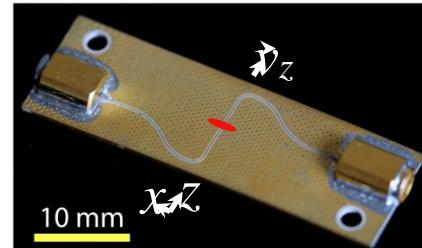
T. Frey *et al.*, PRL 108, 046807 (2012)

K. Petersson *et al.*, Nature 490, 380 (2013)



Rydberg Atoms

S. Hogan *et al.*, PRL 108, 063004 (2012)



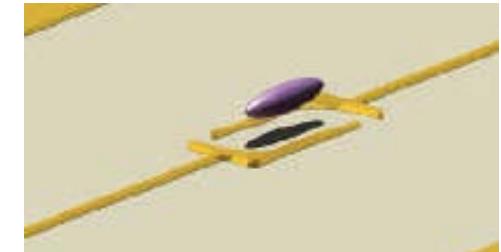
Polar Molecules, Rydberg, BEC

P. Rabl *et al.*, PRL 97, 033003 (2006)

A. Andre *et al.*, Nat. Phys. 2, 636 (2006)

D. Petrosyan *et al.*, PRL 100, 170501 (2008)

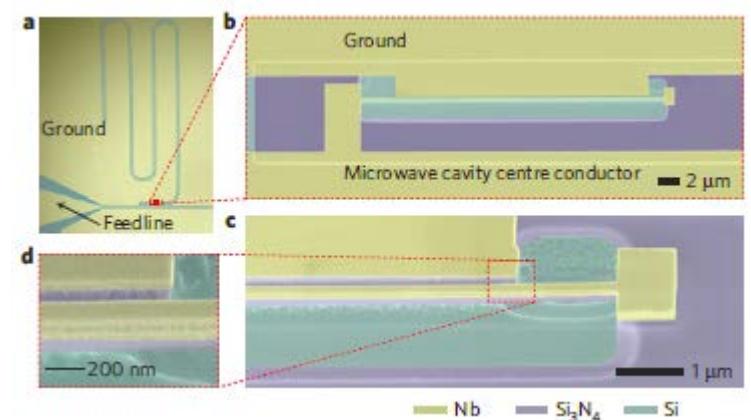
J. Verdu *et al.*, PRL 103, 043603 (2009)



Nano-Mechanics

J. Teufel *et al.*, Nature 475, 359 (2011)

X. Zhou *et al.*, Nat. Phys. 9, 179 (2013)



... and many more

Circuit QED Research

Analog and
Digital Electronics

Microwaves

Cryogenics

Quantum
Information

Quantum
Optics

Quantum Physics
in the Solid State

Hybrid
Systems

Micro- and
Nano-Fabrication

Measurement Technology

Control and
Acquisition

The ETH Zurich Quantum Device Lab

incl. undergrad and summer students



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Swiss Federal Institute of Technology Zurich

QSIT Quantum
Science and
Technology
National Centre of Competence in Research

(-)QED
CIRCUIT AND CAVITY
QUANTUM ELECTRODYNAMICS

FNSNF

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