

# Grover algorithm in superconducting circuits

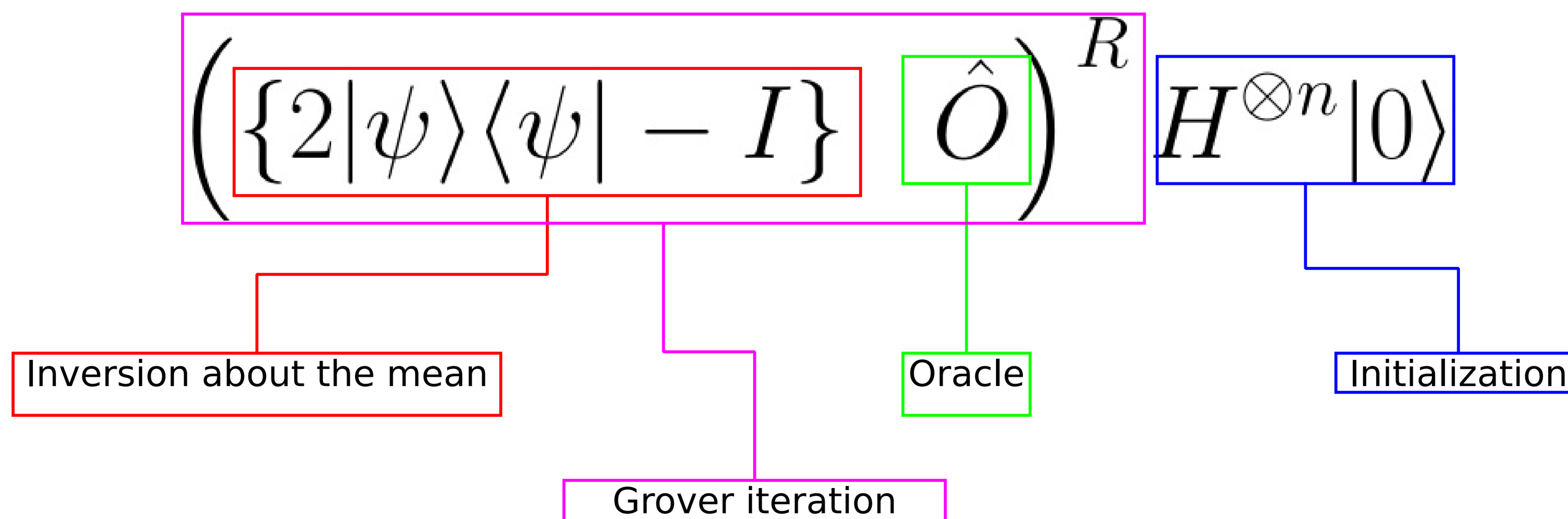
Folkert de Vries & Sergiό Solόrzano



# Grover's search algorithm

Given a name find the phone number ->  
structured search (easy)

Given a phone number, find the matching name ->  
unstructured search (hard)



## The oracle

Don't know the solution but is able to recognize it.

By definition:

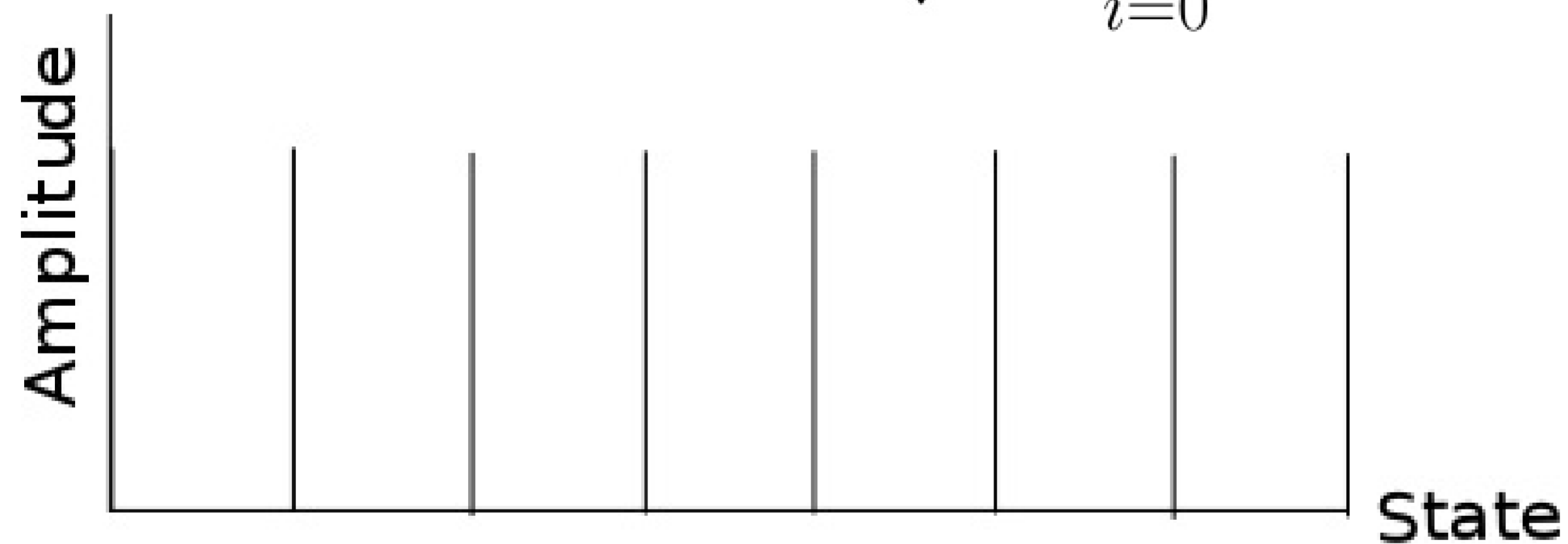
$$\hat{O}|x\rangle = (-1)^{f(x)}|x\rangle$$

$$f(x) = \begin{cases} 1 & x \text{ Is a solution.} \\ 0 & x \text{ Is not a solution.} \end{cases}$$

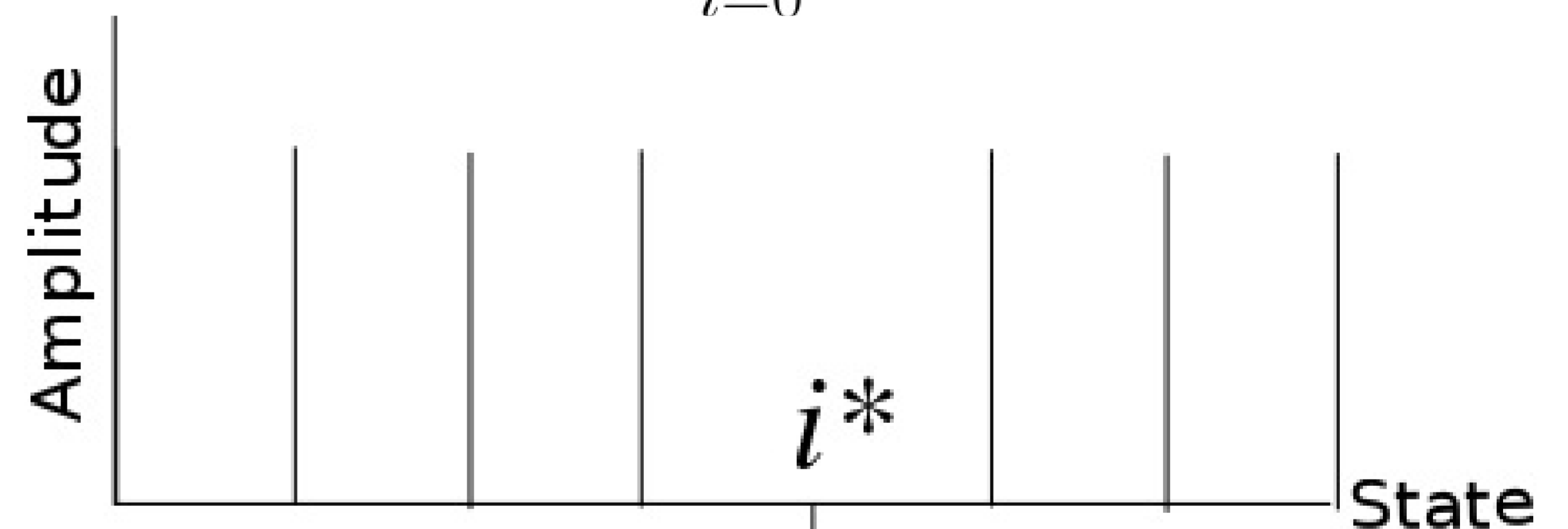
# Initialization

Sets a state that uses the quantum mechanical "built in" parallelization

$$H^{\otimes n}|0\rangle = |\psi\rangle = \frac{1}{\sqrt{N}} \sum_{i=0}^N |i\rangle$$



$$\hat{O}|\psi\rangle = \frac{1}{\sqrt{N}} \sum_{i=0}^N (-1)^{f(i)} |i\rangle$$



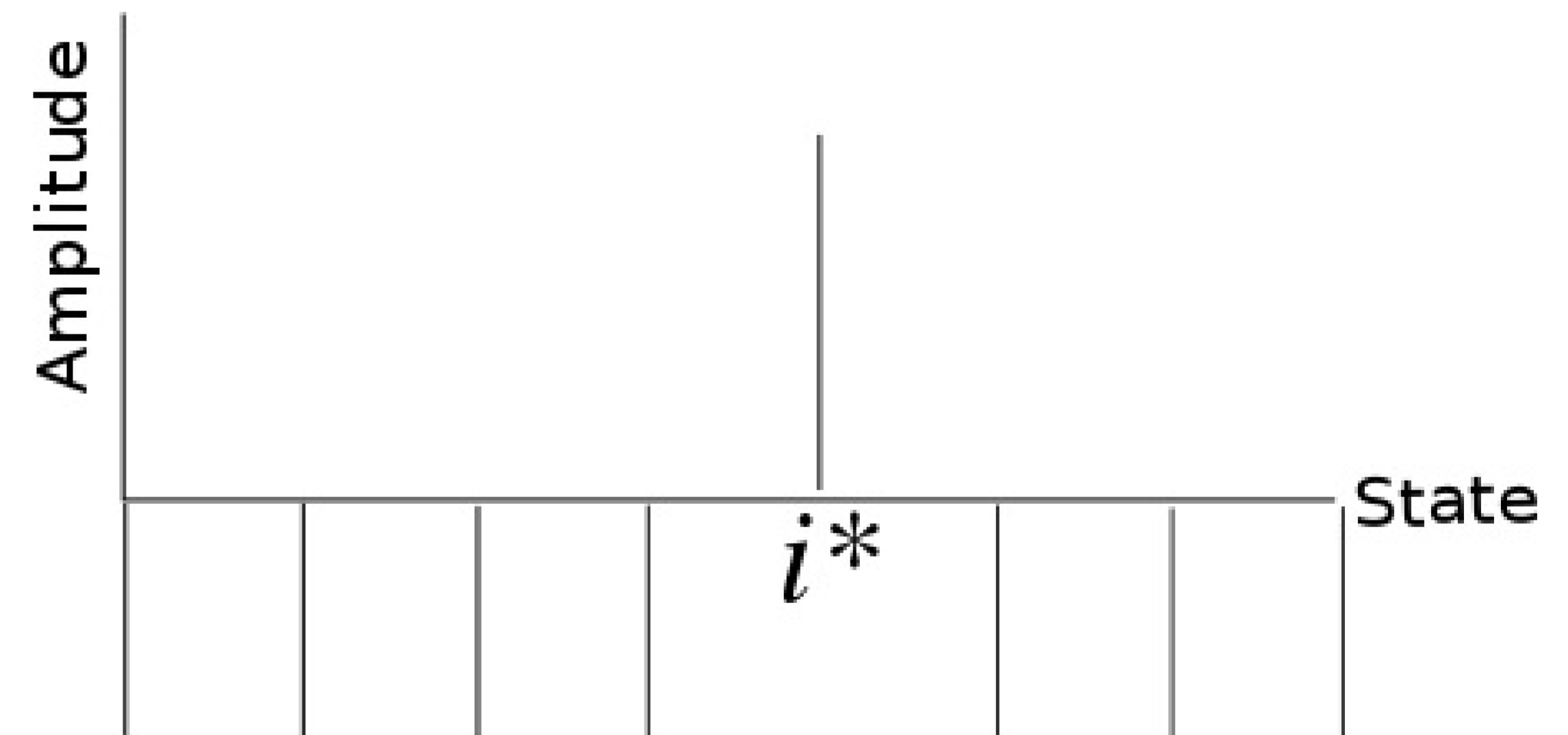
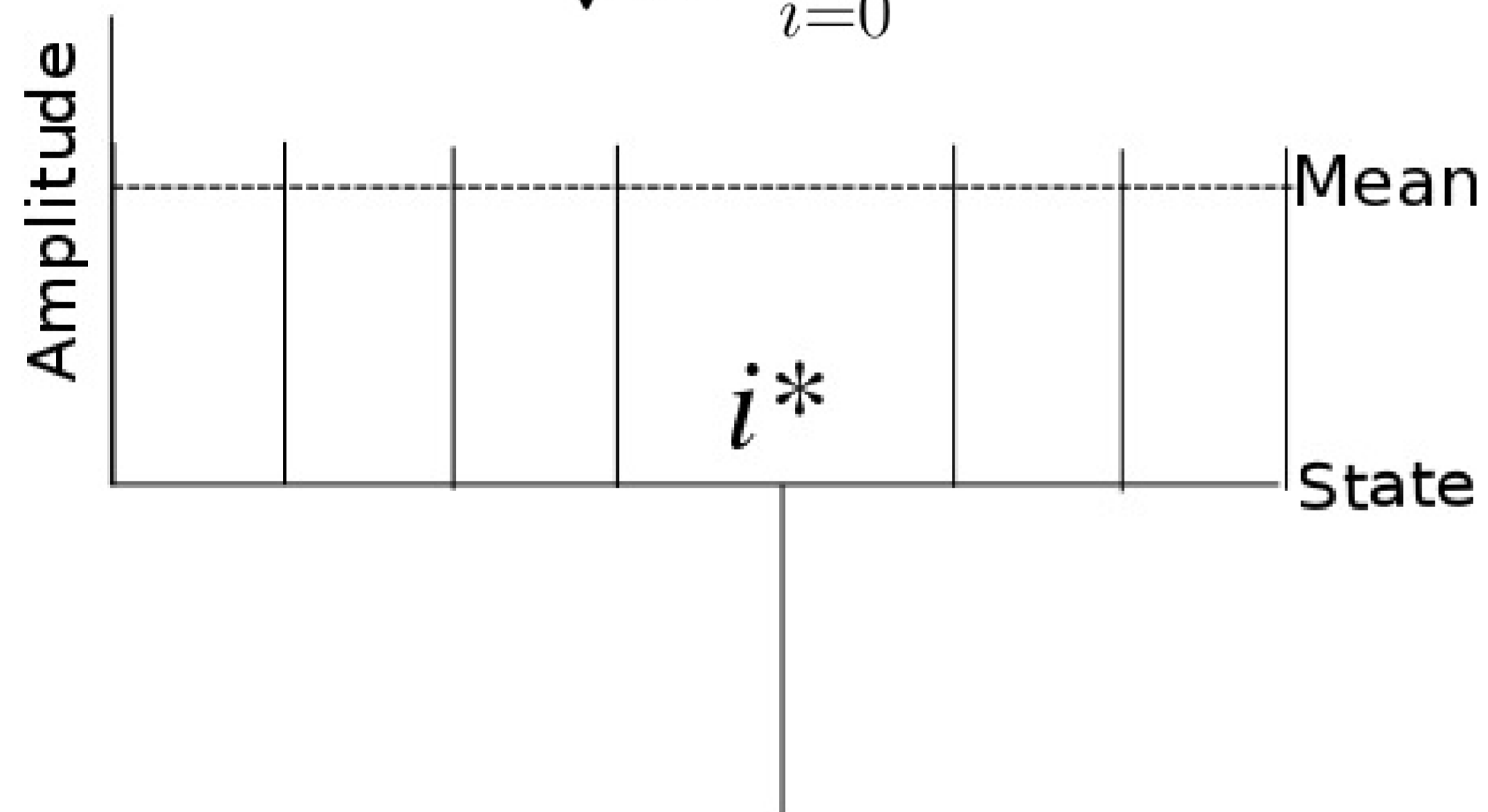
Not useful yet, all states are equally probable.

# Inversion about the mean

$$\{2|\psi\rangle\langle\psi|-I\} \sum_{i=0}^N \alpha_i |i\rangle = \sum_{i=0}^N (-\alpha_i + 2\langle\alpha\rangle) |i\rangle$$

$$\hat{O}|\psi\rangle = \frac{1}{\sqrt{N}} \sum_{i=0}^N (-1)^{f(i)} |i\rangle$$

$$\{2|\psi\rangle\langle\psi|-I\} \hat{O}|\psi\rangle$$

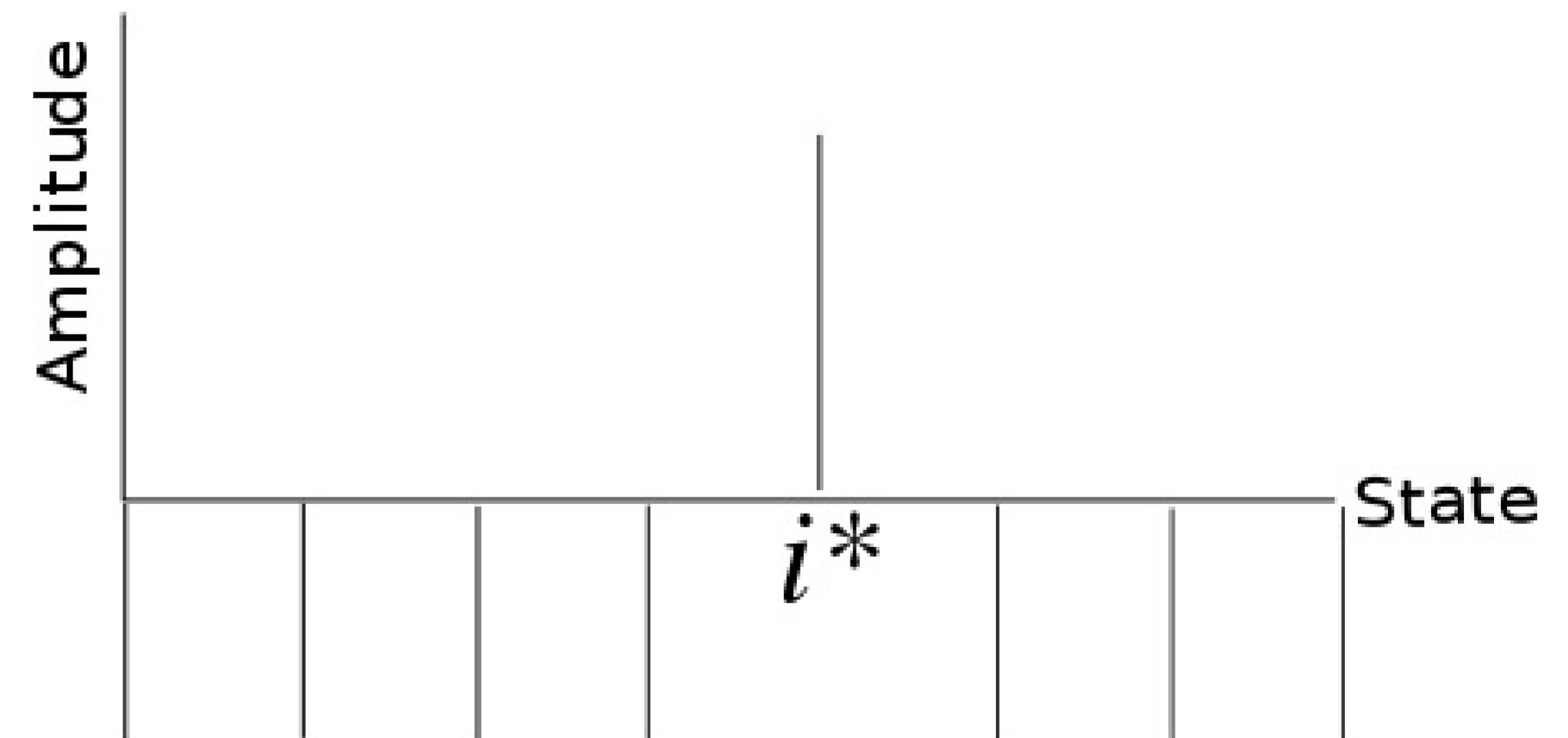
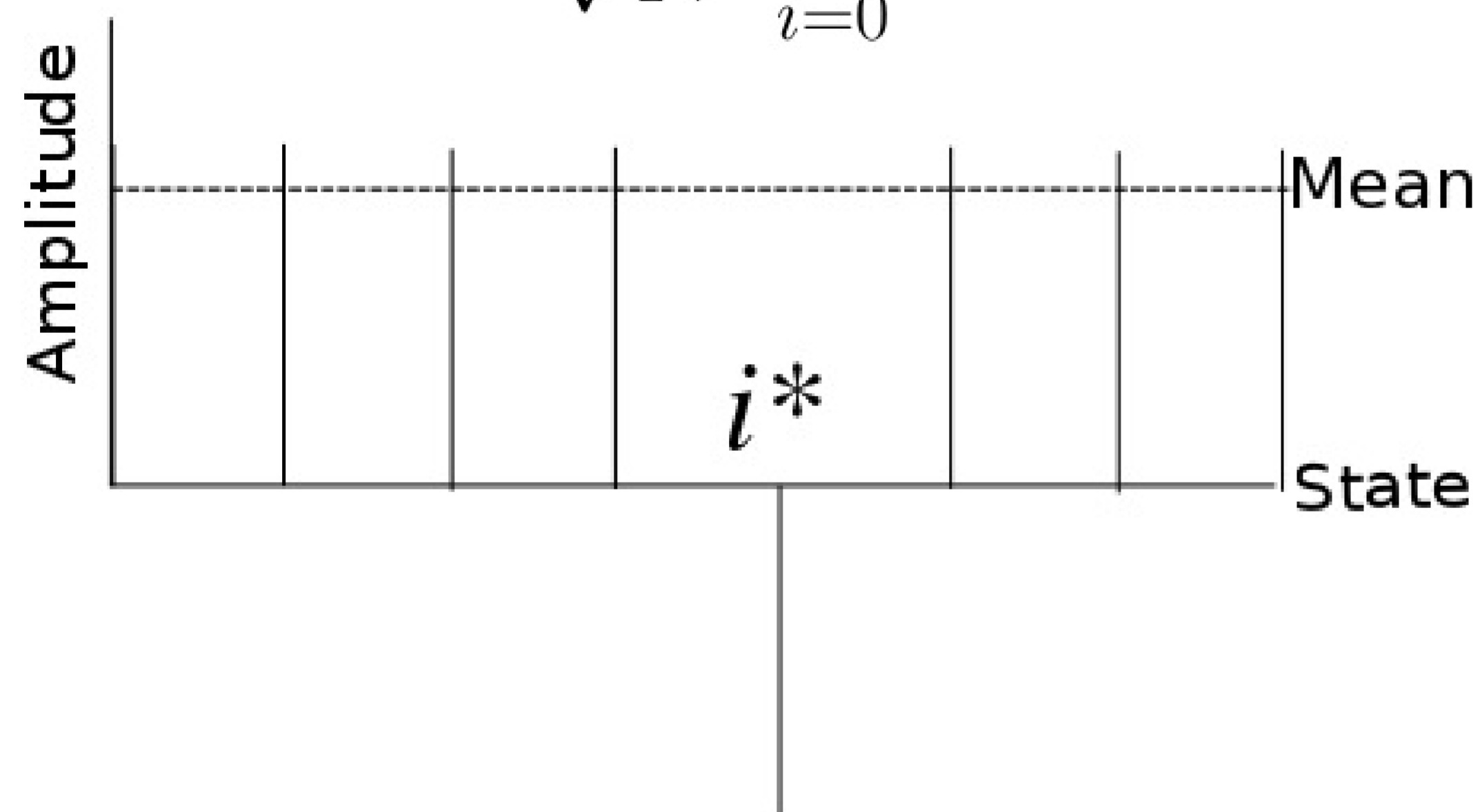


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$$\{2|\psi\rangle\langle\psi|-I\} \hat{O}|\psi\rangle$$



The amplitude of the target state grows while all other amplitudes shrink.

# Inversion about the mean

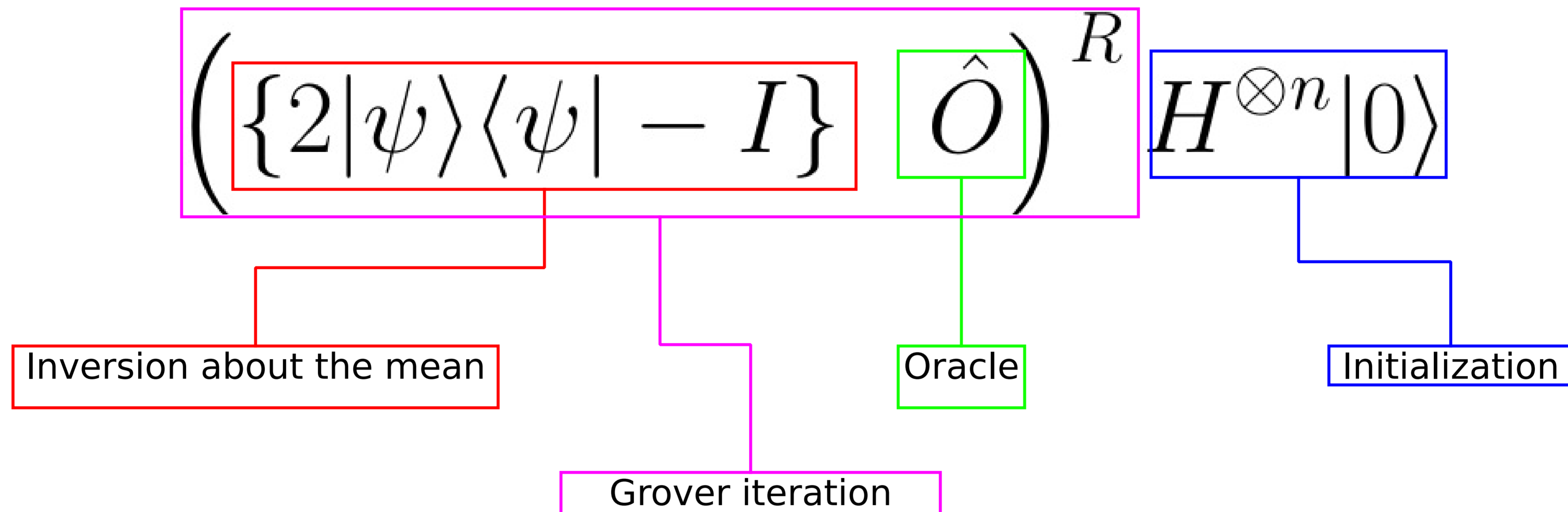
$$2|\psi\rangle\langle\psi| - I = H^{\otimes n} (2|0\rangle\langle 0| - I) H^{\otimes n}$$

$$= H^{\otimes n} C_{00} H^{\otimes n}$$

$$C_{00}|x\rangle = -(-1)^{\delta_{x,0}}|x\rangle$$

The inversion about the mean operation can be decomposed into **one-qubit** and **two-qubit** operations

# Grover's search algorithm

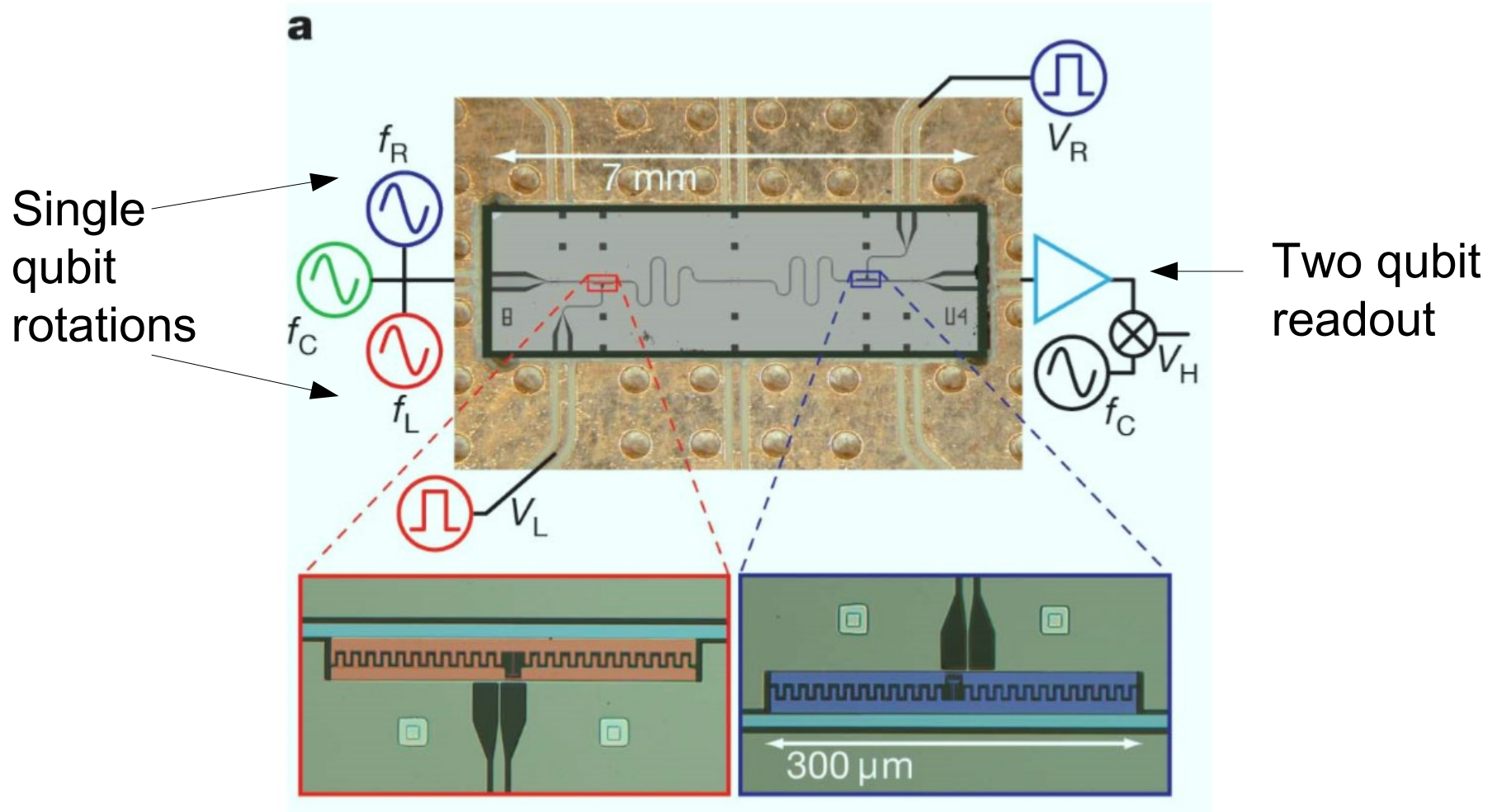


$$R \approx O(\sqrt{N}) \quad \text{Steps}$$

$$O(\log(N)) \quad \text{Qubits}$$



# Experiment 1



DiCarlo et al. Nature 460, 240 (2009)

# Jaynes-Cummings Hamiltonian

$$H = \hbar\omega_r \left( a^\dagger a + \frac{1}{2} \right) - \sum_{i=1,2} \frac{\hbar\omega_i}{2} \sigma_{i,z} + \hbar \sum_{i=1,2} g_i (a^\dagger \sigma_i^- + a \sigma_i^+).$$

↓ Dispersive regime

$$H = \hbar (\omega_r + \chi_1 \sigma_{1,z} + \chi_2 \sigma_{2,z}) a^\dagger a + \sum_{i=1,2} \hbar \frac{\chi_i - \omega_i}{2} \sigma_{i,z} + \hbar J (\sigma_1^+ \sigma_2^- + \sigma_1^- \sigma_2^+).$$

# Jaynes-Cummings Hamiltonian

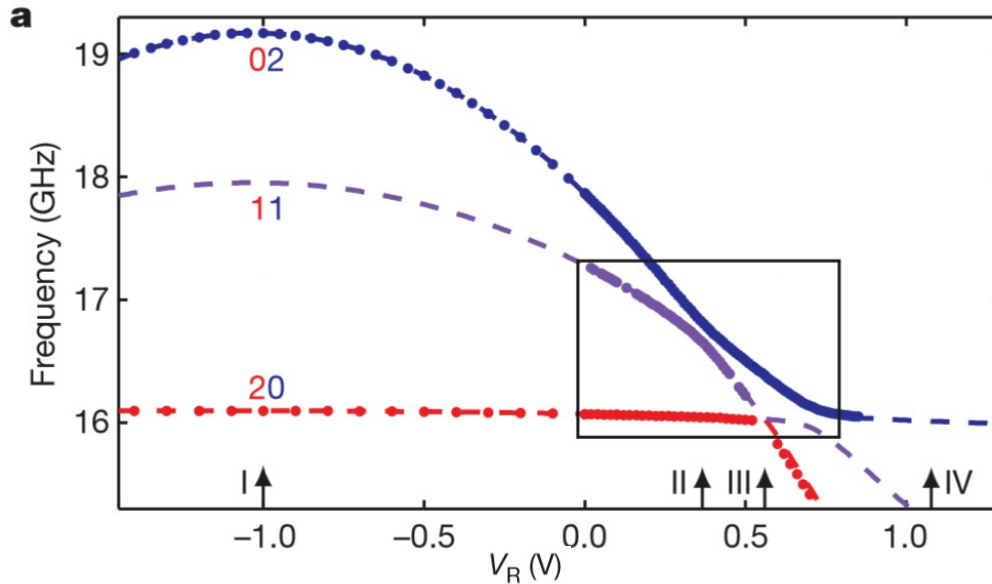
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$$J = \frac{1}{2} g_1 g_2 \left( \frac{1}{\omega_1 - \omega_r} + \frac{1}{\omega_2 - \omega_r} \right)$$

# Oracle: Controlled-phase gate



$$|00\rangle \rightarrow |00\rangle$$

$$|01\rangle \rightarrow |01\rangle$$

$$|10\rangle \rightarrow |10\rangle$$

$$|11\rangle \rightarrow -|11\rangle$$

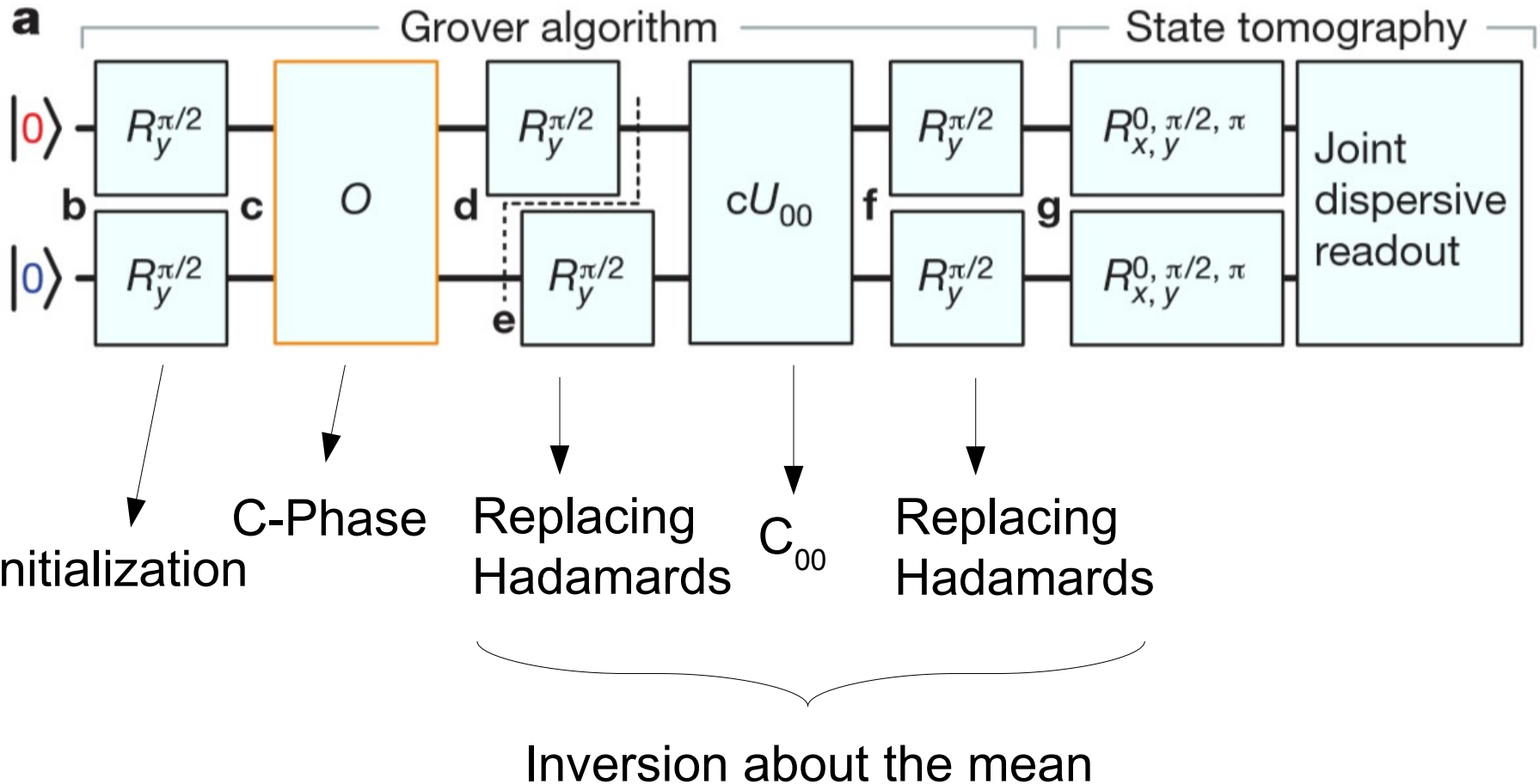
C-Phase gate

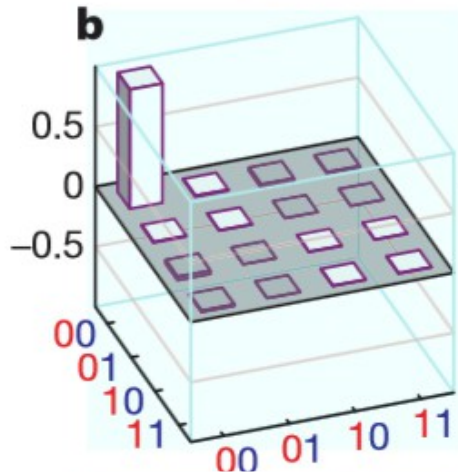
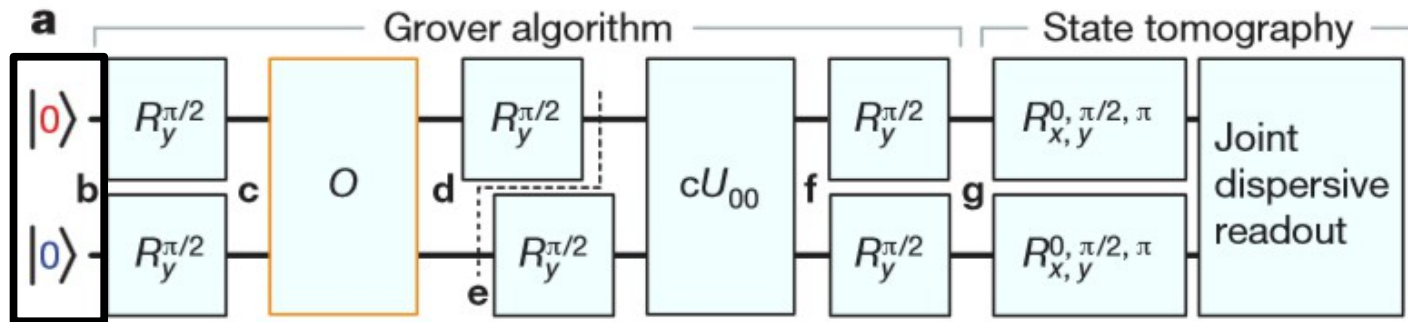
$$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & e^{i\phi_{01}} & 0 & 0 \\ 0 & 0 & e^{i\phi_{10}} & 0 \\ 0 & 0 & 0 & e^{i\phi_{11}} \end{bmatrix}$$



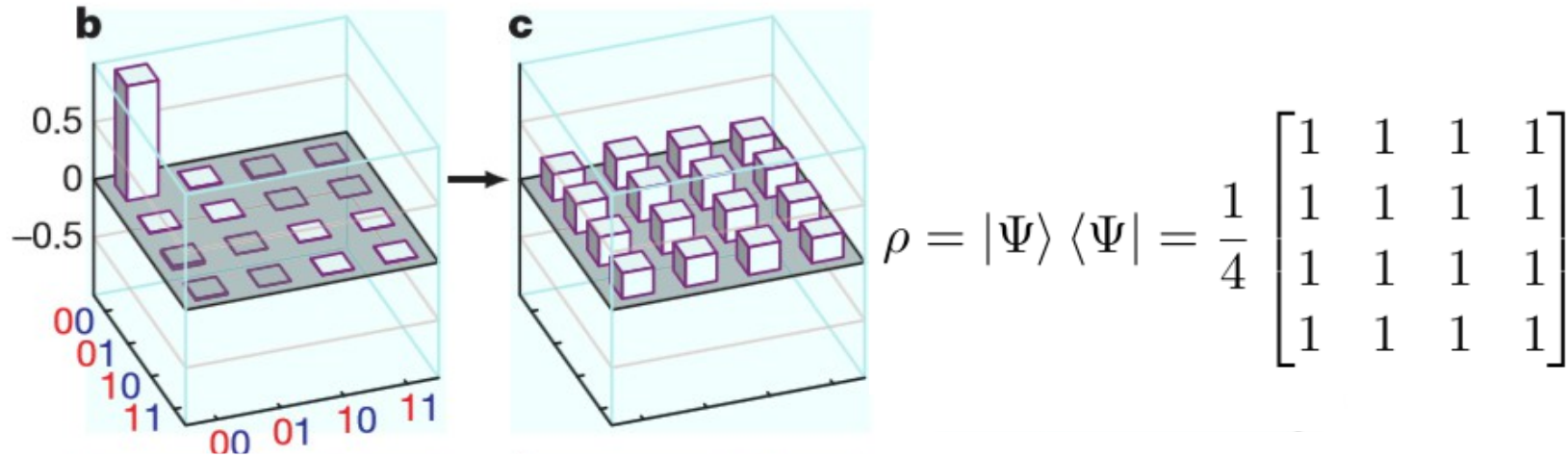
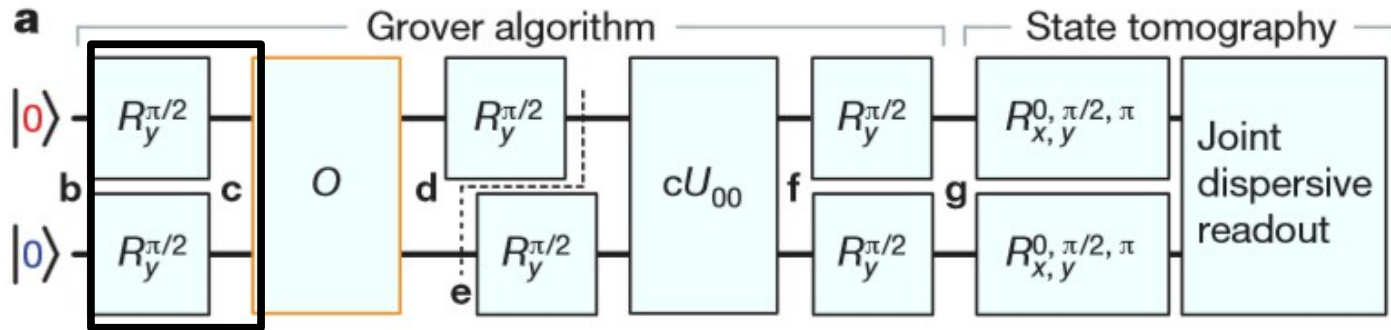
$$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix}$$

# Schematic setup

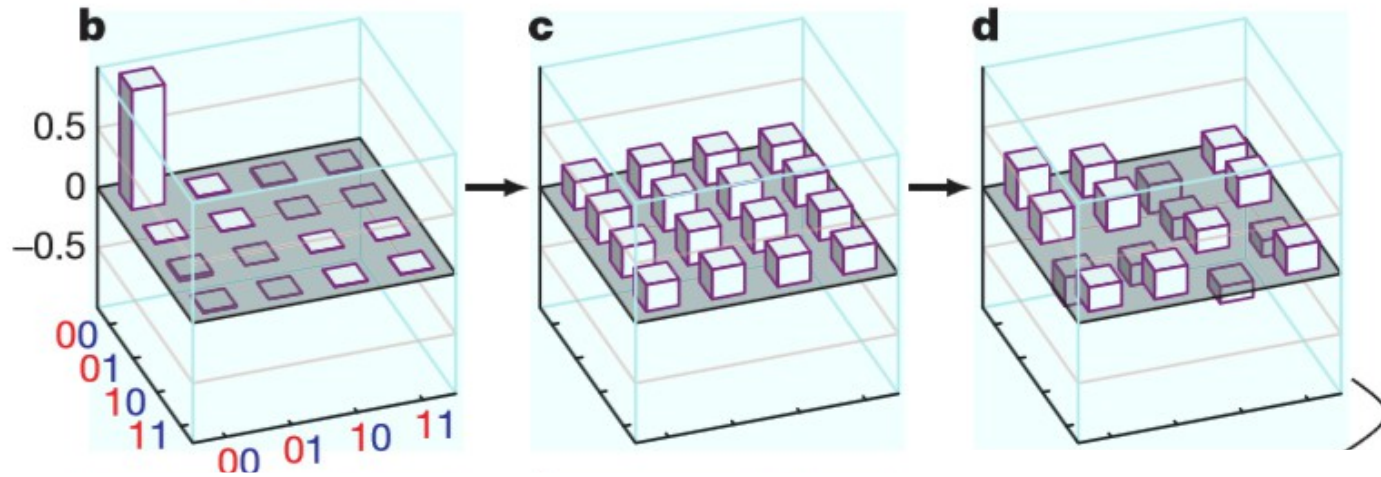
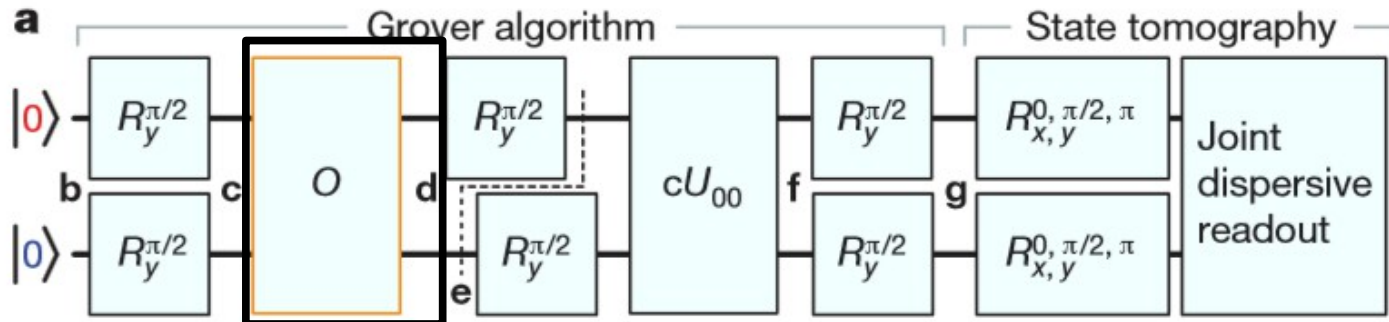




$$|\Psi\rangle = |00\rangle = |0\rangle_1 \otimes |0\rangle_2$$

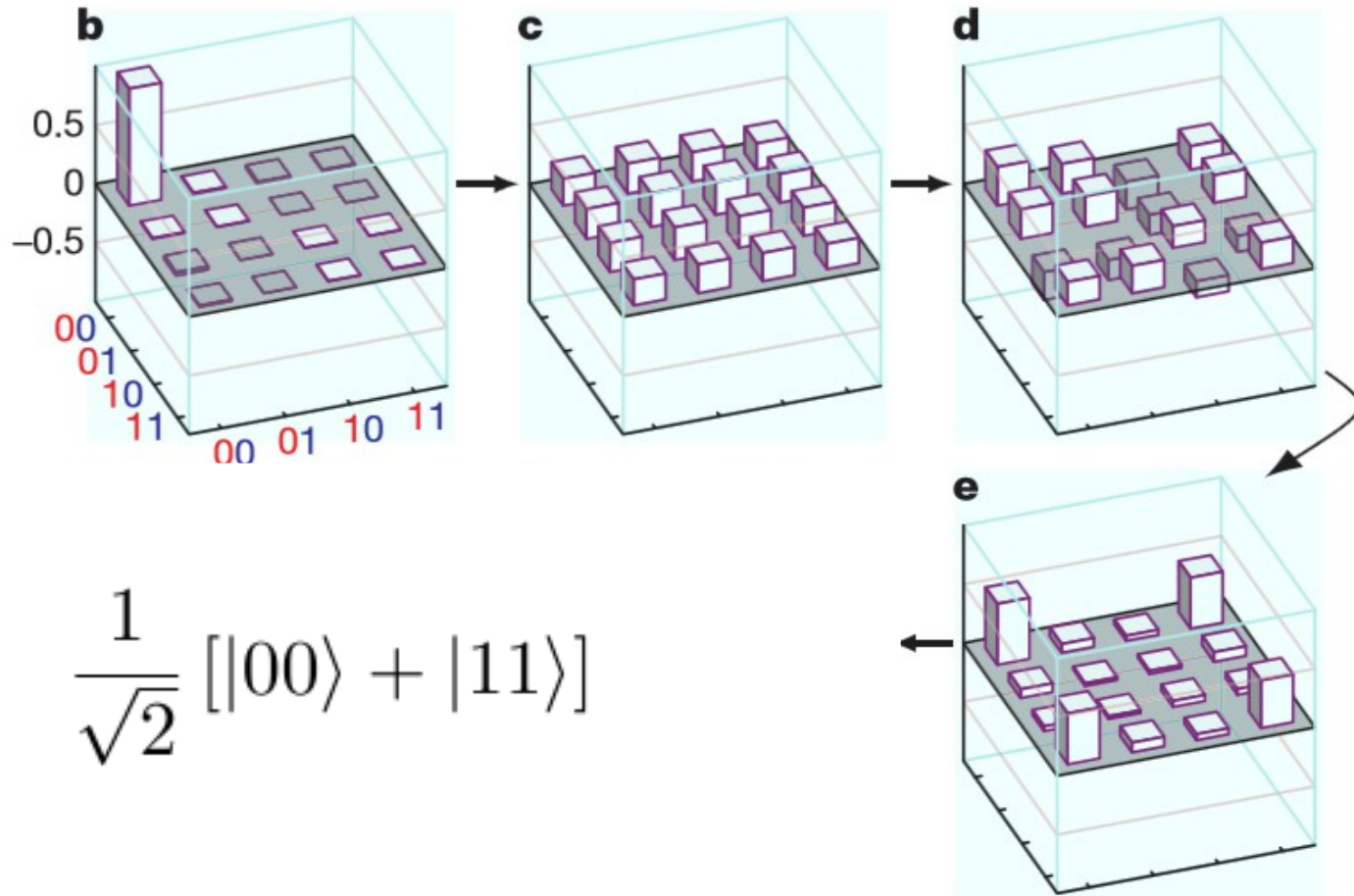
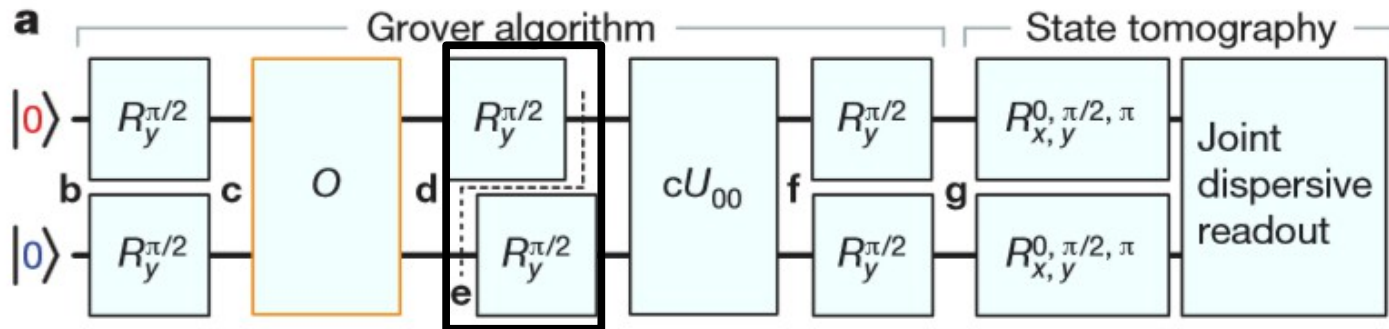


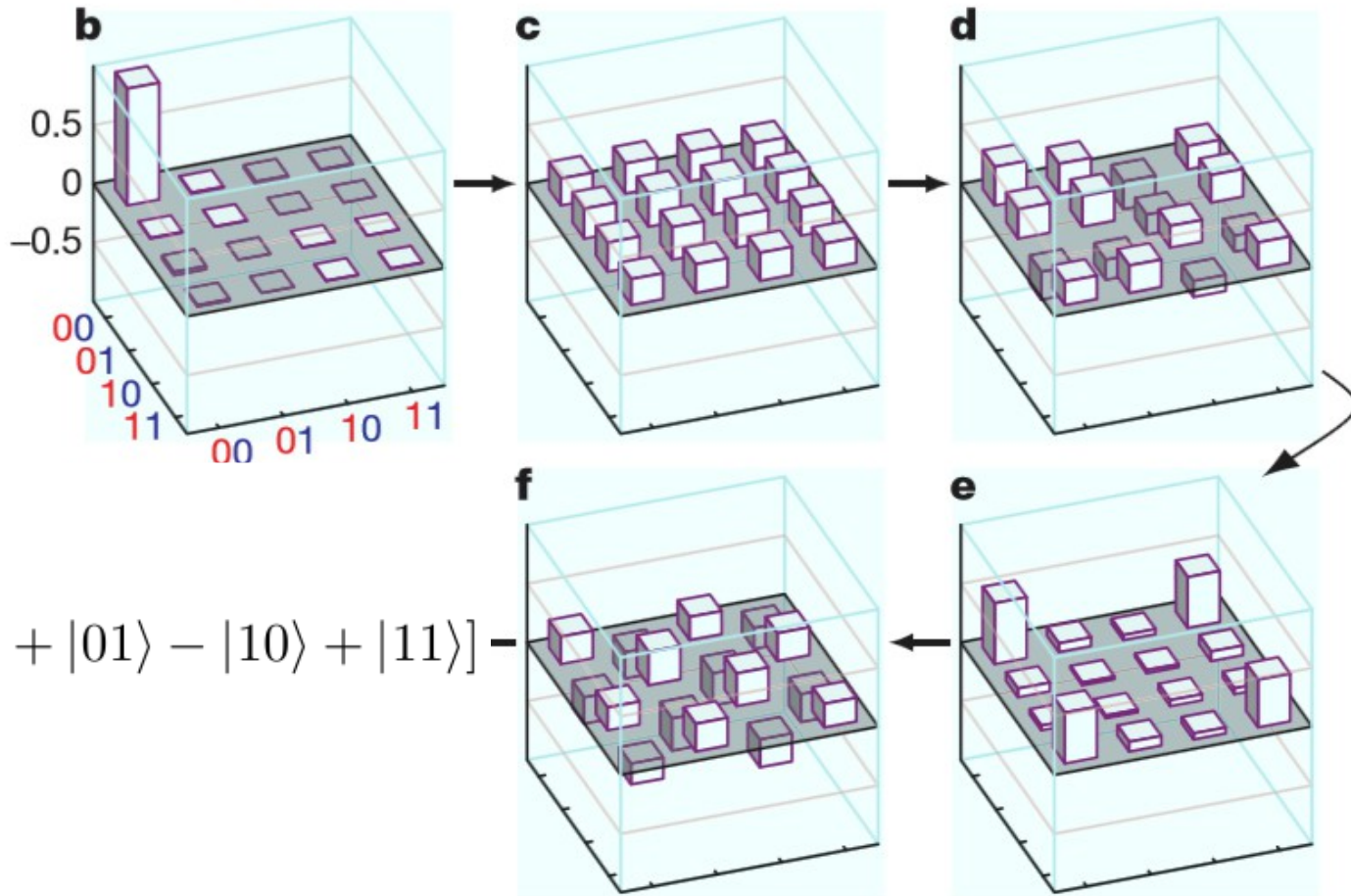
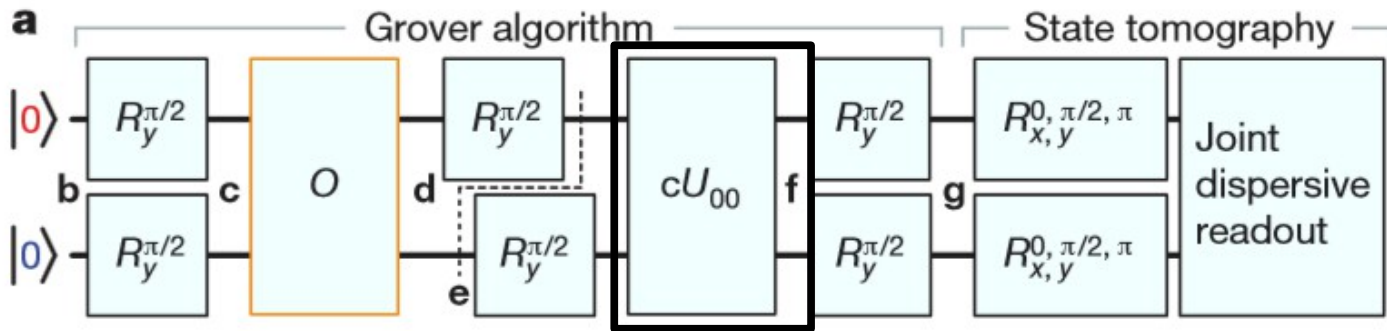
$$R_y^{\pi/2} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & -1 \\ 1 & 1 \end{pmatrix} \longrightarrow \frac{1}{2} [ |00\rangle + |01\rangle + |10\rangle + |11\rangle ]$$



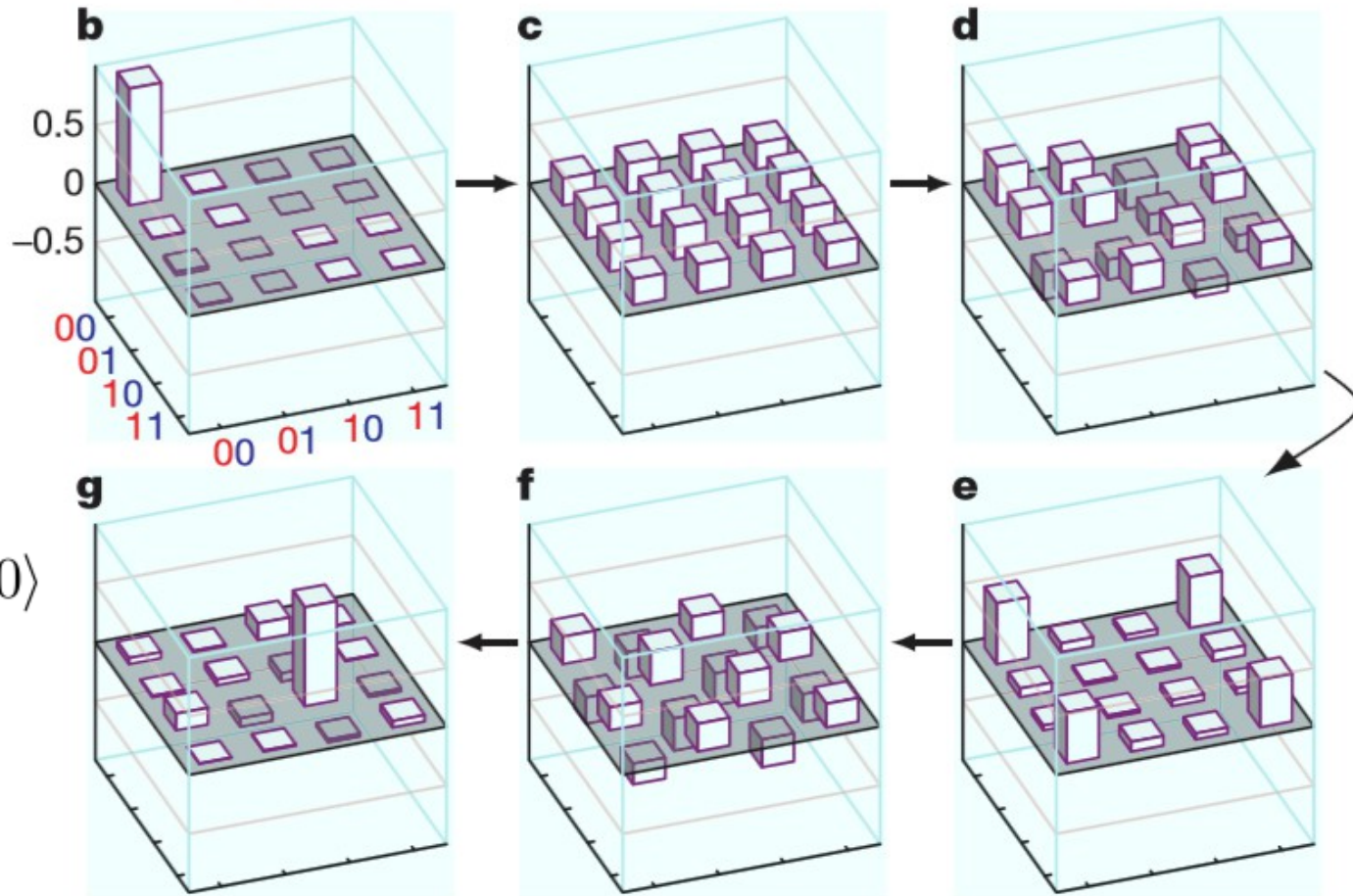
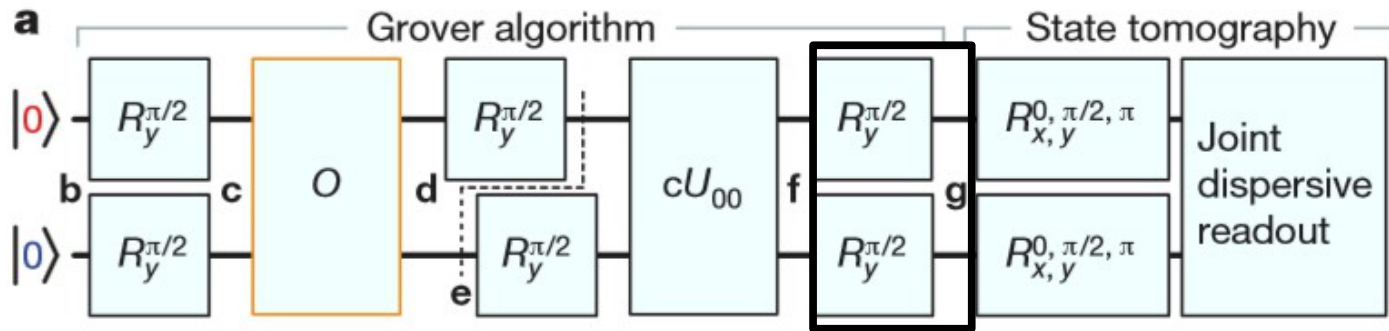
$$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \longrightarrow \frac{1}{2} [ |00\rangle + |01\rangle - |10\rangle + |11\rangle ]$$





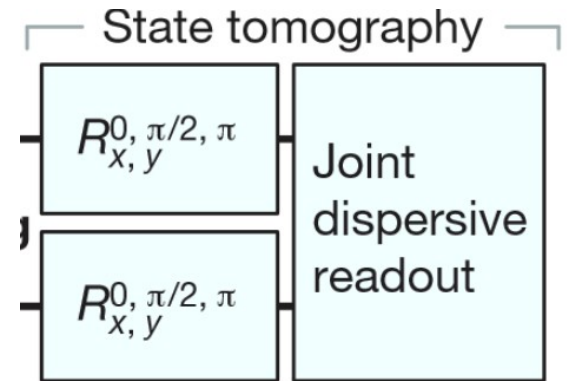


$$\frac{1}{\sqrt{2}} [-|00\rangle + |01\rangle - |10\rangle + |11\rangle]$$



# Tomography: Joint dispersive read out

$$M = \beta_1 \sigma_z^L + \beta_2 \sigma_z^R + \beta_{12} \sigma_z^L \otimes \sigma_z^R$$



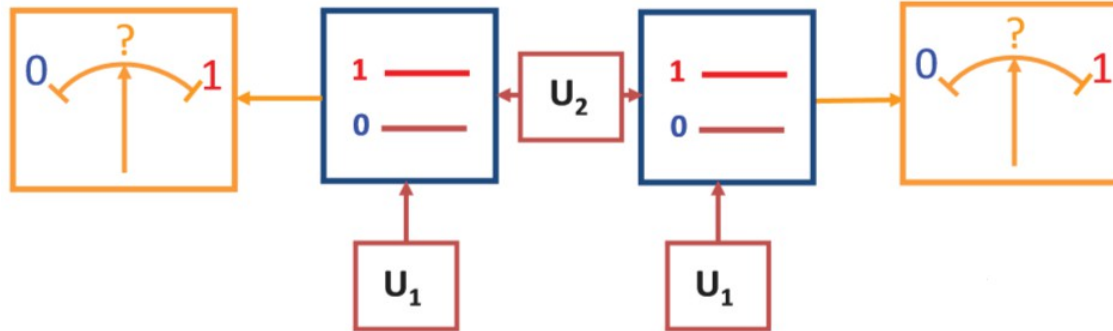
- Dispersive regime, all  $\beta$ 's are approximately equal.
- Single qubit rotations:  $0 \ R_x^{\pi/2} \ R_y^{\pi/2} \ R_x^{\pi}$
- 15 observables (16 -1) by single qubit rotations before measuring M.
- Likelyhood estimator for the averages over 450.000 runs.

# Results DiCarlo et al.

$$F(\rho, \psi) = \langle \psi | \rho | \psi \rangle$$

Element		Grover search oracle*			
		$f_{00}$	$f_{01}$	$f_{10}$	$f_{11}$
$\langle 0,0   \rho   0,0 \rangle$	Ideal	1	0	0	0
	Measured	0.81(1)	0.08(1)	0.07(2)	0.065(7)
$\langle 0,1   \rho   0,1 \rangle$	Ideal	0	1	0	0
	Measured	0.066(7)	0.802(9)	0.05(1)	0.054(8)
$\langle 1,0   \rho   1,0 \rangle$	Ideal	0	0	1	0
	Measured	0.08(1)	0.05(1)	0.82(2)	0.07(1)
$\langle 1,1   \rho   1,1 \rangle$	Ideal	0	0	0	1
	Measured	0.05(2)	0.07(1)	0.06(1)	0.81(1)

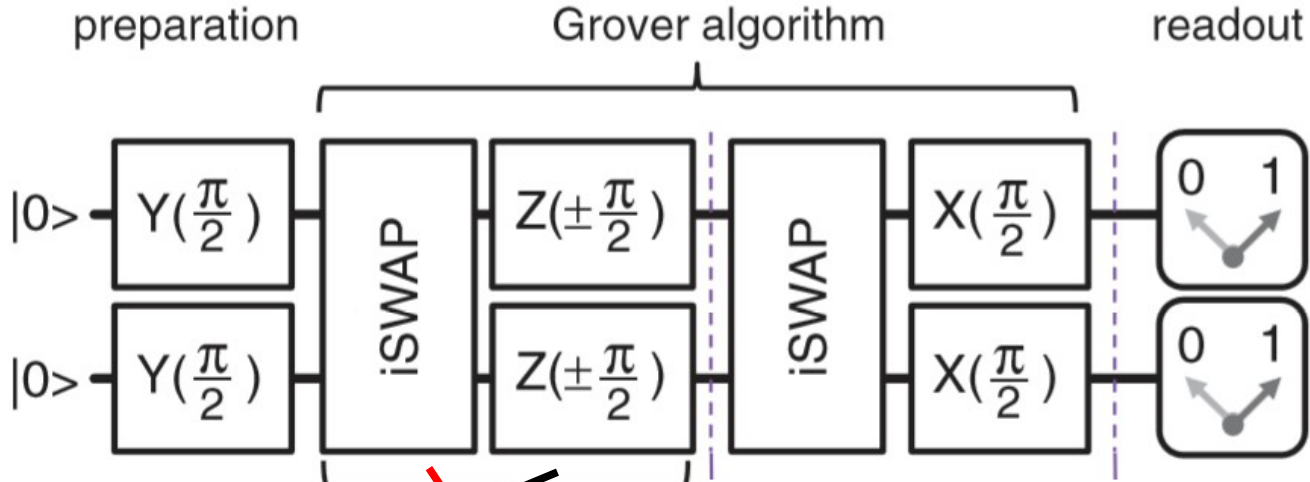
# Experiment 2 : single qubit read out



Dewes et al. Phys. Rev. B 85, 140503(R)



# Schematic setup



oracle

C-Phase & SWAP

$$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & i & 0 & 0 \\ 0 & 0 & -i & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & i & 0 \\ 0 & i & 0 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$



# Results Dewes et al.

State	Tomography (F)	Single shot (P)
00	0.70	0.67
01	0.62	0.55
10	0.67	0.62
11	0.66	0.52

Single shot measurement ~200 ns

# Conclusion

- Grover search algorithm

$$\left( \{2|\psi\rangle\langle\psi| - I\} \hat{O} \right)^R H^{\otimes n} |0\rangle$$

- Experimentally realizable in superconducting circuits
  - Oracle ~ C-Phase
  - average about mean ~ H + C-Phase + H
- Tomography fidelity > 0.8, qubit-resonator-qubit
- Single shot probability > 0.5, qubit-capacitor-qubit

# References

- *Quantum speeding-up of computation demonstrated in a superconducting two-qubit processor.* A. Dewes et al. Phys. Rev. B 85, 140503(R) (2012), arxiv:1110.5170
- *Demonstration of two-qubit algorithms using a superconducting quantum processor,* L. DiCarlo et al. Nature 460, 240 (2009).
- *Quantum Computation and Quantum Information,* Nielsen and Chuang.