

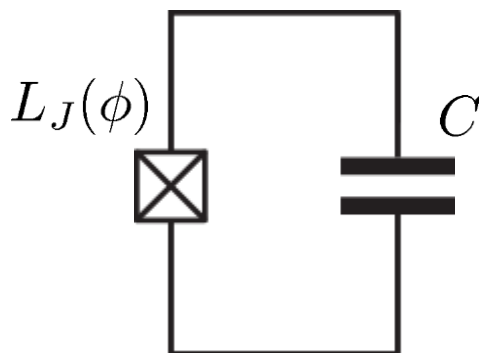
Constructing Non-Linear Quantum Electronic Circuits

circuit elements:



Josephson junction:
a non-dissipative nonlinear
element (inductor)

anharmonic oscillator:



$$L_J(\phi) = \left(\frac{\partial I}{\partial \phi} \right)^{-1}$$
$$= \frac{\phi_0}{2\pi I_c} \frac{1}{\cos(2\pi\phi/\phi_0)}$$

non-linear energy
level spectrum:



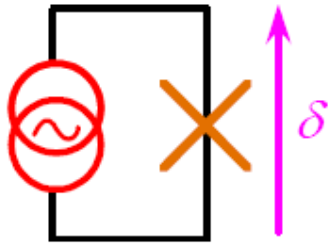
electronic
artificial atom

A Classification of Josephson Junction Based Qubits

How to make use in of Jospelson junctions in a qubit?

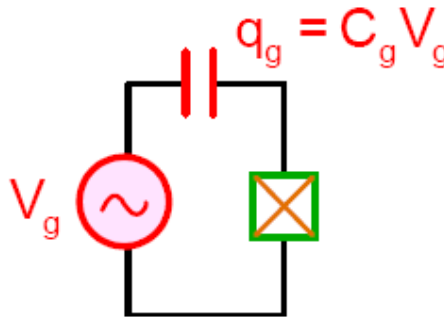
Common options of bias (control) circuits:

phase qubit



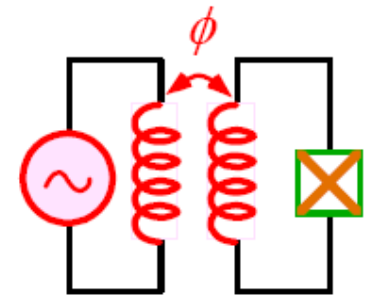
current bias

charge qubit
(Cooper Pair Box, Transmon)



charge bias

flux qubit

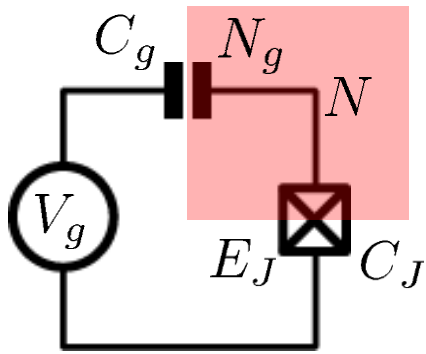


flux bias

How is the control circuit important?

The Cooper Pair Box Qubit

A Charge Qubit: The Cooper Pair Box



discrete charge on island:

$$N = \frac{Q}{2e}$$

continuous gate charge:

$$N_g = \frac{C_g V_g}{2e}$$

total box capacitance

$$C_\Sigma = C_g + C_J$$

Hamiltonian: $H = H_{\text{el}} + H_{\text{mag}}$

electrostatic part:

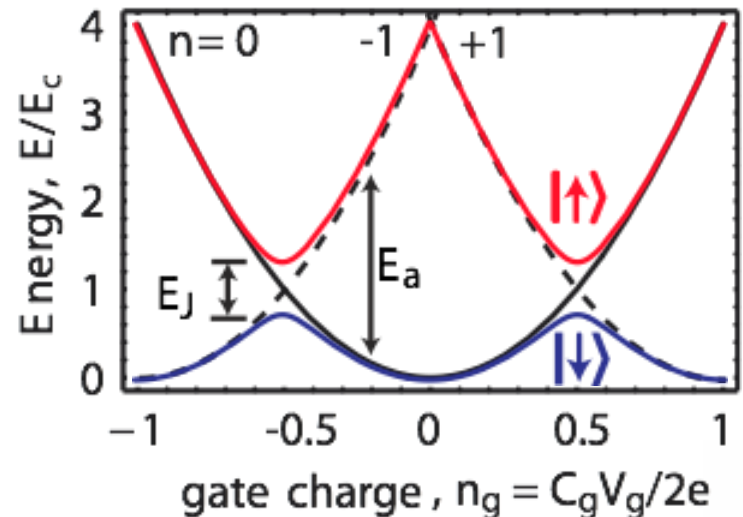
$$H_{\text{el}} = \frac{Q^2}{2C} = \frac{(2e)^2}{2C_\Sigma} (N - N_g)^2$$

charging energy E_C

magnetic part:

$$H_{\text{mag}} = -E_J \cos \delta \approx \frac{\phi^2}{2L_{J0}}$$

Josephson energy



Hamilton Operator of the Cooper Pair Box

Hamiltonian: $\hat{H} = \hat{H}_{\text{el}} + \hat{H}_{\text{mag}} = E_C(\hat{N} - N_g)^2 - E_J \cos \hat{\delta}$

commutation relation: $[\hat{\delta}, \hat{N}] = i$ $\cos \hat{\delta} = \frac{1}{2}(e^{i\hat{\delta}} + e^{-i\hat{\delta}})$

charge number operator: $\hat{N}|N\rangle = N|N\rangle$ eigenvalues, eigenfunctions

$$\sum_N |N\rangle\langle N| = 1 \quad \text{completeness}$$

$$\langle N|M\rangle = \delta_{NM} \quad \text{orthogonality}$$

phase basis: $|\delta\rangle = \frac{1}{\sqrt{2\pi}} \sum_N e^{iN\delta} |N\rangle$ basis transformation

$$e^{\pm i\hat{\delta}} |N\rangle = |N \pm 1\rangle$$

Solving the Cooper Pair Box Hamiltonian

Hamilton operator in the **charge basis** N :

$$\hat{H} = \sum_N \left[E_C (N - N_g)^2 |N\rangle\langle N| - \frac{E_J}{2} (|N\rangle\langle N+1| + |N+1\rangle\langle N|) \right]$$

solutions in the charge basis:

$$\hat{H}|\psi_n(N)\rangle = E_n|\psi_n(N)\rangle$$

Hamilton operator in the **phase basis** δ :

$$\hat{H} = E_C (\hat{N} - N_g)^2 - E_J \cos \hat{\delta} = E_C \left(-i \frac{\partial}{\partial \delta} - N_g \right)^2 - E_J \cos \hat{\delta}$$

transformation of the number operator:

$$\hat{N} = \frac{\hat{Q}}{2e} = -i\hbar \frac{1}{2e} \frac{\partial}{\partial \phi} = -i \frac{\partial}{\partial \delta}$$

solutions in the phase basis:

$$\hat{H}|\psi_n(\delta)\rangle = E_n|\psi_n(\delta)\rangle$$

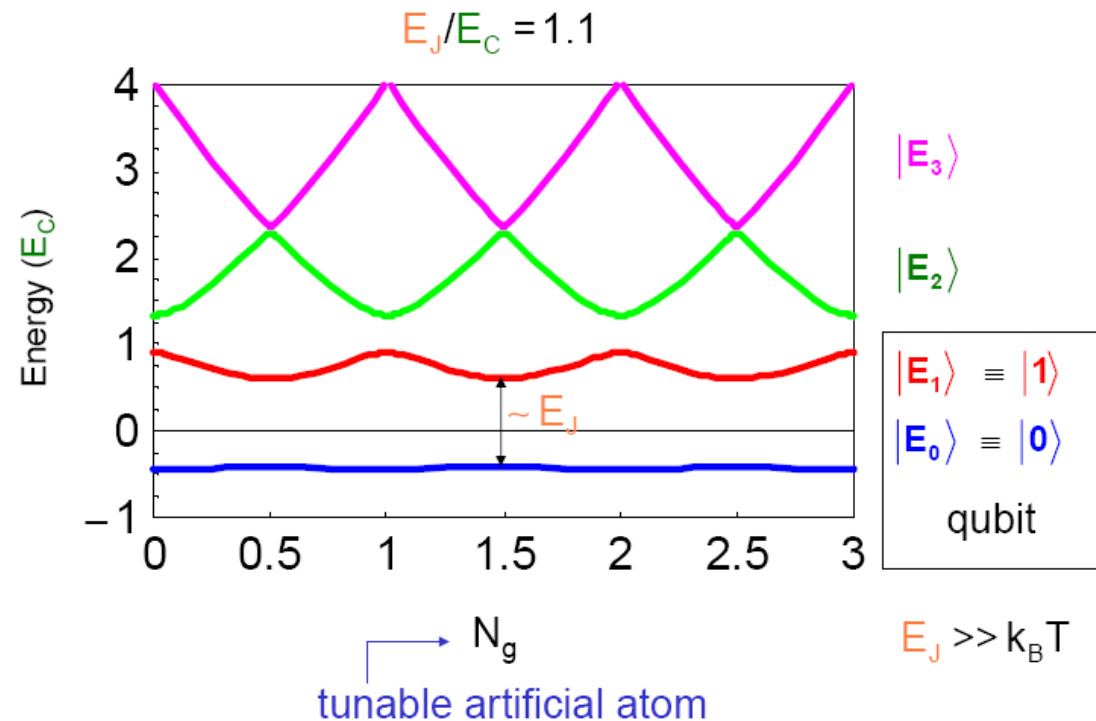
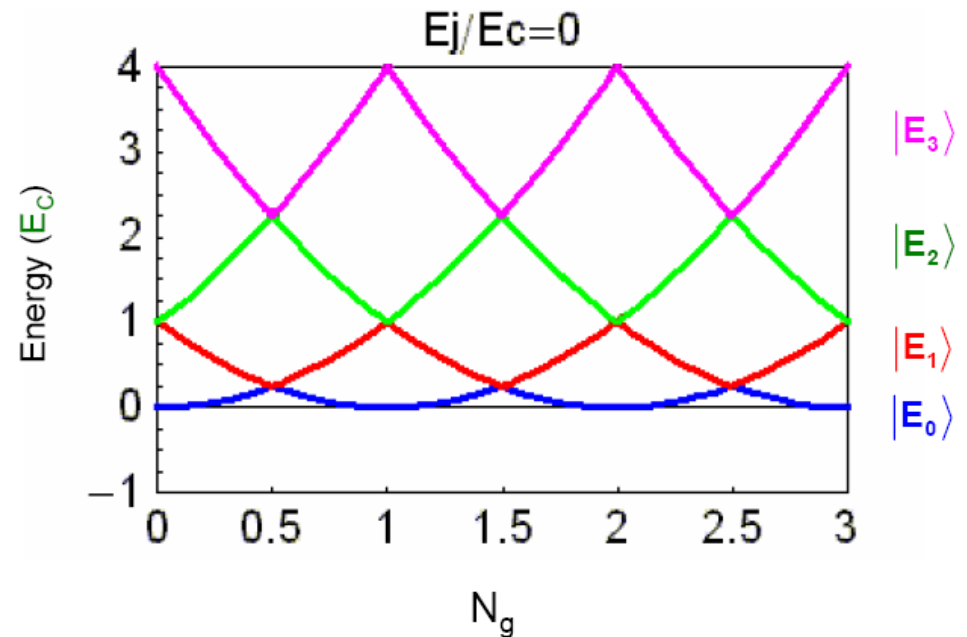
Energy Levels

energy level diagram for $E_J=0$:

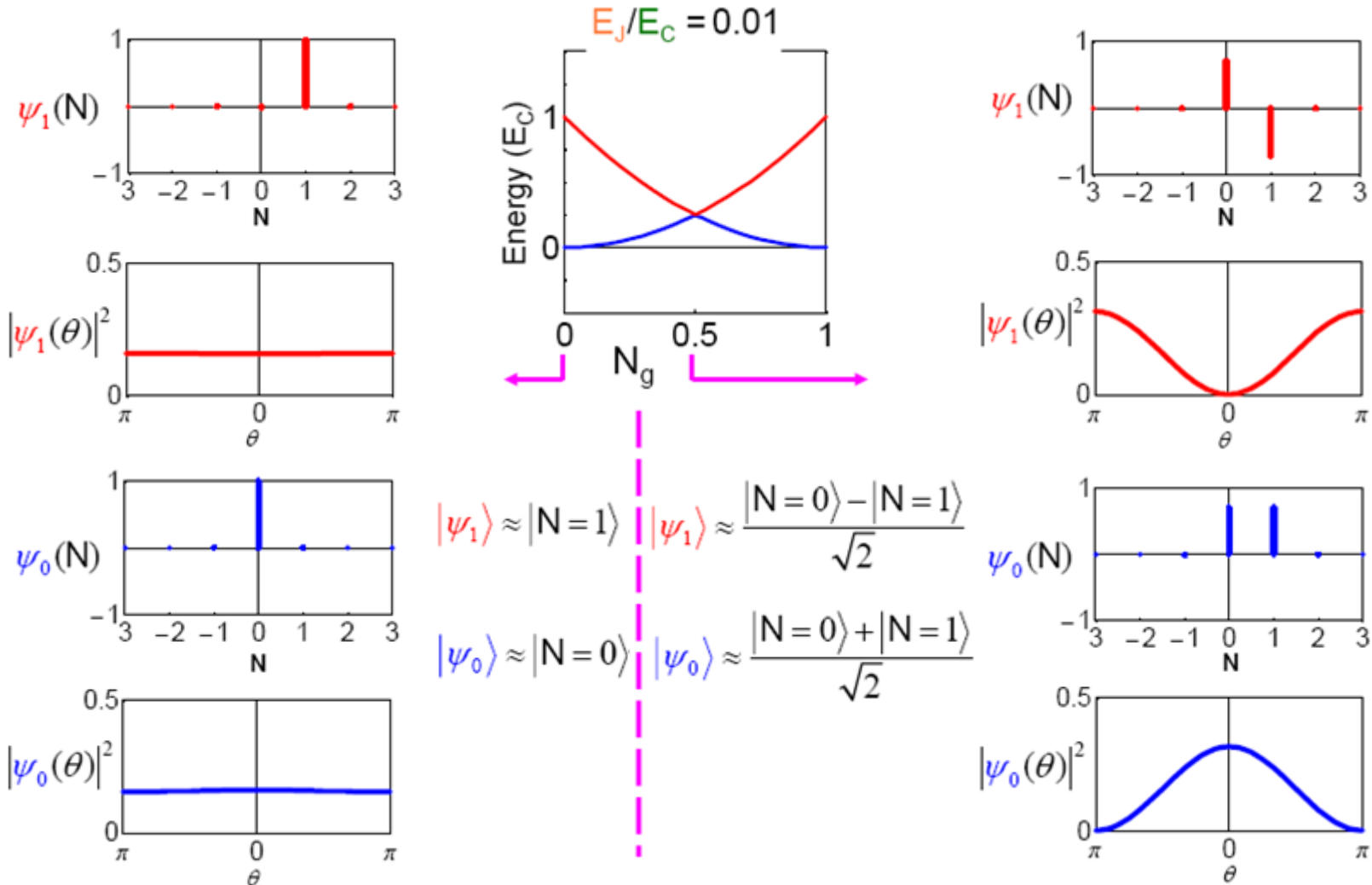
- energy bands are formed
- bands are periodic in N_g

energy bands for finite E_J

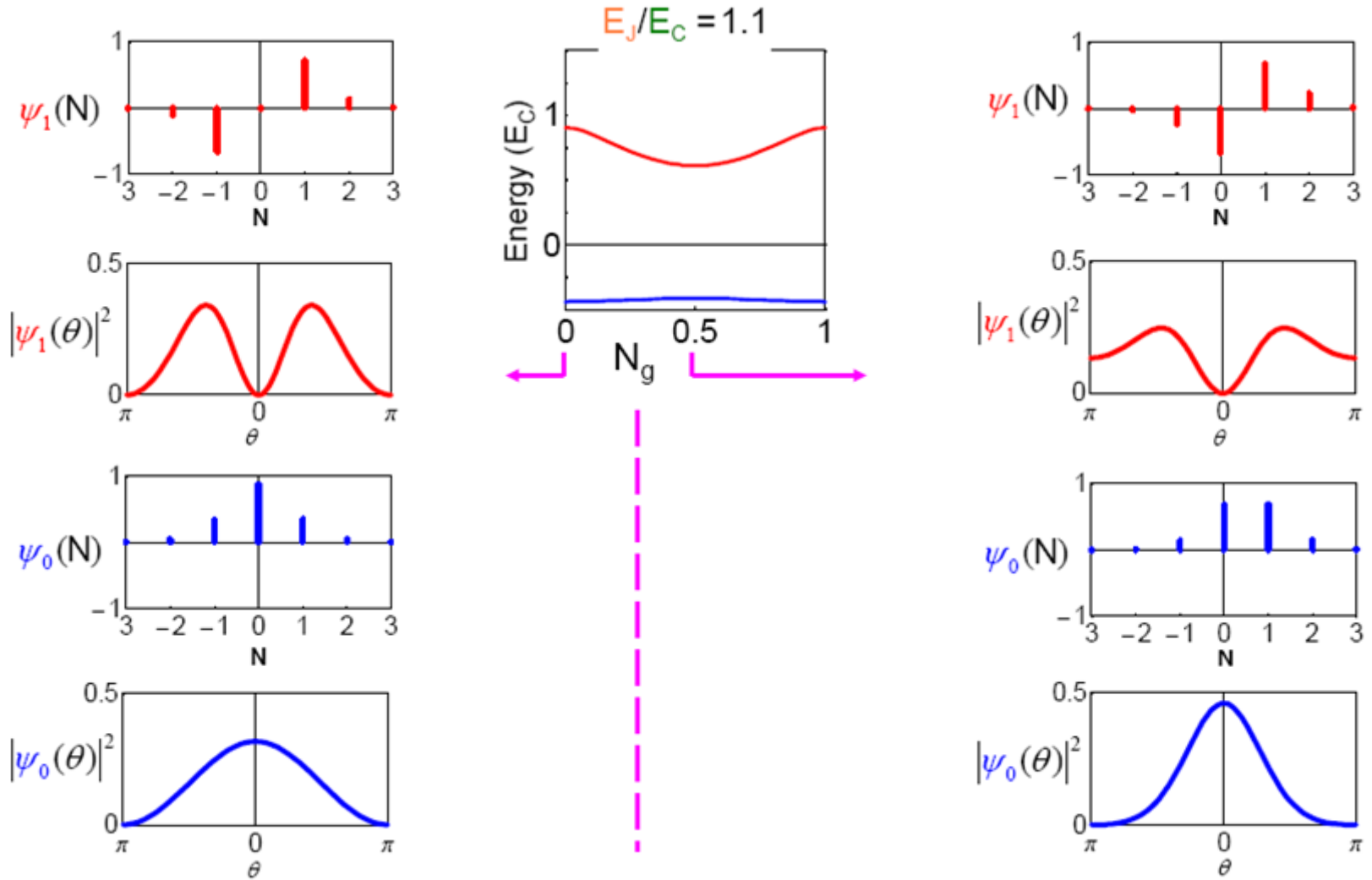
- Josephson coupling lifts degeneracy
- E_J scales level separation at charge degeneracy



Charge and Phase Wave Functions ($E_J \ll E_C$)

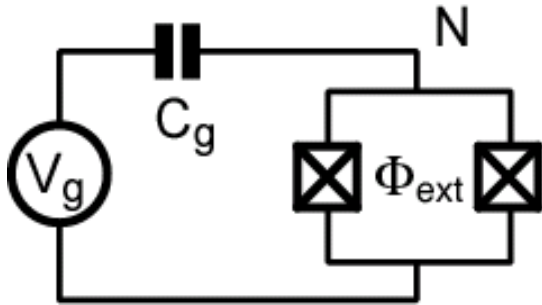


Charge and Phase Wave Functions ($E_J \sim E_C$)



Tuning the Josephson Energy

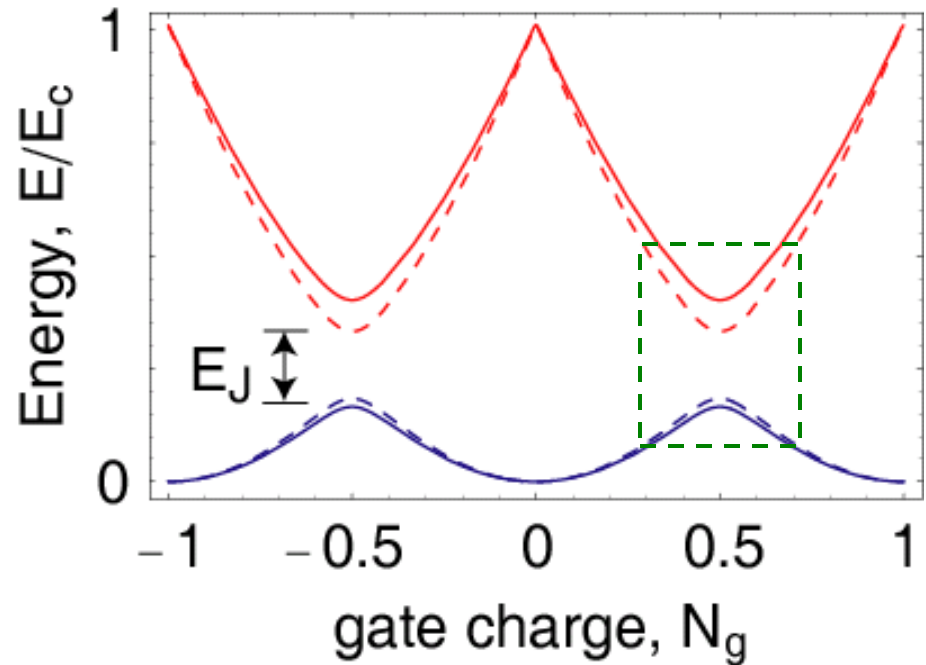
split Cooper pair box in perpendicular field



$$H = E_C (N - N_g)^2 - E_{J,\max} \cos \left(\pi \frac{\phi_{\text{ext}}}{\phi_0} \right)$$

SQUID modulation of Josephson energy

$$E_J = E_{J,\max} \cos \left(\pi \frac{\phi_{\text{ext}}}{\phi_0} \right)$$



consider two state approximation

Two-State Approximation

$$\hat{H}_{\text{CPB}} = \hat{H}_{\text{el}} + \hat{H}_J = E_C(\hat{N} - N_g)^2 - E_J \cos \hat{\delta}$$

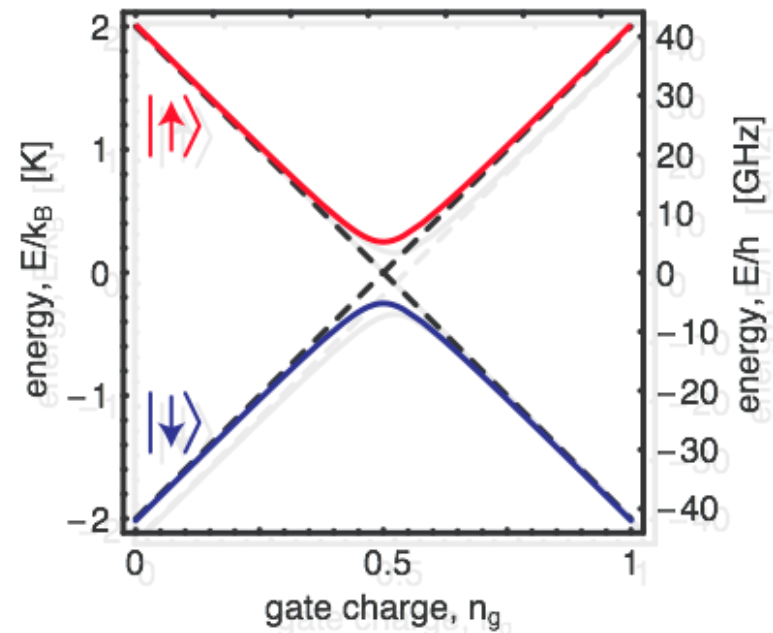
$$\hat{H}_{\text{CPB}} = \sum_N \left[E_C(N - N_g)^2 |N\rangle \langle N| - \frac{E_J}{2} (|N\rangle \langle N+1| + |N+1\rangle \langle N|) \right]$$

Restricting to a two-charge Hilbert space:

$$\hat{N} = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} = \frac{1 - \hat{\sigma}_z}{2}$$

$$\cos \hat{\delta} = \frac{\hat{\sigma}_x}{2}$$

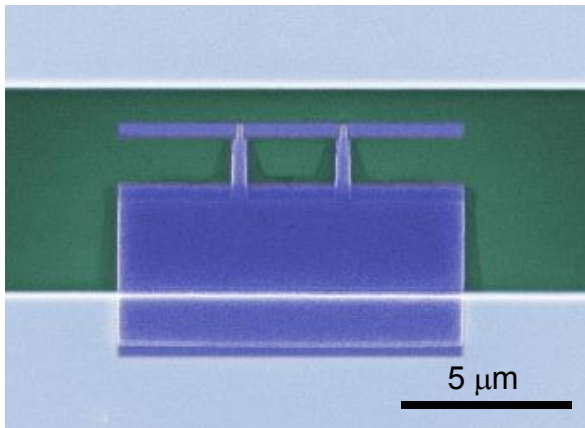
$$\begin{aligned} \hat{H} &= -\frac{E_C}{2}(1 - 2N_g)\hat{\sigma}_z - \frac{E_J}{2}\hat{\sigma}_x \\ &= -\frac{1}{2}(E_{\text{el}}\hat{\sigma}_z + E_J\hat{\sigma}_x) \end{aligned}$$



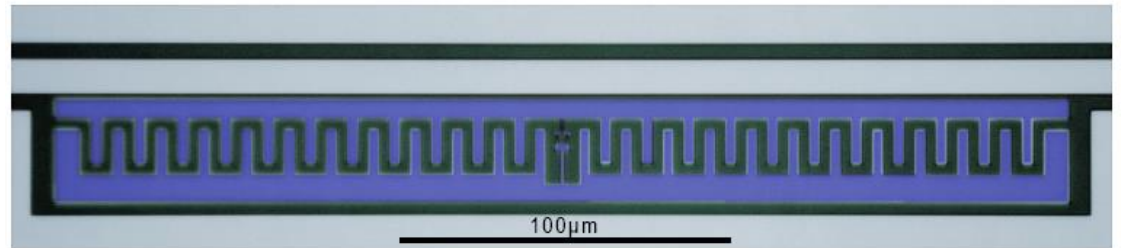
A Variant of the Cooper Pair Box

a Cooper pair box with a small charging energy

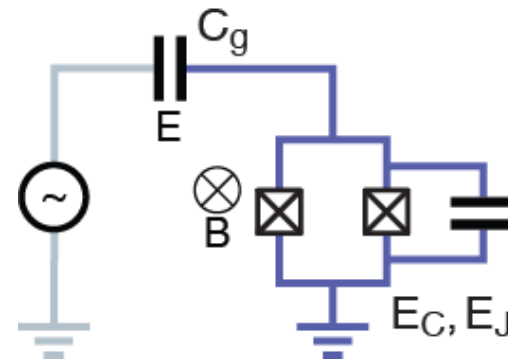
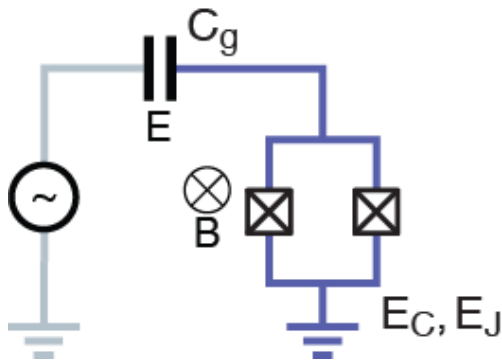
standard CPB:



Transmon qubit:



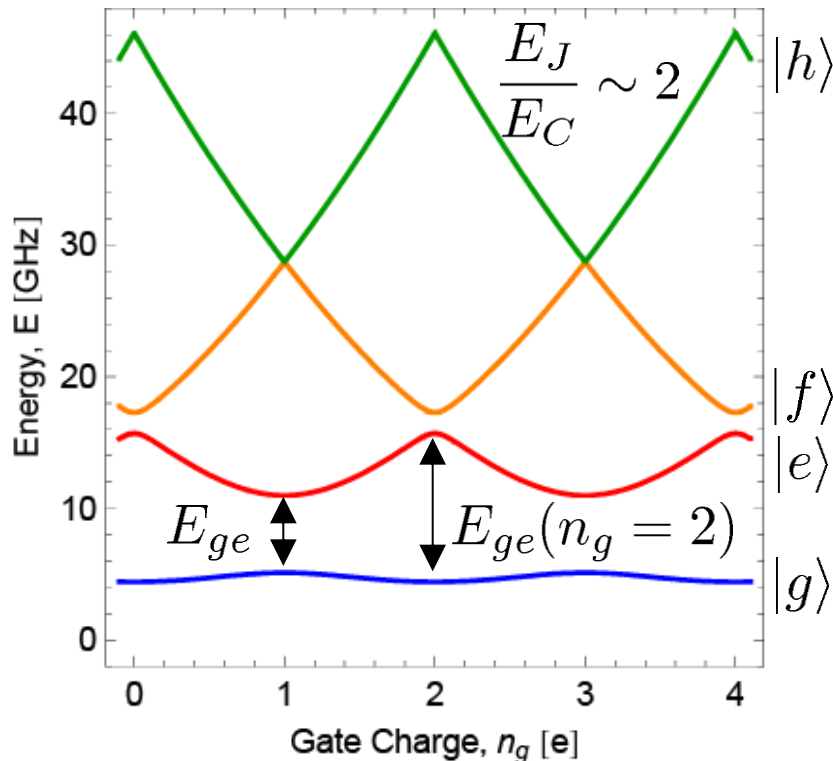
circuit diagram:



J. Koch *et al.*, Phys. Rev. A **76**, 042319 (2007)
J. Schreier *et al.*, Phys. Rev. B **77**, 180502 (2008)

The Transmon: A Charge Noise Insensitive Qubit

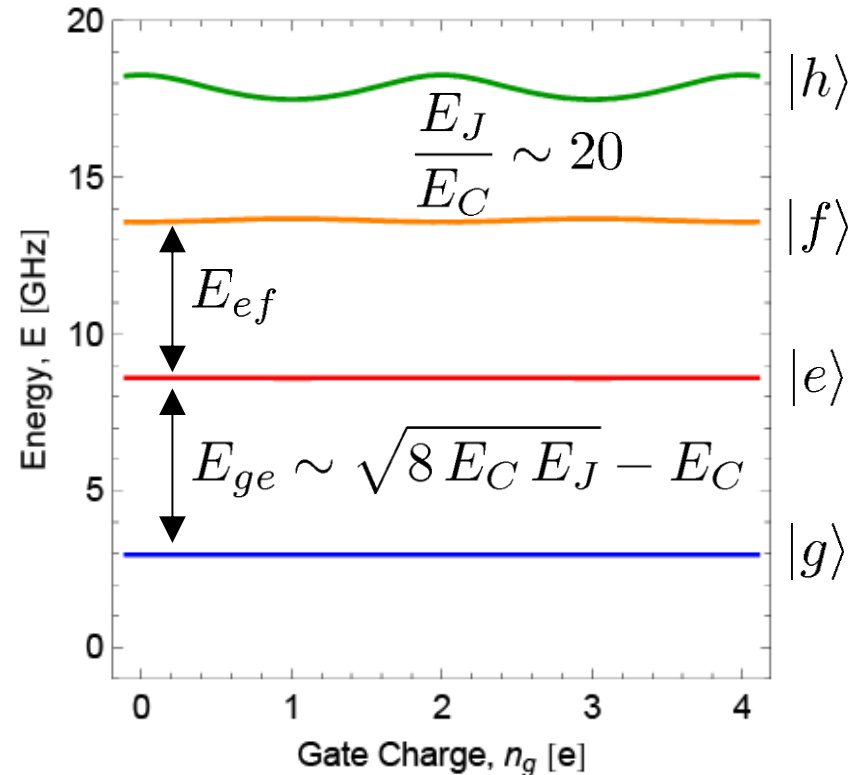
Cooper pair box energy levels:



dispersion:

$$\epsilon = E_{ge}(n_g = 1) - E_{ge}(n_g = 2)$$

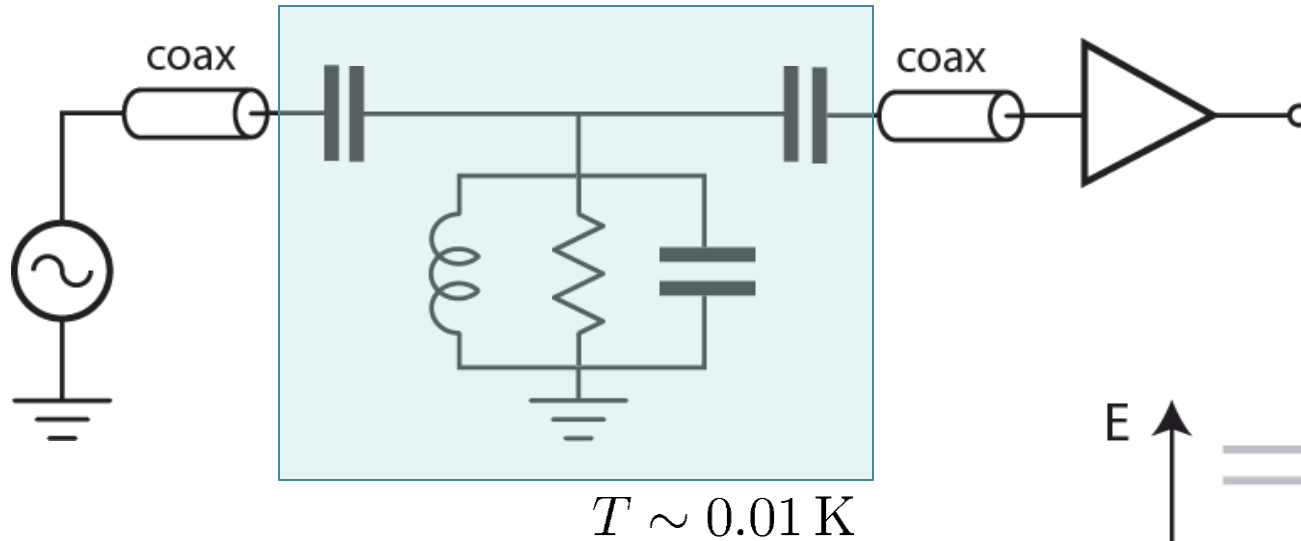
Transmon energy levels:



relative anharmonicity:

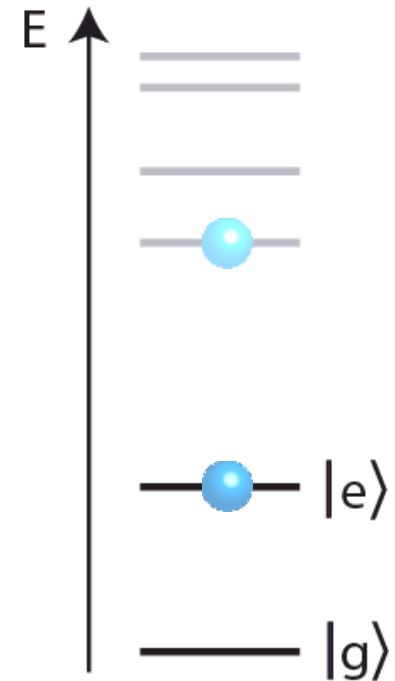
$$\alpha_r = \frac{E_{ef} - E_{ge}}{E_{ge}}$$

How to Operate Circuits Quantum Mechanically?



recipe:

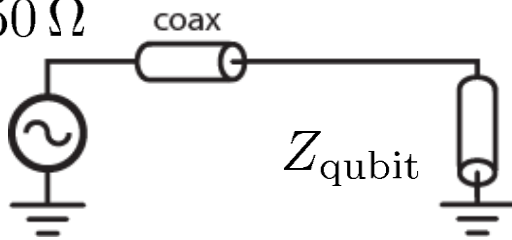
- avoid dissipation
- work at low temperatures
- isolate quantum circuit from environment



Control of Coupling to Electromagnetic Environment

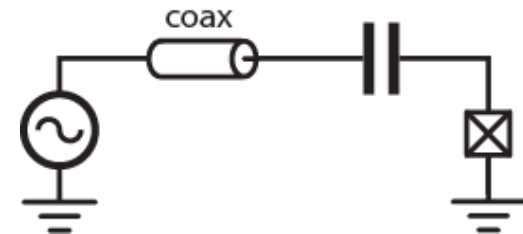
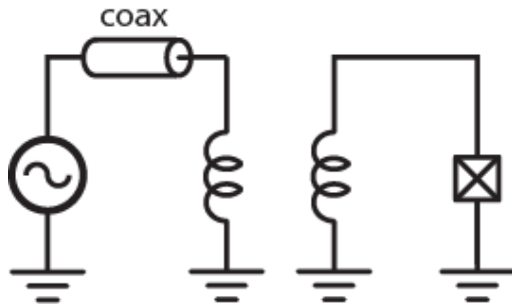
coupling to environment (bias wires):

$$Z_{\text{line}} \sim 50 \Omega$$

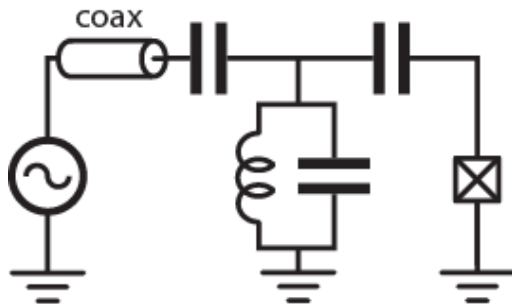


decoherence due to energy relaxation stimulated by the vacuum fluctuations of the environment (spontaneous emission)

Decoupling schemes using non-resonant impedance transformers ...

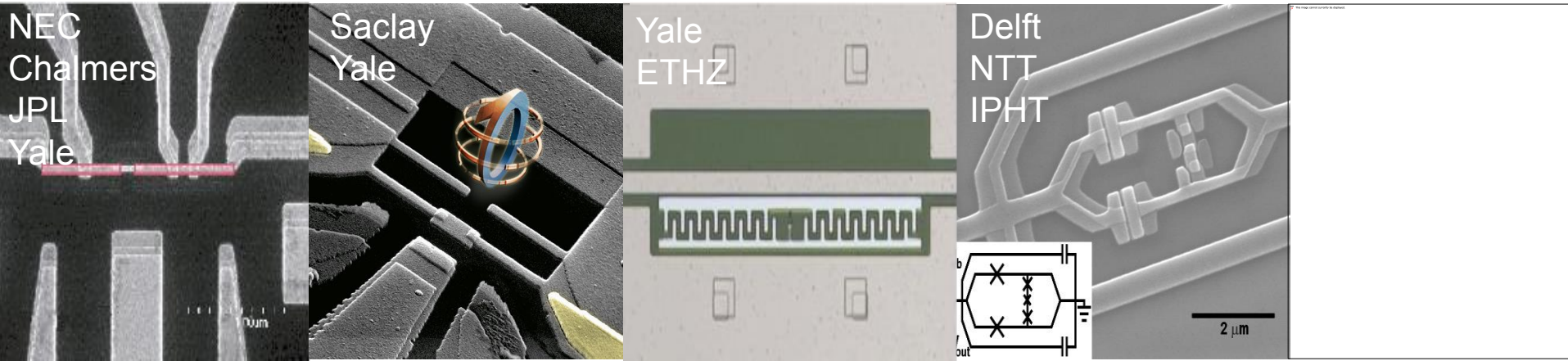


... or resonant impedance transformers



control spontaneous emission by circuit design

Realizations of Superconducting Artificial Atoms



'artificial atoms' -- single superconducting qubits

review:

J. Clarke and F. Wilhelm
Nature 453, 1031 (2008)

'artificial molecules' -- coupled superconducting qubits

