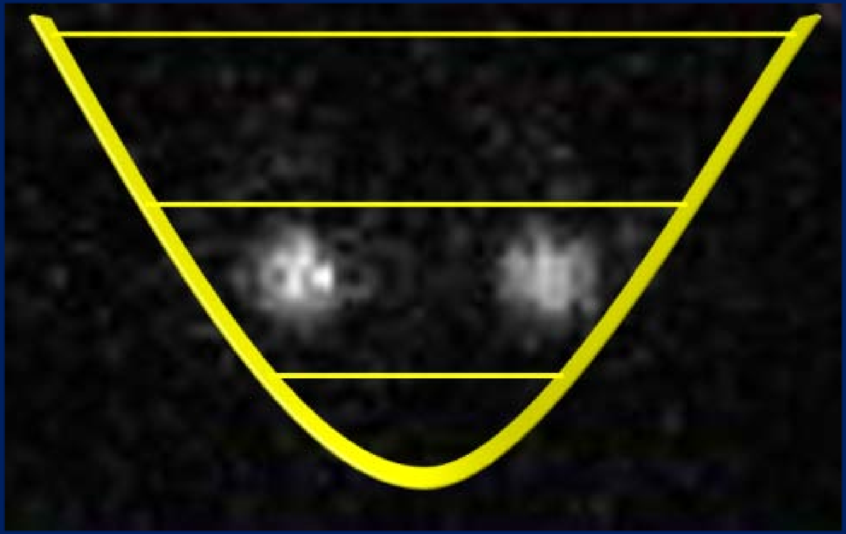


ETH
Eidgenössische Technische Hochschule Zürich
Swiss Federal Institute of Technology Zurich

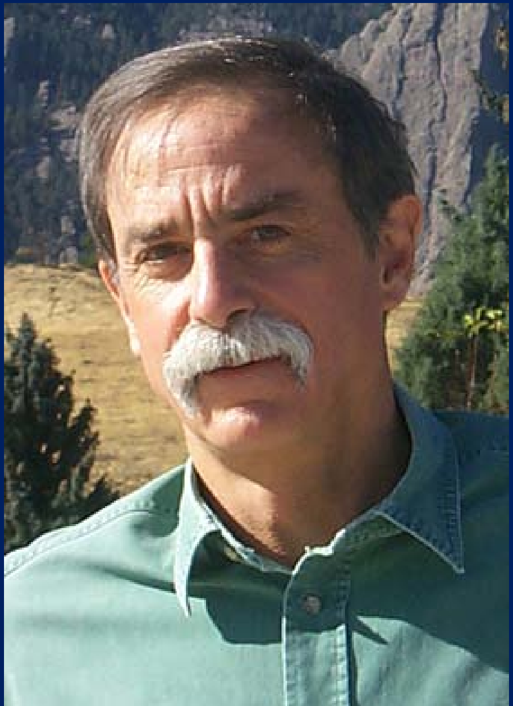
Quantum computation and simulation with cold ions

Jonathan Home

Dave Wineland and Serge Haroche

Nobel Prize in Physics 2012

"for ground-breaking experimental methods that enable measuring and manipulation of individual quantum systems"

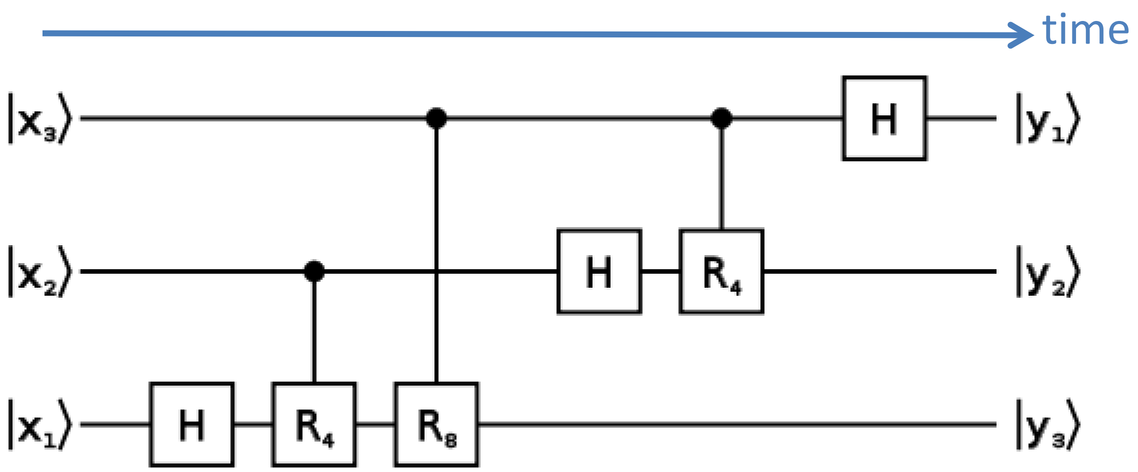


Quantum computation

(Precision control of large-scale quantum mechanical systems)

David Deutsch: Collection of two-state quantum systems (qubits)

Initialise qubit to single state



Detect qubit state

Operations which manipulate isolated qubits or pairs of qubits

Large scale device:

Transport information around processor/distribute entangled states

Perform operations accurately enough to achieve fault-tolerant error-correction (reach arbitrary size system) – accuracy of 0.9999 required

Isolating single charged atoms

Laplace‘s equation
– no chance to trap with static fields

$$\frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2} = 0$$

Paul trap: Use a ponderomotive potential – change potential fast compared to speed of ion

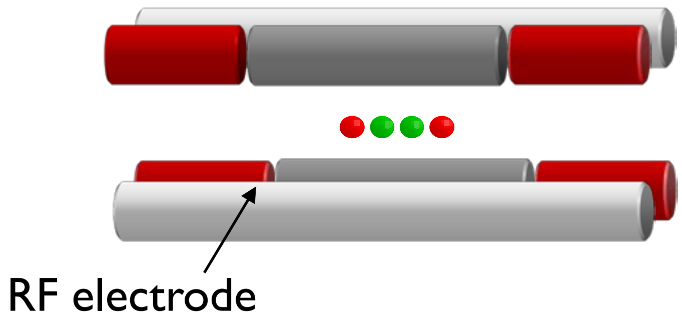
$$\frac{\partial^2 V}{\partial x^2} + \left(\frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2}\right)\cos(\Omega t)$$

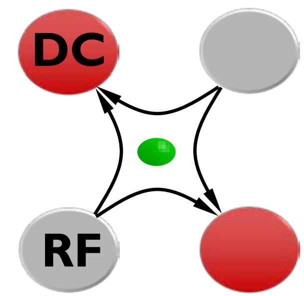
$$M\frac{d^2 x}{dt^2} = qE\cos\Omega t$$

$$\frac{1}{2}M\left(\frac{dx}{dt}\right)^2 = U_{\mathrm{PP}} = \frac{q^2 E^2}{2M\Omega^2}\sin^2\Omega t$$

Time average - Effective potential energy which is minimal at minimum *E*

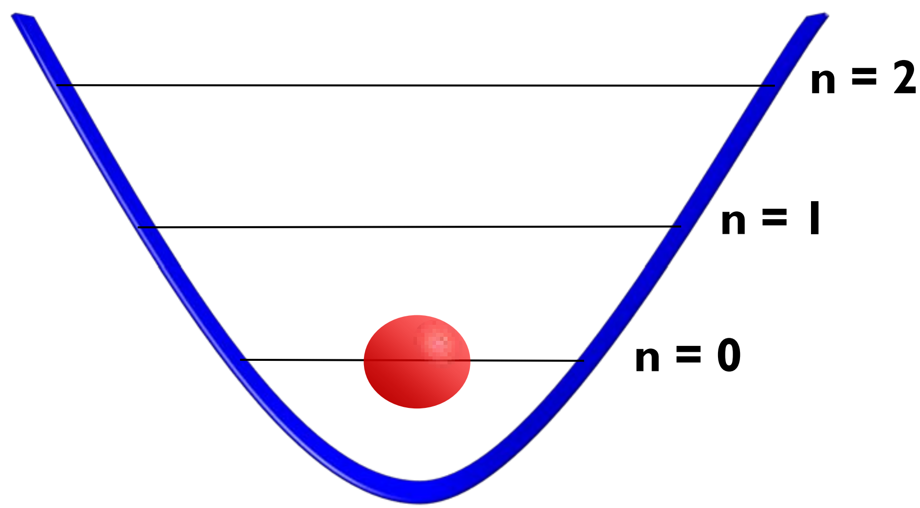
Traps – traditional style



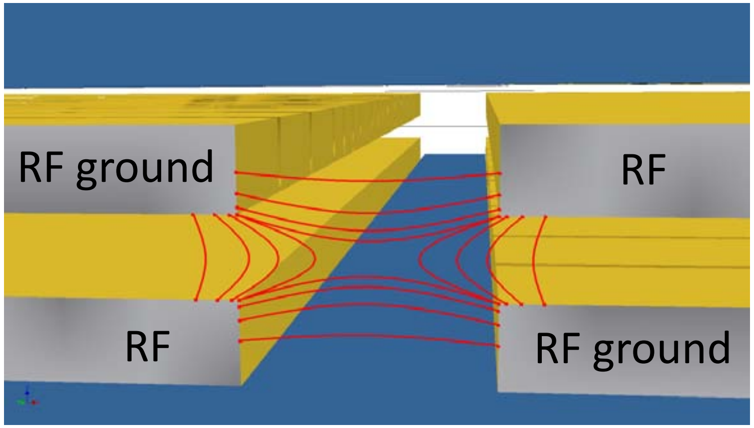


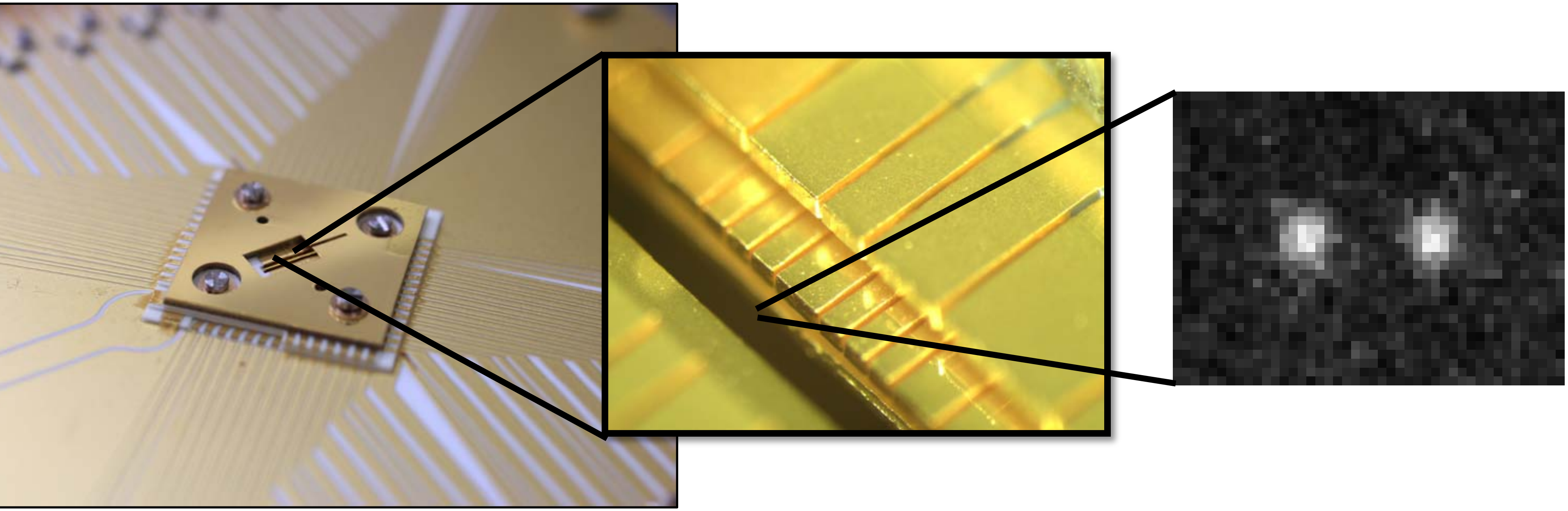
Axial potential gives almost ideal harmonic behaviour

$$\hat{H} = \hbar\omega(\hat{a}^{\dagger}\hat{a} + 1/2)$$



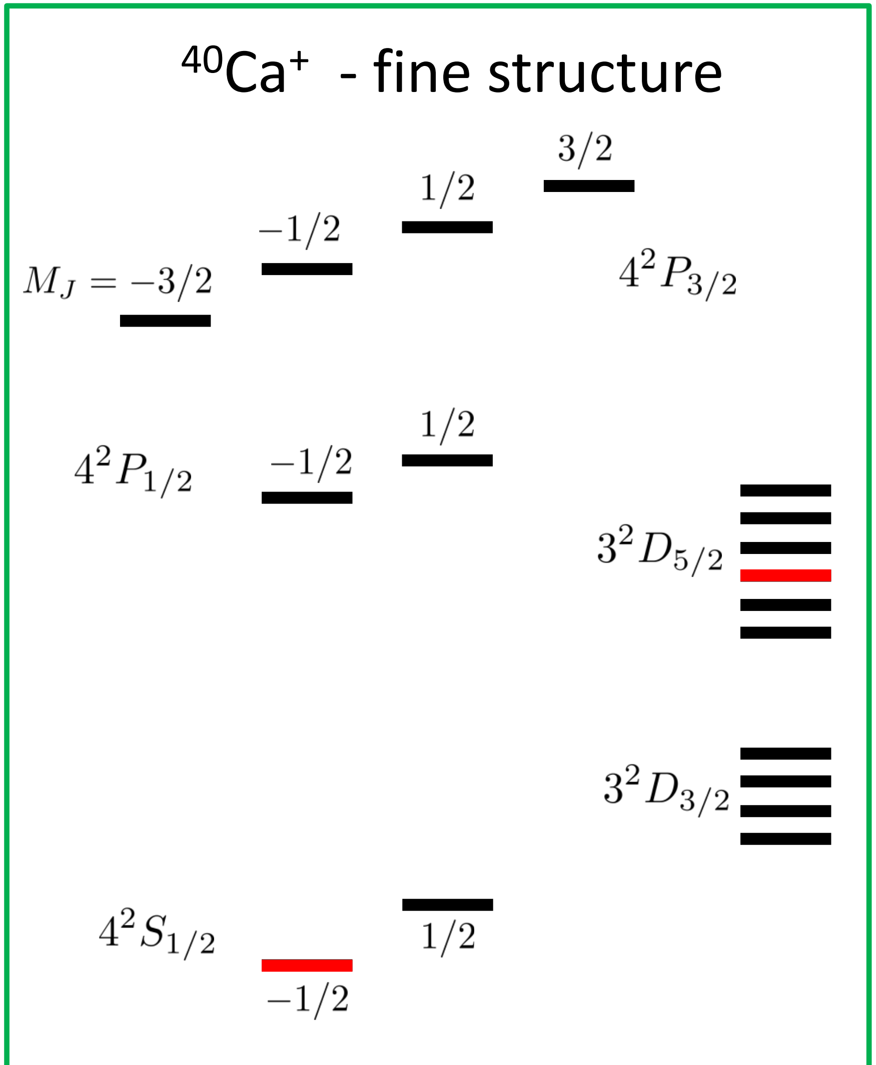
Ion traps at ETH Zurich

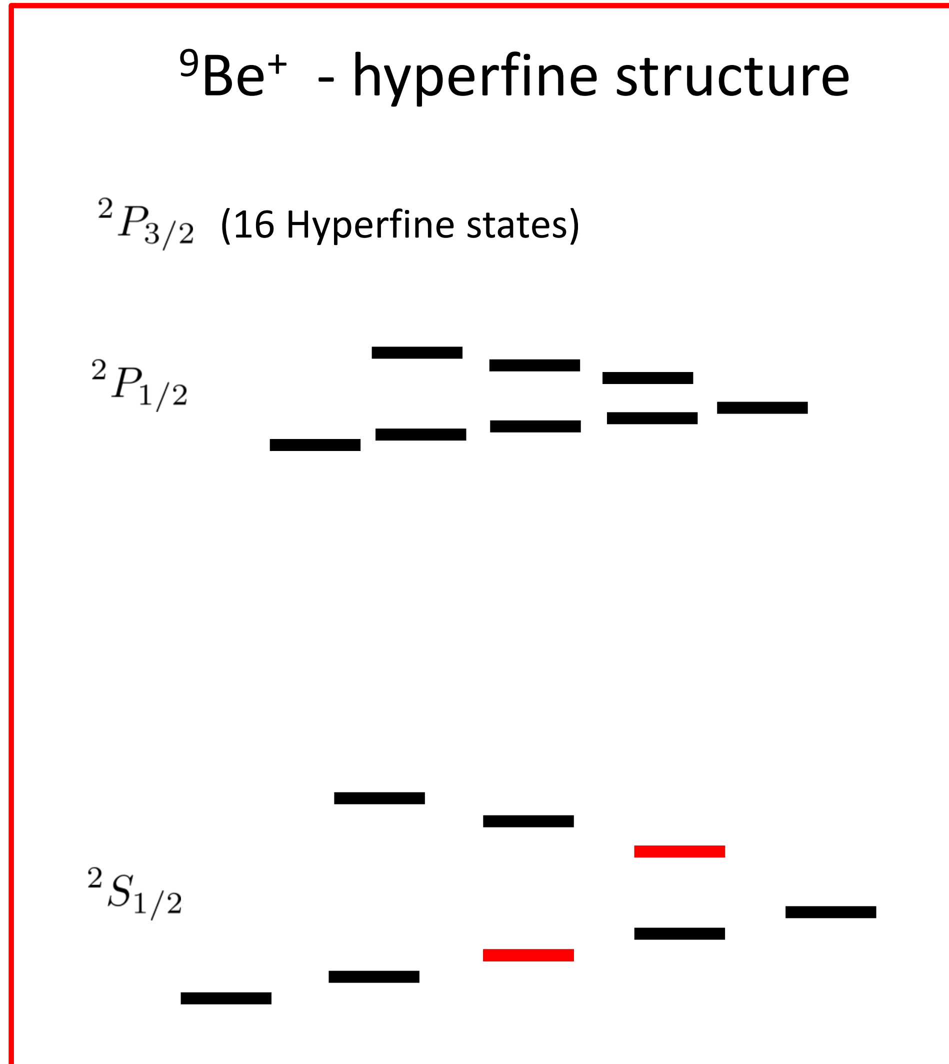




Storing a qubit in multi-level atoms

$$|\psi\rangle = (a|0\rangle + b|1\rangle)$$



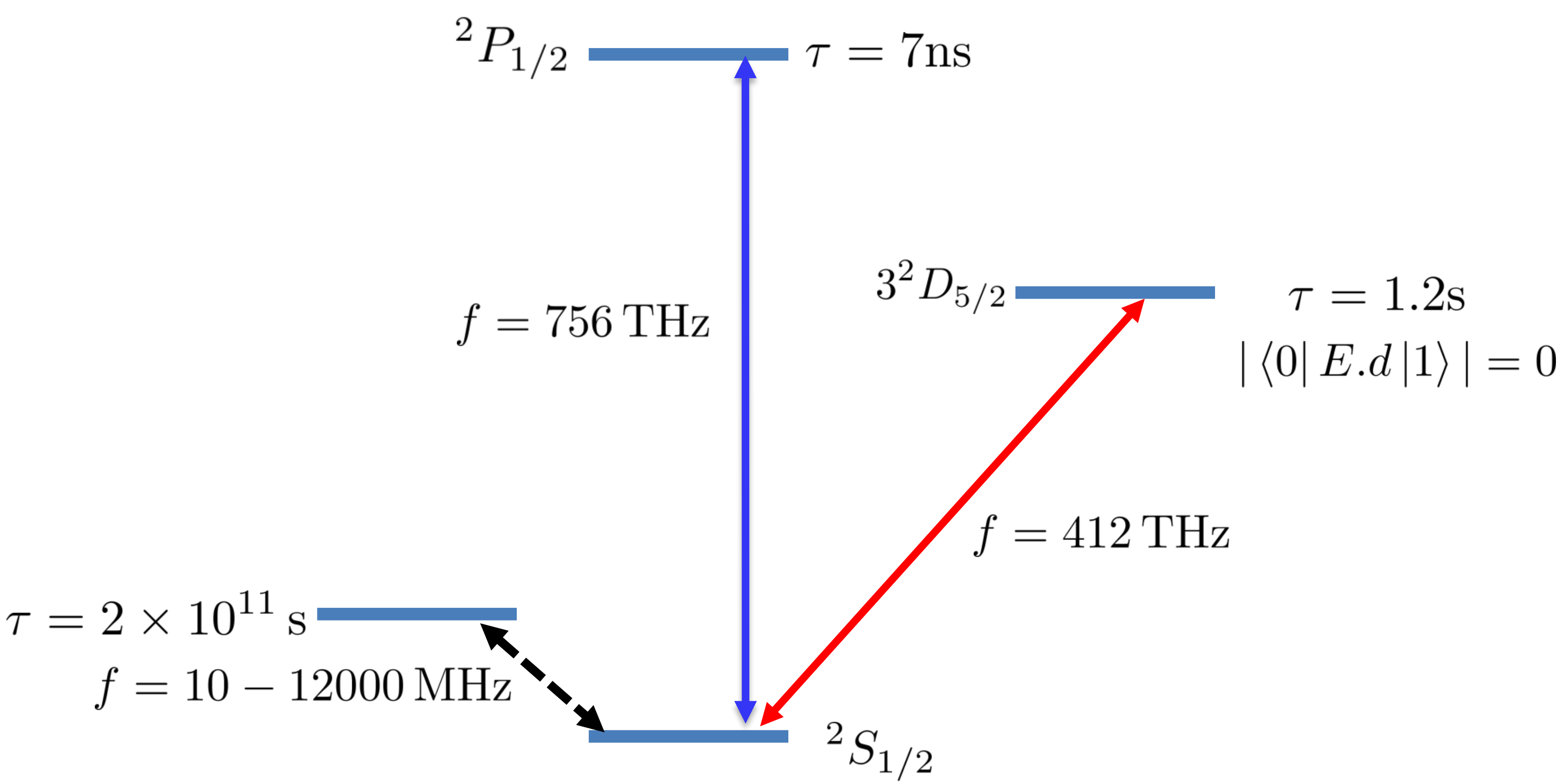


Choice of qubit levels

$$|\psi\rangle = (a|0\rangle + b|1\rangle)$$

Requirement: long decay time for upper level.

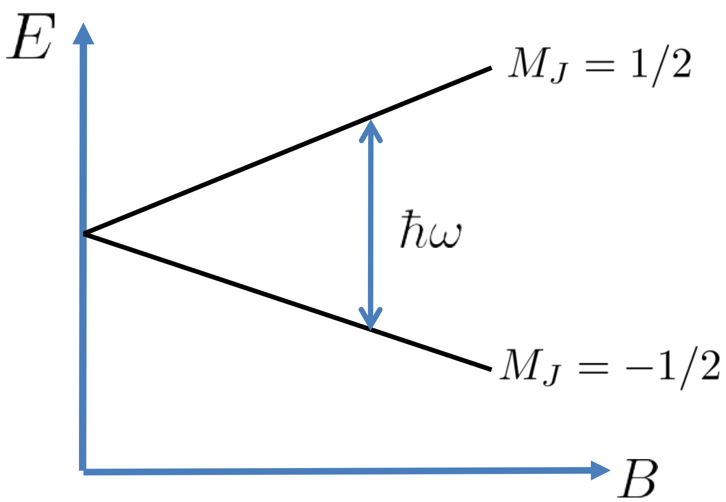
$$\tau \propto \frac{1}{\omega^3} \frac{1}{|\langle 0| E.d |1\rangle|^2}$$



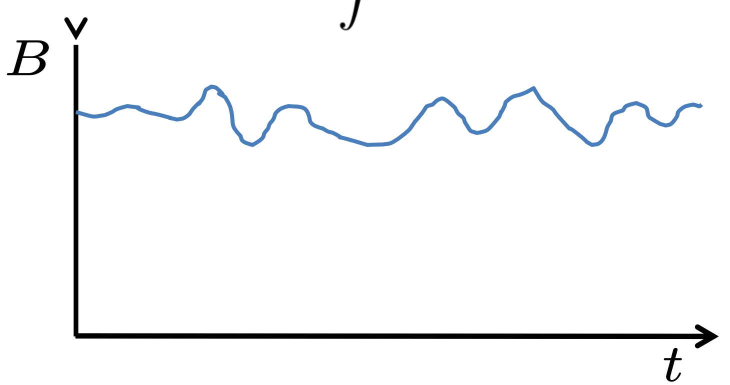
Storing qubits in an atom - phase coherence

$$|\psi\rangle = (a|0\rangle + be^{i\phi}|1\rangle)$$

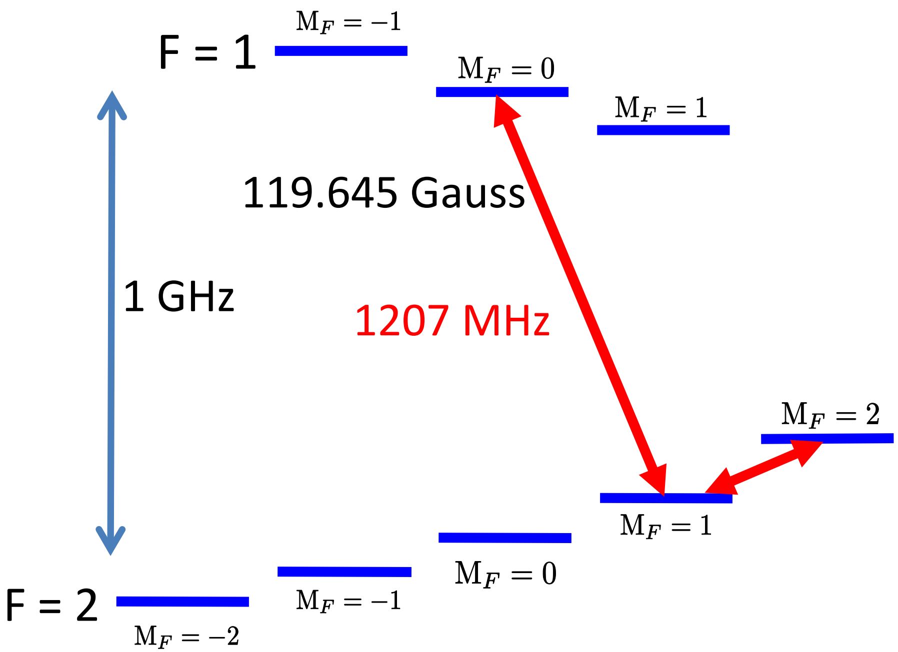
Problem: noise! – mainly from classical fields

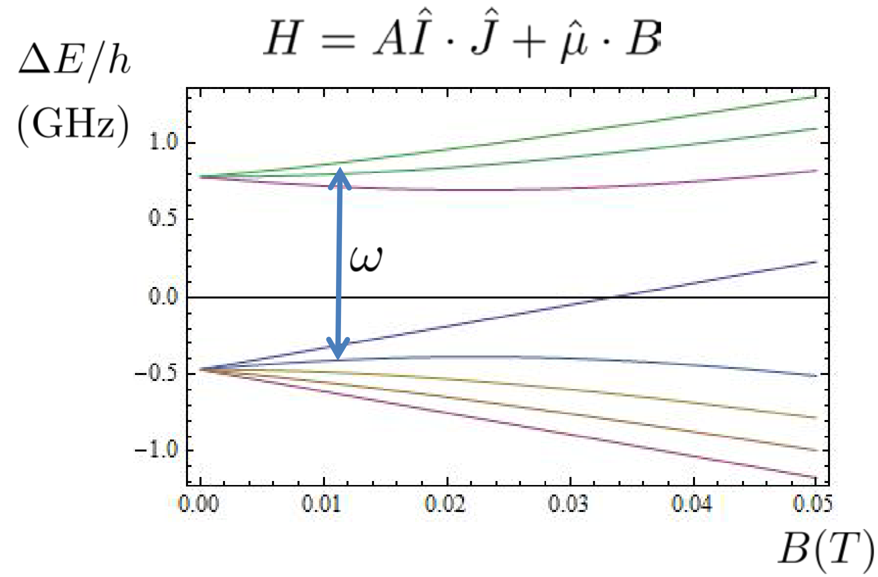


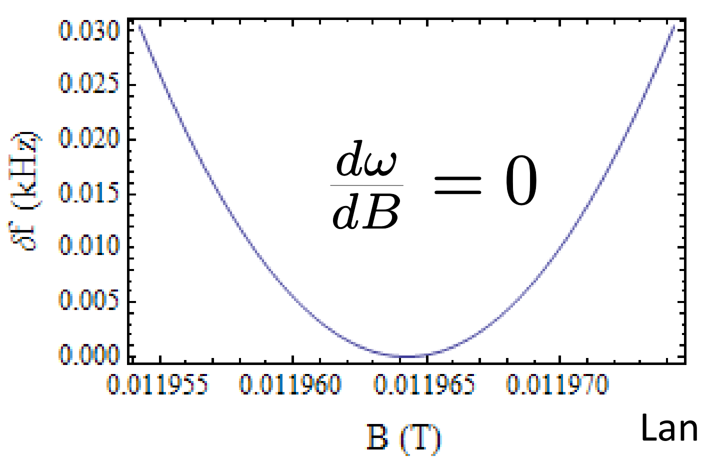
$$\phi \propto \int B(t)dt$$

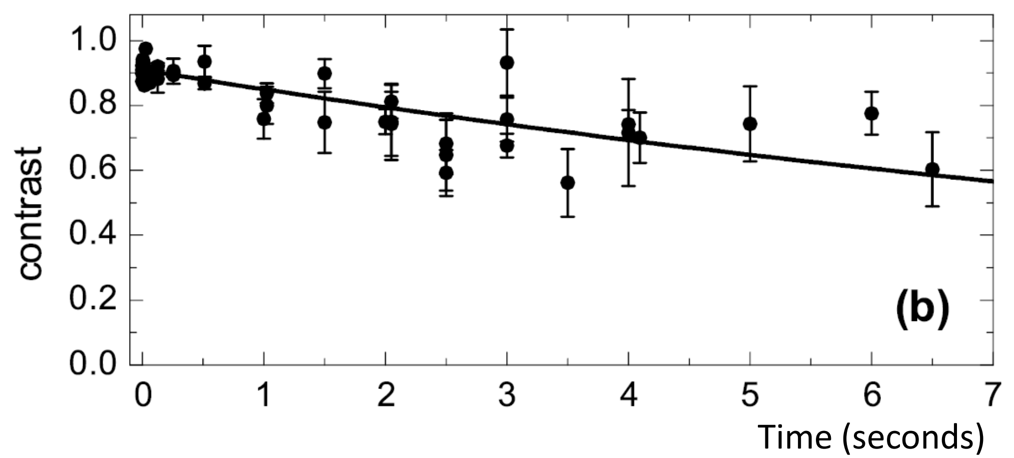


Storing qubits in an atom
Field-independent transitions









Langer et al. PRL 95, 060502 (2005), >50 seconds Alcock et al. (2012)

Entanglement for protection

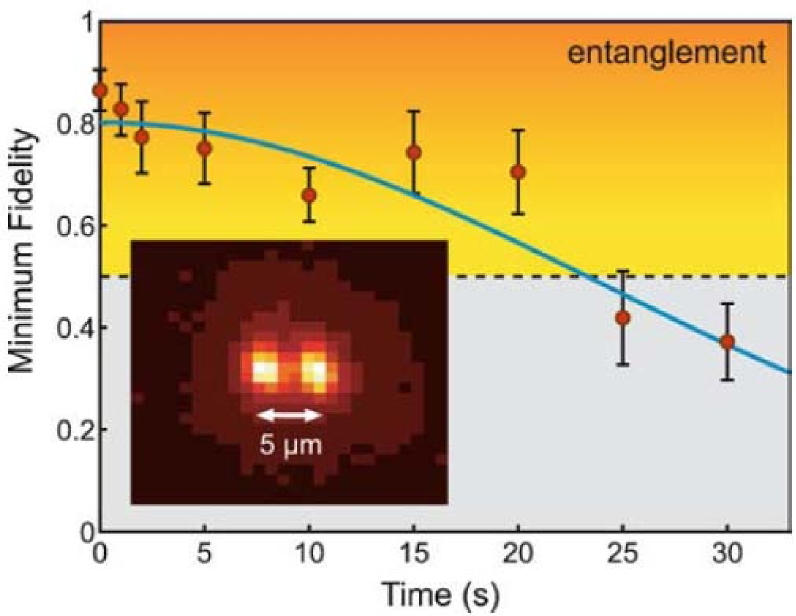
Rejection of common-mode noise

$$|0\rangle + e^{i\omega'(t)t}\,|1\rangle \qquad\qquad |0\rangle + e^{i\omega(t)t}\,|1\rangle$$

Now consider entangled state

$$e^{i\omega(t)t}\,|01\rangle + e^{i\omega'(t)t}\,|10\rangle = e^{i\omega(t)t}\left(|01\rangle + e^{i(\omega'(t)-\omega(t))t}\,|10\rangle\right)$$

If noise is common mode, entangled states can have very long coherence times



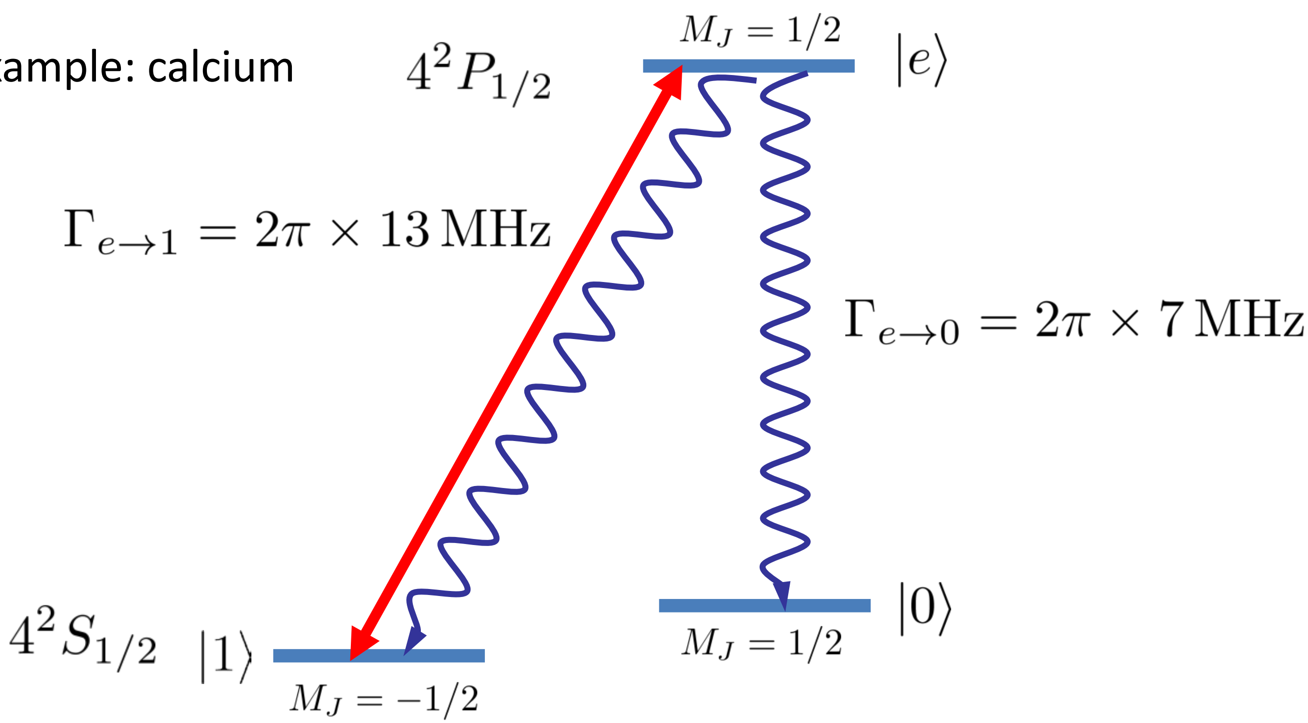
Haffner et al., Appl. Phys. B 81, 151-153 (2005)

Preparing the states of ions

Optical pumping – state initialisation

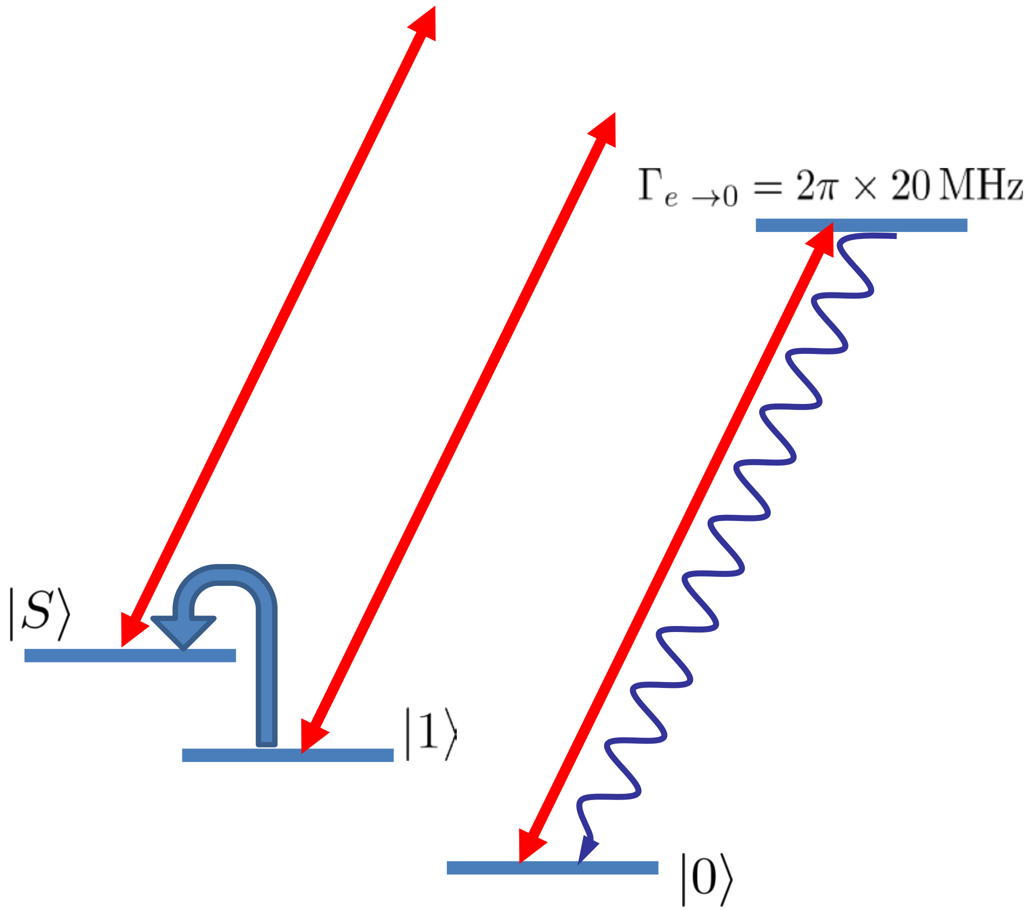
Use a dipole transition for speed

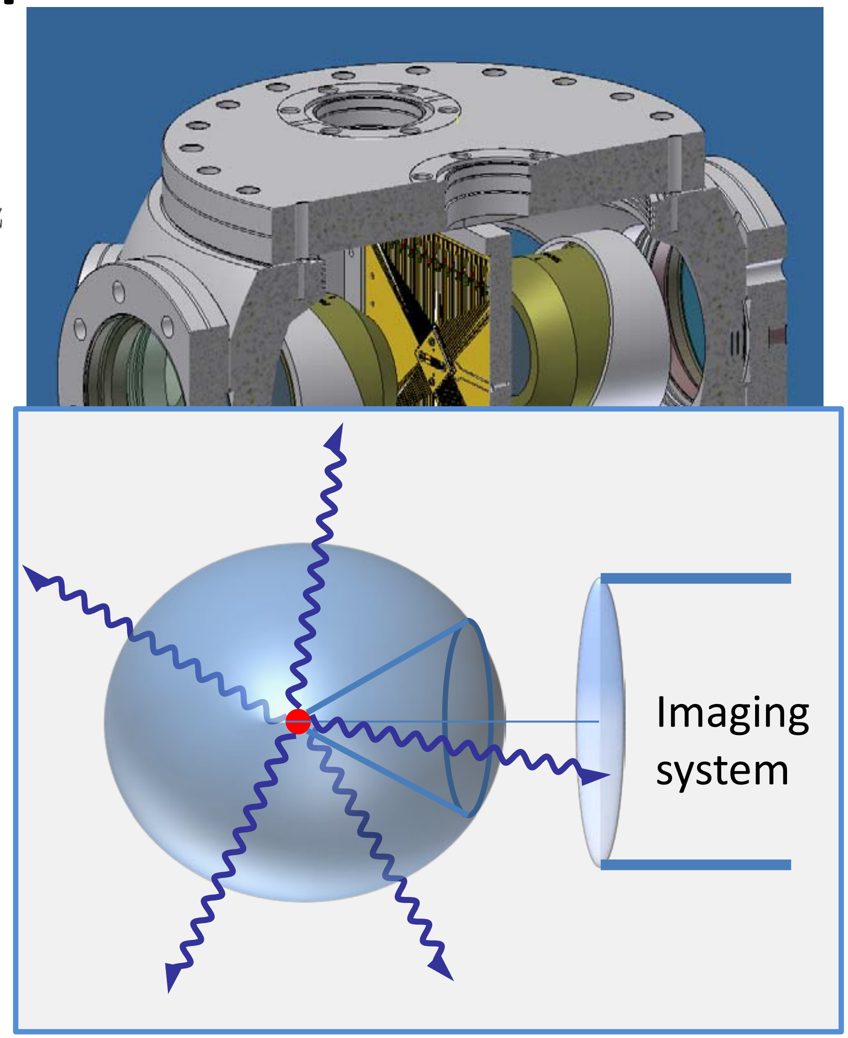
Example: calcium



Calcium: scatter around 3 photons to prepare $|0\rangle$ $\tau_{\mathrm{prep}} \sim 20\,\mathrm{ns}$

Reading out the quantum state





Photon scattered every 7 ns
BUT
we only collect a small fraction of these

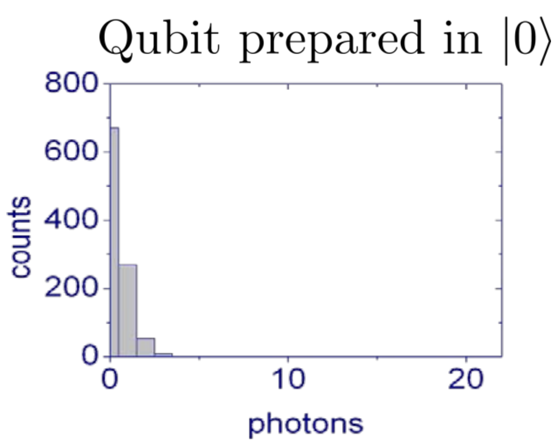
Need to scatter 1000 photons to detect atom $T_{\text{readout}} \sim 100 \to 1000\,\mu\text{s}$

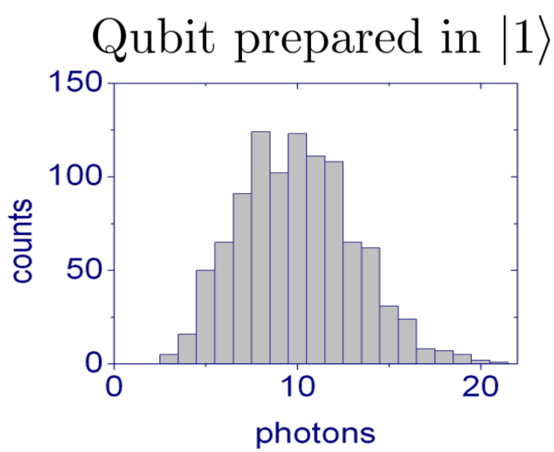
Measurement – experiment sequence

How many photons?

Statistics: repeat the experiment many (1000) times

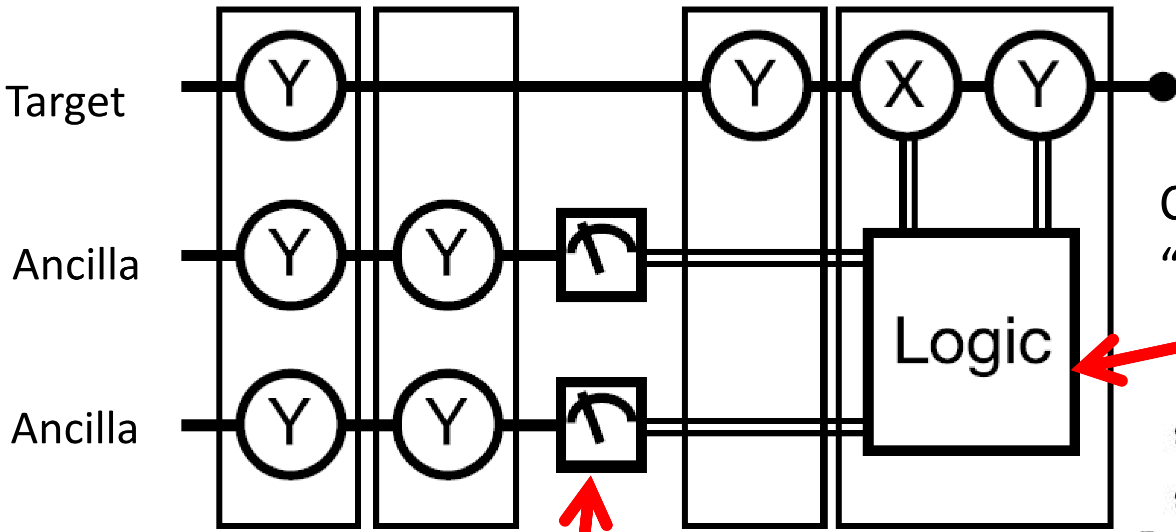
Number of photons = 8, 4, 2, 0, 0, 1, 5, 0, 0, 8 ....





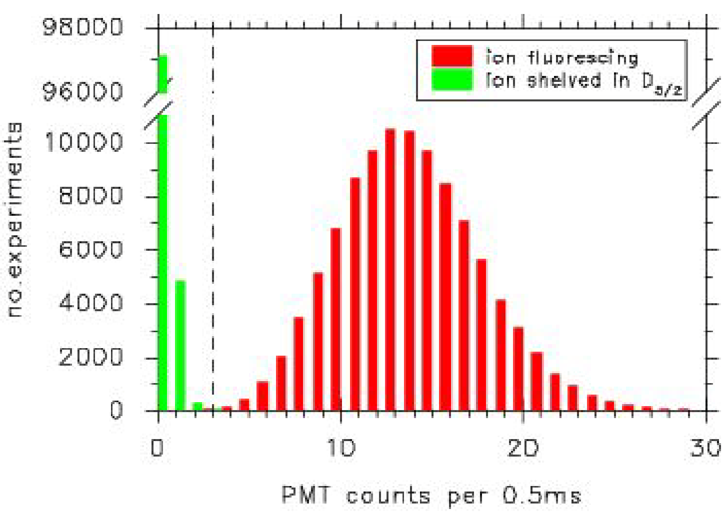
Single shot measurement

"Realization of quantum error-correction", Chiaverini et al., Nature 432, 602, (2004)



Classical processing
"If you get 1, 0, do Y, else do X"

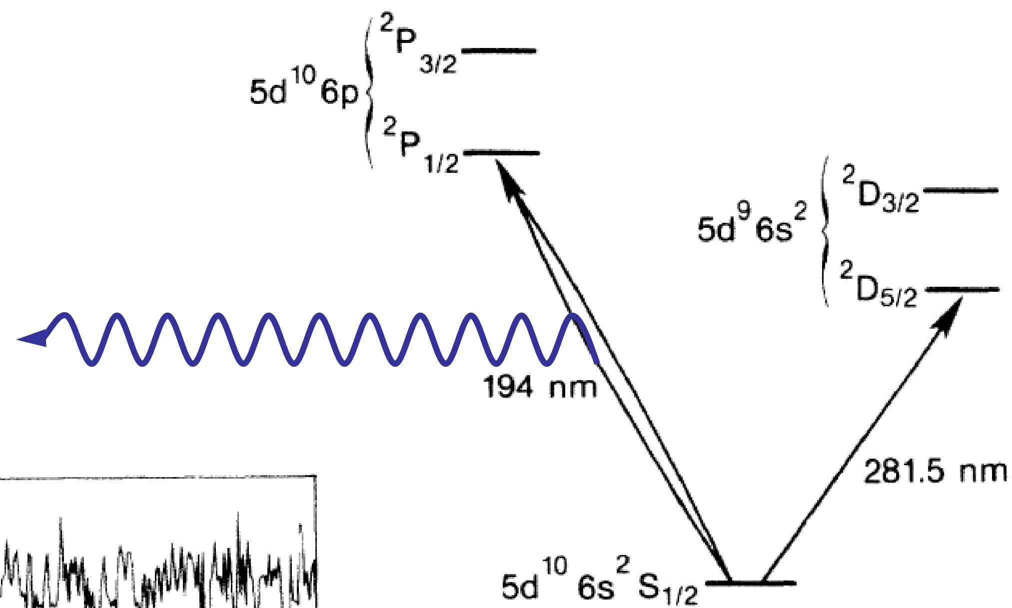
Measurement: "8 counts, this qubit is 1!"

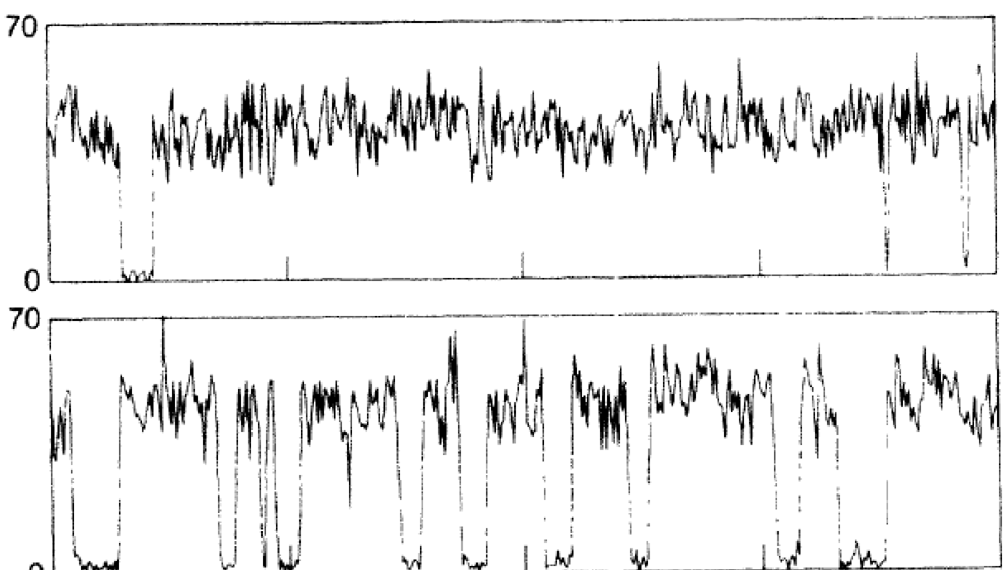


Accuracy of 0.9999 achieved in 150 microseconds
Myerson et al. Phys. Rev. Lett. 100, 200502, (2008)

Quantum jumps of a single atom

Bergquist et al. PRL 57, 14, (1986)



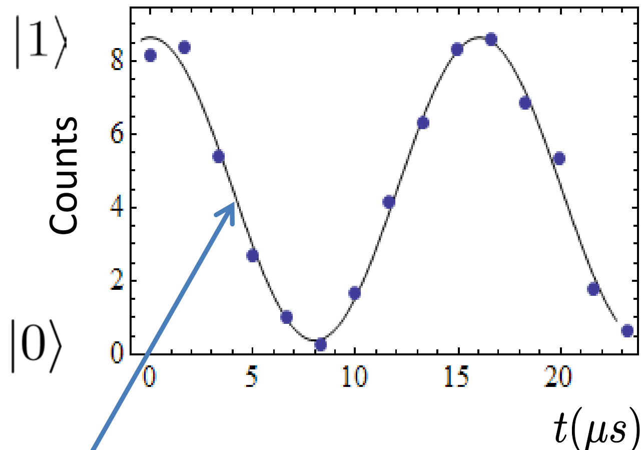


281.5 nm laser on resonance

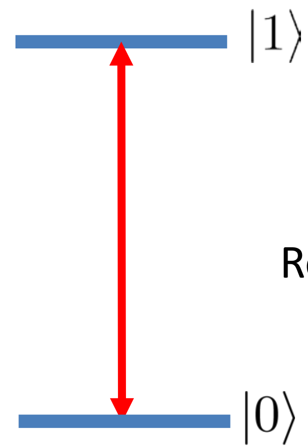
time

Manipulating single qubits

$$H = \Omega\left(|0\rangle\langle 1| + |1\rangle\langle 0|\right)\cos\left(\omega t + \phi\right)$$

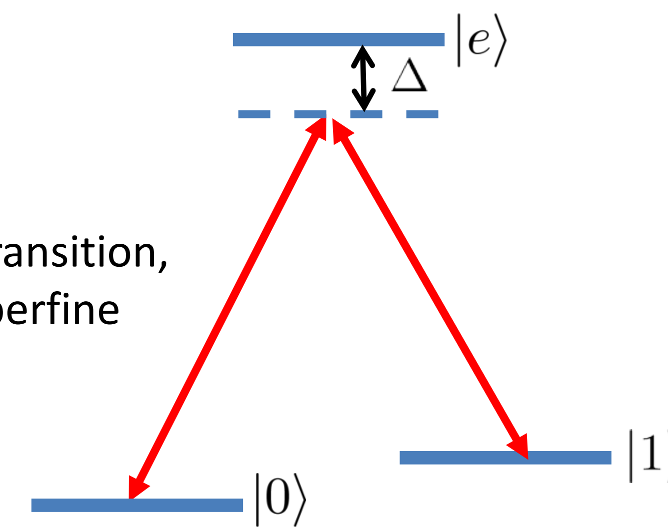


$$|\psi\rangle = (|0\rangle + e^{i\phi}|1\rangle)/\sqrt{2}$$



Laser-driven, quadrupole

Resonant microwaves, hyperfine



Raman transition, hyperfine

$$\text{Lowest error rate} \sim 1 \times 10^{-6}$$

Addressing individual qubits

Intensity addressing

Shine laser beam at one ion in string

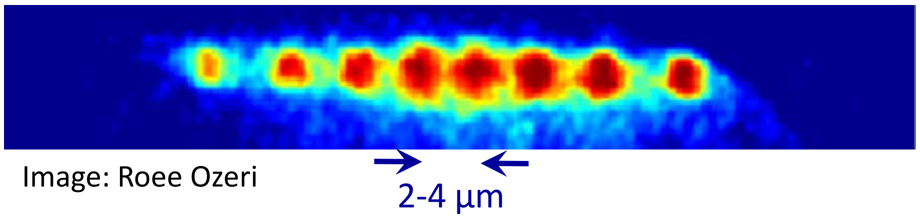
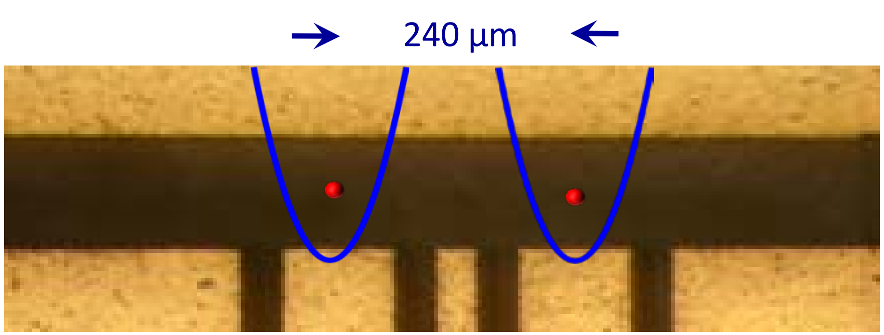
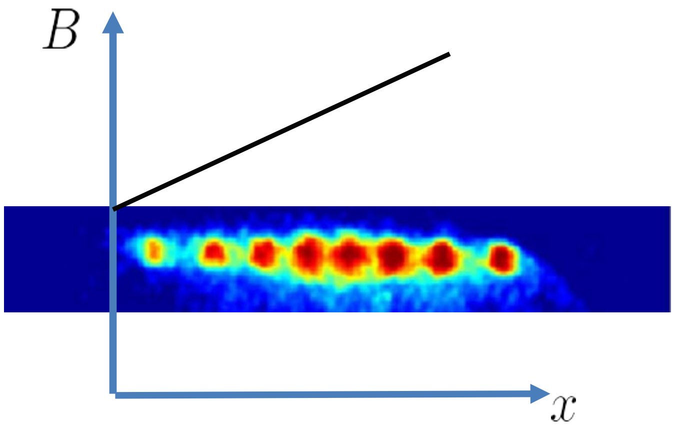


Image: Roee Ozeri

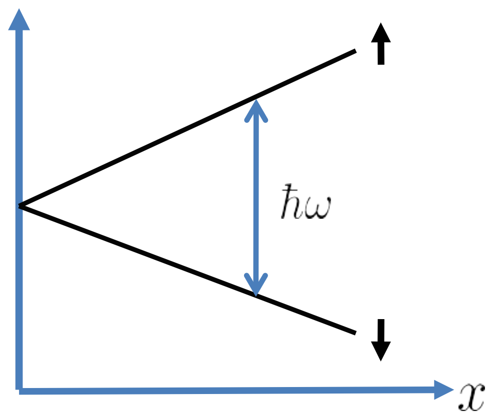
Separate ions by a distance much larger than laser beam size



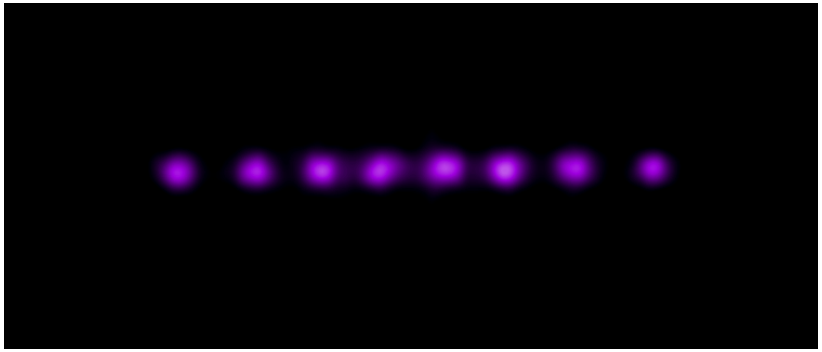
Frequency addressing



$E = -\mu \cdot B$

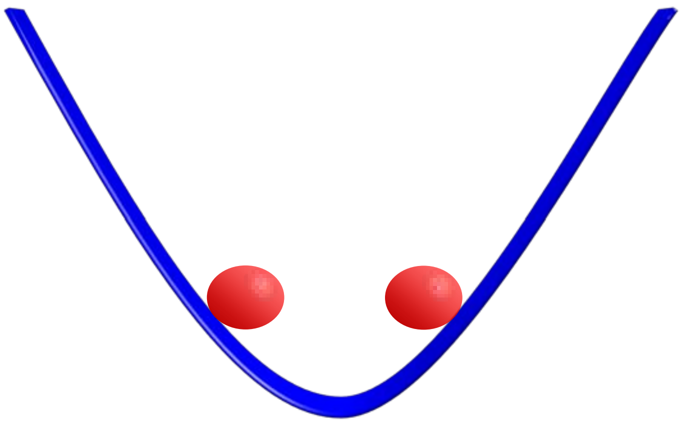


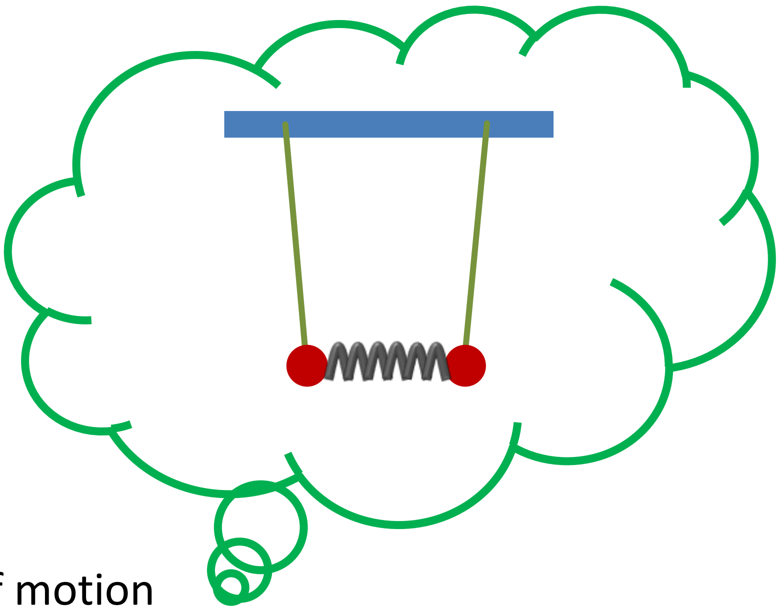
Multiple qubits: interactions



Multiple ions: coupled harmonic oscillators

$$V = \frac{k}{2}z_1^2 + \frac{k}{2}z_2^2 + \frac{\alpha}{|z_1 - z_2|}$$

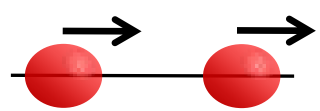




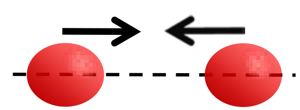
Expand about equilibrium – equation of motion

$$\begin{pmatrix} \ddot{\epsilon}_1 \\ \ddot{\epsilon}_2 \end{pmatrix} = -\omega_z^2 \begin{pmatrix} \epsilon_1 \\ \epsilon_2 \end{pmatrix} = \begin{pmatrix} k & \alpha \\ \alpha & k \end{pmatrix} \begin{pmatrix} \epsilon_1 \\ \epsilon_2 \end{pmatrix}$$

$\hbar\omega_c$



$\hbar\omega_s$

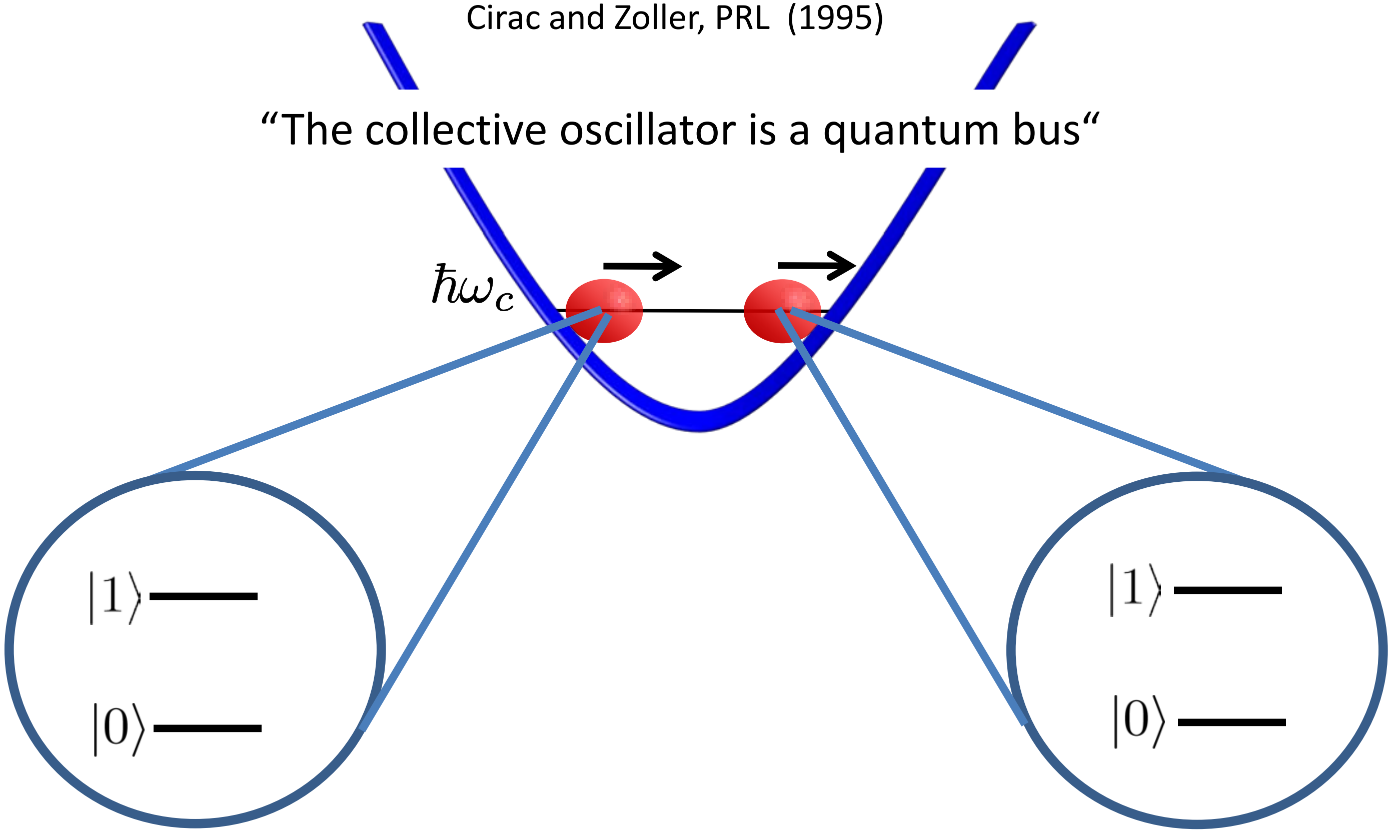


Independent oscillators
- shared motion

The original thought

Cirac and Zoller, PRL (1995)

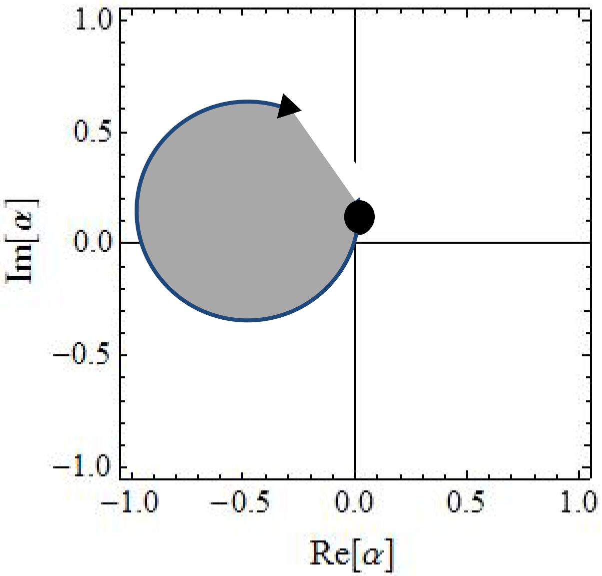
“The collective oscillator is a quantum bus“



$$G = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & i & 0 & 0 \\ 0 & 0 & i & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} |0_\phi 0_\phi\rangle \\ |1_\phi 0_\phi\rangle \\ |0_\phi 1_\phi\rangle \\ |1_\phi 1_\phi\rangle \end{pmatrix}$$

The forced harmonic oscillator

Classical forced oscillator $$\frac{d^2x}{dt^2} = -\omega_z^2 x + \frac{F}{m}\cos(\omega t + \phi)$$



“returns“ after $t = \frac{2\pi}{\delta}$

Radius of loop $\propto \frac{F}{\delta}$

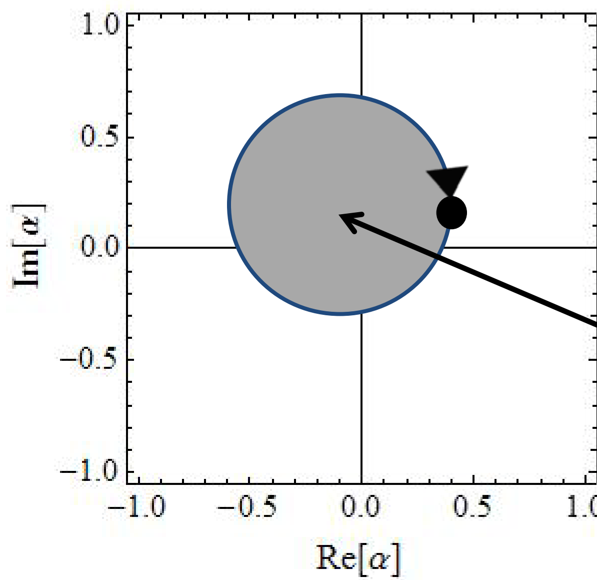
Forced quantum oscillators

$$\begin{aligned} H(t) &= \Omega\cos(\omega t)e^{ikz} \\ &\simeq \Omega\cos(\omega t)\left(1 + ikz_0\left(\hat{a}e^{i\omega_z t} + \hat{a}^\dagger e^{-i\omega_z t}\right)\right) \end{aligned}$$

$$[H(t), H(t')] \neq 0$$

$$U = \exp\left(\frac{i}{\hbar}\int^t H(t')dt' - \frac{1}{2\hbar^2}\int^t\int^{t'}[H(t'), H(t'')]dt'dt'' + \ldots\right)$$

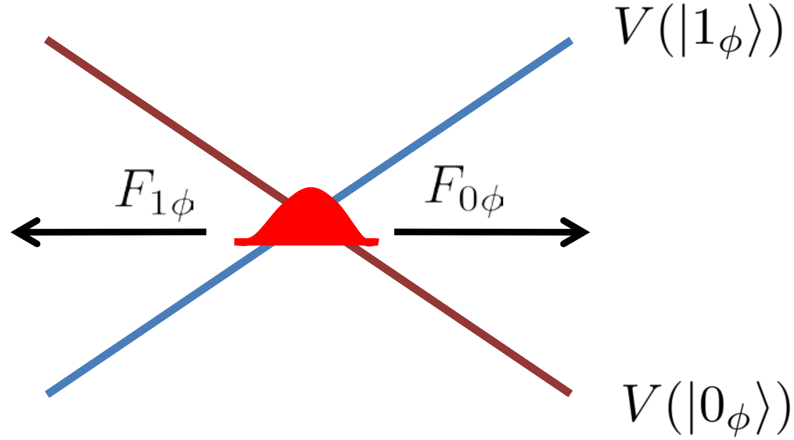
Transient excitation, phase acquired



$$|1\rangle|\alpha_0\rangle \to e^{i\Phi}|1\rangle|\alpha_0\rangle$$

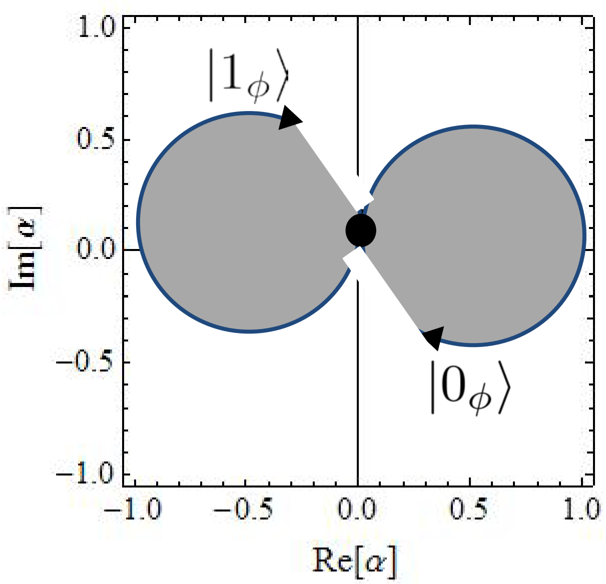
$$\Phi \propto A$$

State-dependent excitation

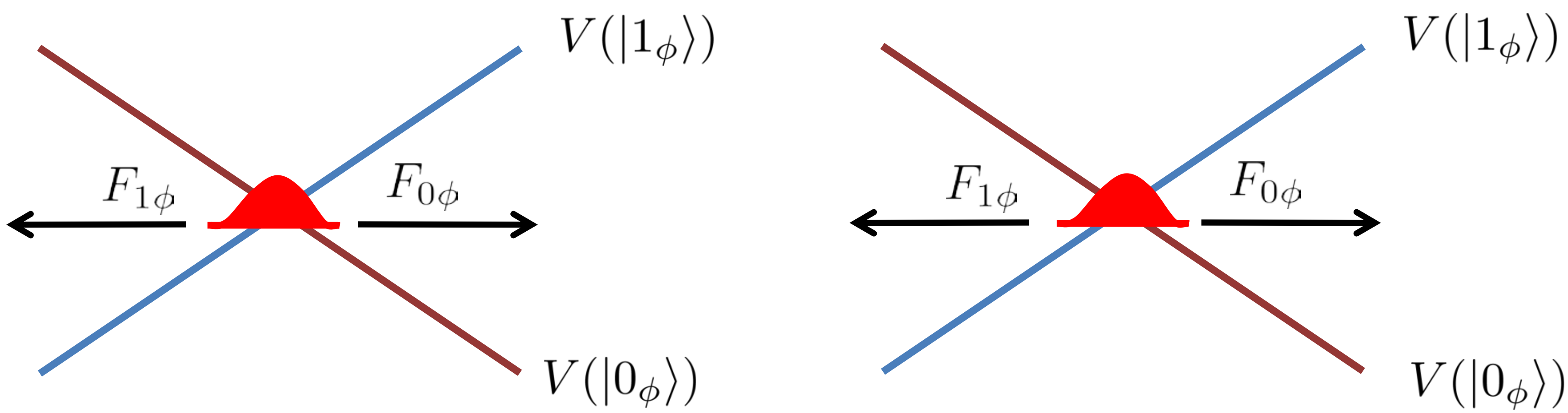


$$|\psi\rangle = |1_\phi\rangle\,|\alpha\rangle + |0_\phi\rangle\,|-\alpha\rangle$$

Entangled spin-motion state
“Schrodinger‘s cat“

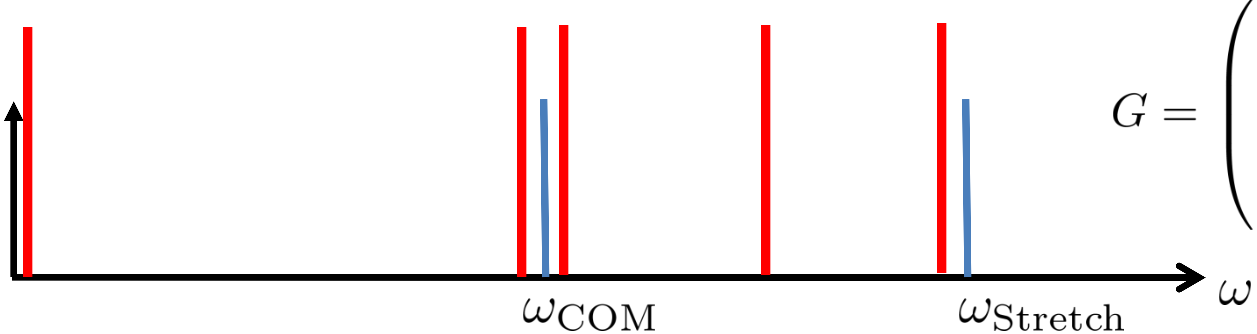


Two-qubit gate, state-dependent excitation



$|0_\phi 1_\phi\rangle\,,|1_\phi 0_\phi\rangle$ → Force is out of phase; excite Stretch mode

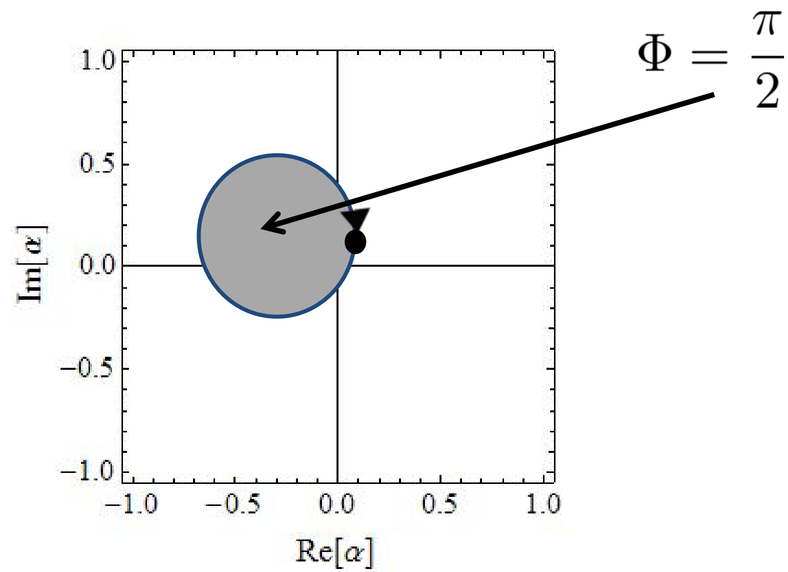
$|1_\phi 1_\phi\rangle\,,|0_\phi 0_\phi\rangle$ → Force is in-phase; excite COM mode



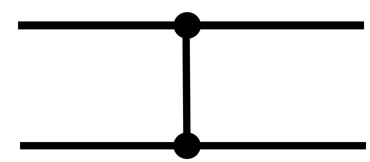
$$G = \begin{pmatrix} e^{i\phi_{00}} & 0 & 0 & 0 \\ 0 & e^{i\phi_{01}} & 0 & 0 \\ 0 & 0 & e^{i\phi_{01}} & 0 \\ 0 & 0 & 0 & e^{i\phi_{01}} \end{pmatrix}$$

Examples: quantum computing

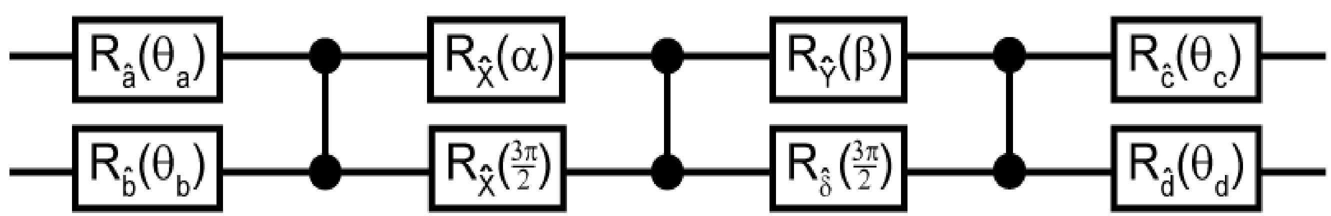
Choose the duration and power: $t_g = 2\pi/\delta \sim 7 \rightarrow 100\mu\text{s}$



$$G = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & i & 0 & 0 \\ 0 & 0 & i & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$



$\hat{U} =$

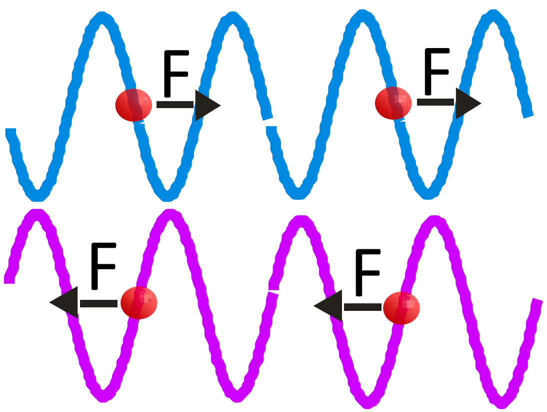


Universal two-qubit ion trap quantum processor: Hanneke et al. Nature Physics 6, 13-16 (2010)

Laser-driven multi-qubit gates

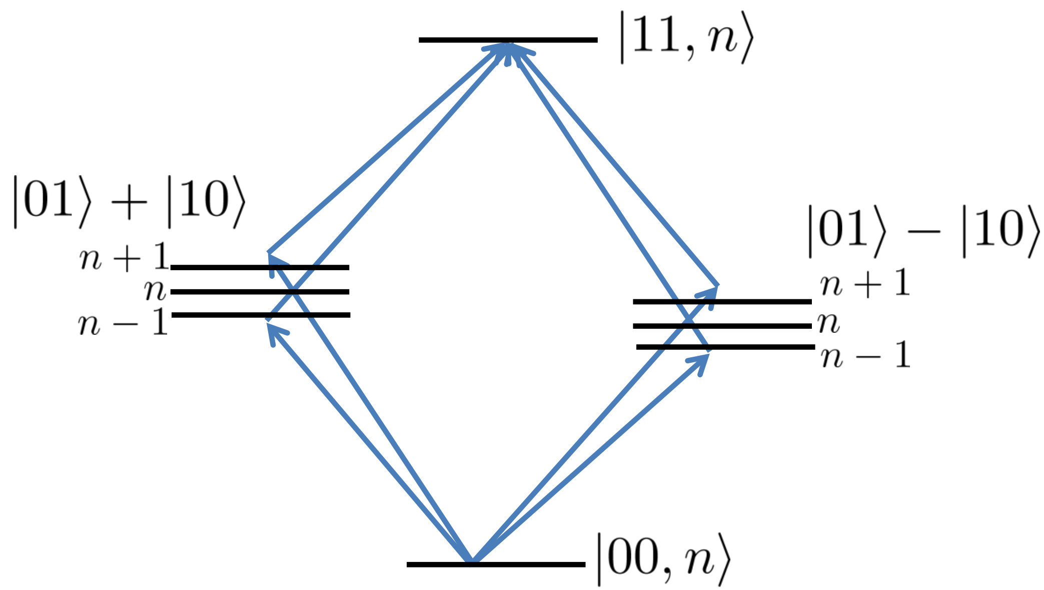
σ_z basis, polarisation standing wave

Leibfried et al. Nature 422, 412-415 (2003)

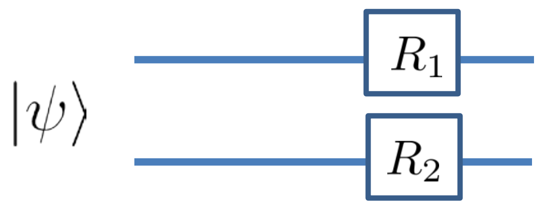


$$G = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & i & 0 & 0 \\ 0 & 0 & i & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

σ_x, σ_y basis, interference effect



State and entanglement characterisation



Detect

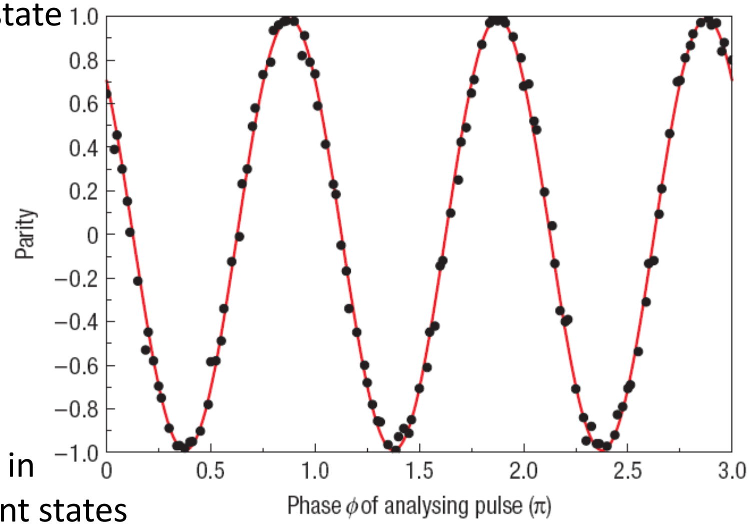
8, 6, 7, 4, 9, 0, 0, 1, 1, 6, 1, 9, 0, 0…
5, 4, 3,11, 4, 1, 0, 0, 1, 8, 0, 8, 1, 0…

Entanglement – correlations…

$$(|11\rangle + |00\rangle)\sqrt{2}$$

$$R_1 = R_2 = R(\theta = \pi/2, \phi)$$

Qubits in the same state



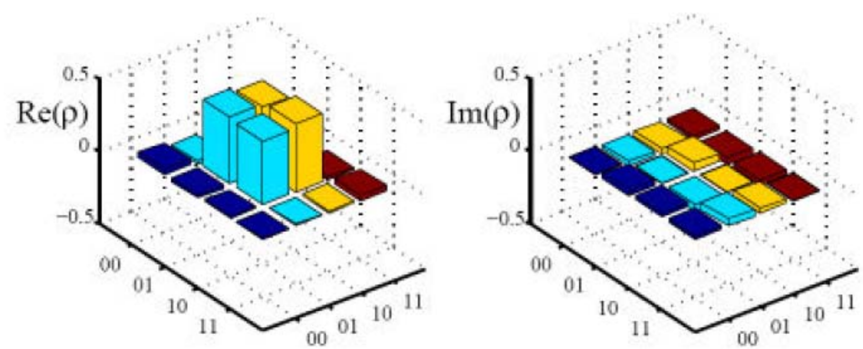
Qubits in different states

F = 0.993 (Innsbruck)

Benhelm et al. Nat. Phys 4, 463(2008)

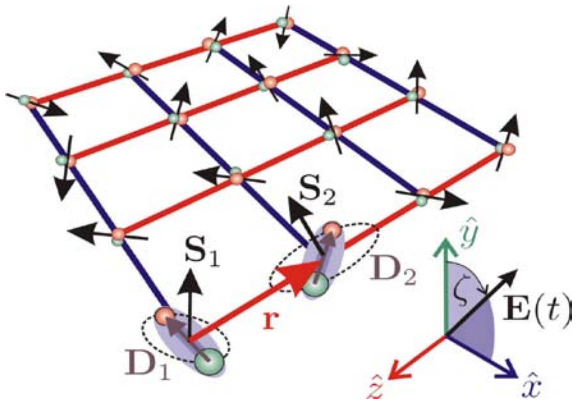
Choose 12 different settings of R_1, R_2

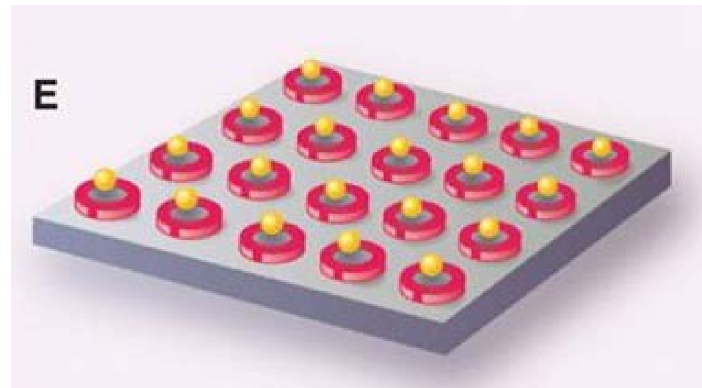
Reconstruct density matrix



Quantum simulation with trapped-ions

Creation of “condensed-matter“ Hamiltonians with long-range interactions
(Friedenauer et al. Nat. Phys 4, 757-761 (2008), Kim et al. Nature 465, 7298 (2010))





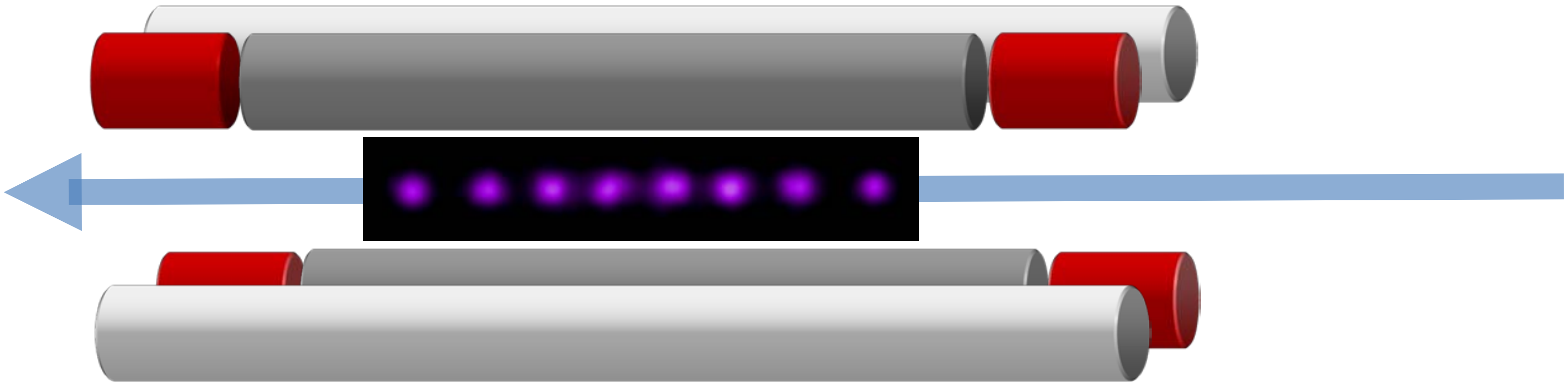
Go to limit of large motional detuning
(very little entanglement between spin and motion)

$$\Omega \ll \delta$$

$$\Phi_{10} = \Phi_{01} \simeq i\frac{\Omega^2}{\delta}t$$

$$H_{\text{eff}} \simeq \frac{\Omega^2}{\delta} s_1^x s_2^x \qquad H_{\text{eff}} \simeq \frac{\Omega^2}{\delta} \sum_{i \neq j}^{N} s_i^x s_j^x$$

Dealing with large numbers of ions

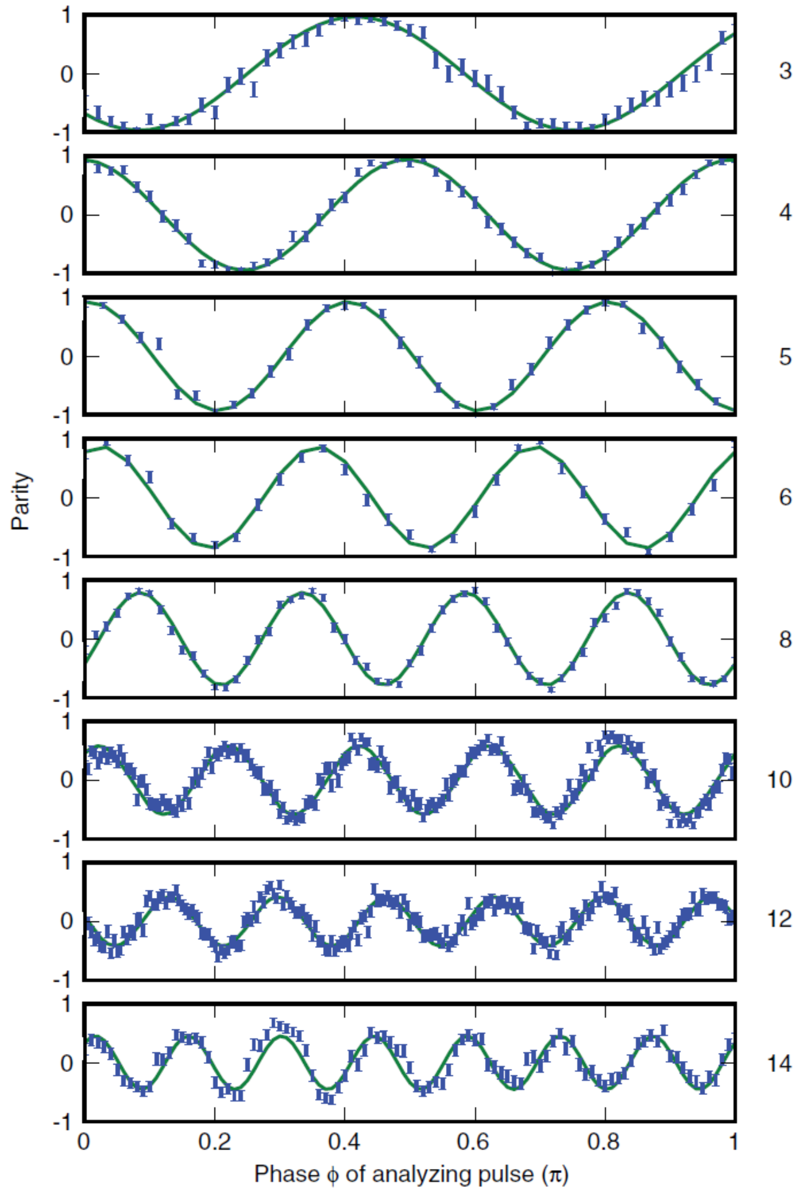


Technical requirement	Limitation
Spectral mode addressing	Mode density increases
Many ions	Heating rates proportional to N
Simultaneous laser addressing	Ions take up space (separation > 2 micron) Laser beams are finite-size

Entanglement of up to 14 ions

$$\left(|11...1\rangle + |00...0\rangle\right)/\sqrt{2}$$

Monz et al., PRL 106, 130506 (2011)

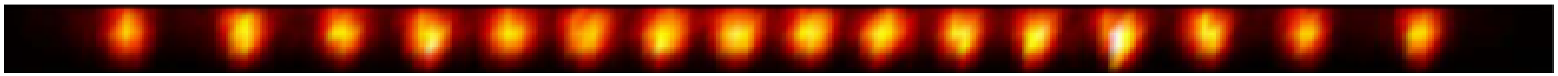


High contrast – 3 ions

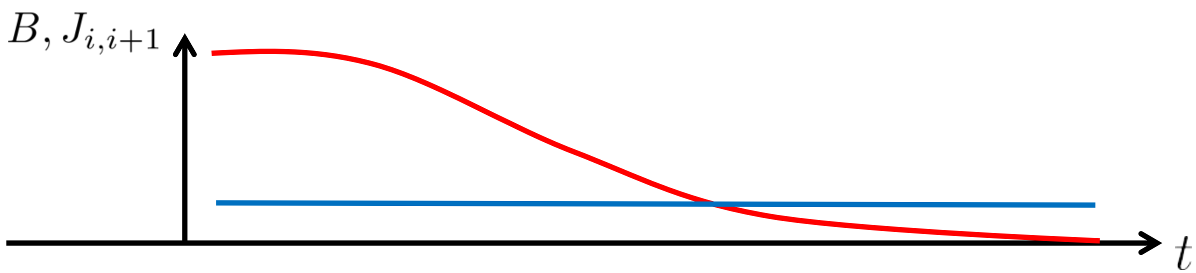
Reduced contrast – 14 ions

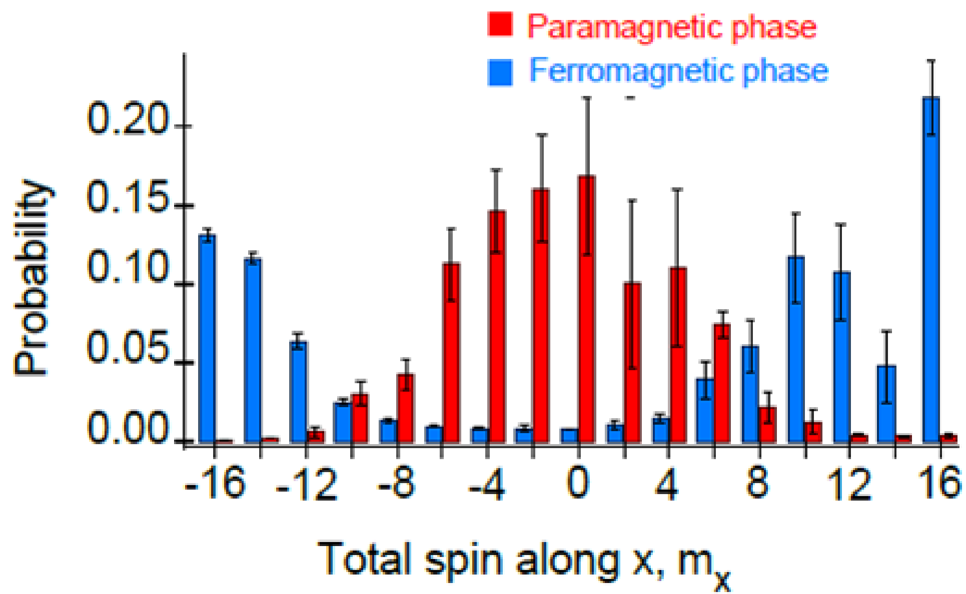
Quantum simulations in ion strings

(up to 18 ions Islam et al. arxiv:1210.0142)



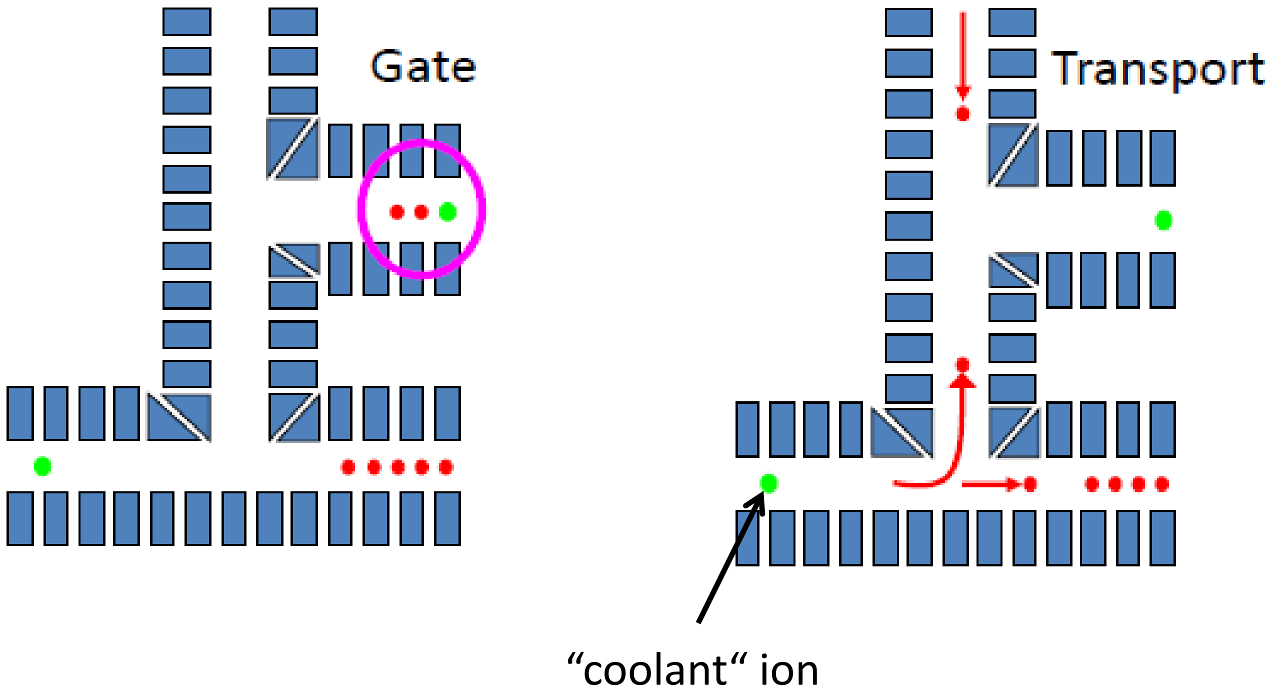
$$\hat{H} = \sum_{j<i} J_{ij}\sigma_x^i\sigma_y^j - B\sum_i \sigma_y^i$$





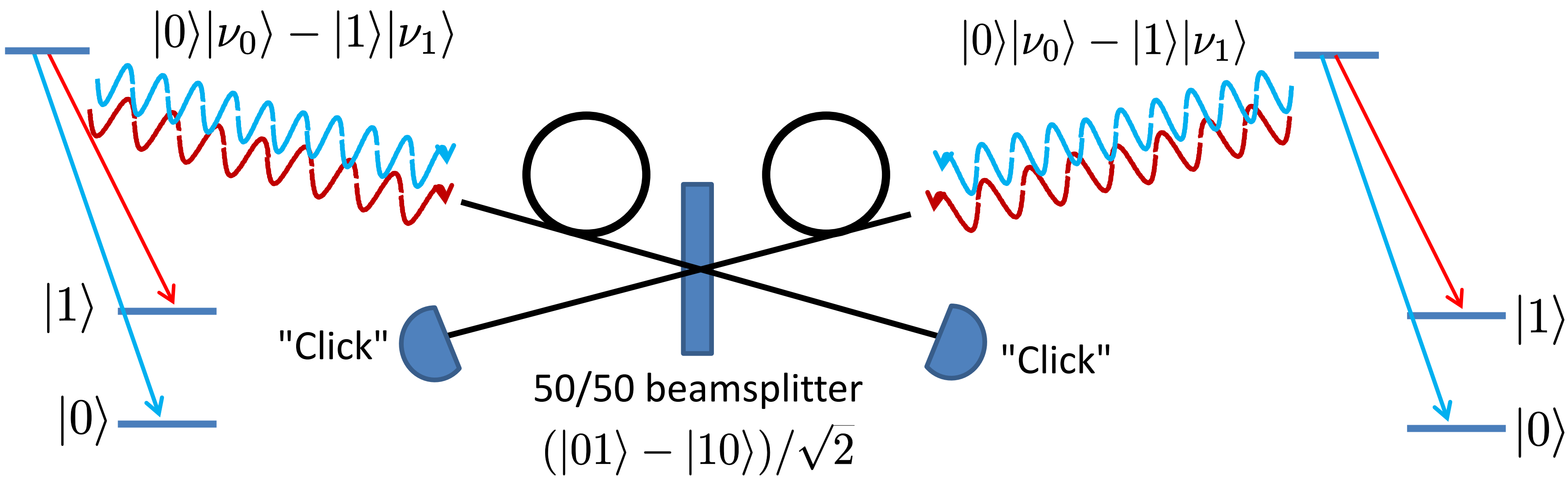
Isolate small numbers of ions

Wineland et al. J. Res. Nat. Inst. St. Tech, (1998)

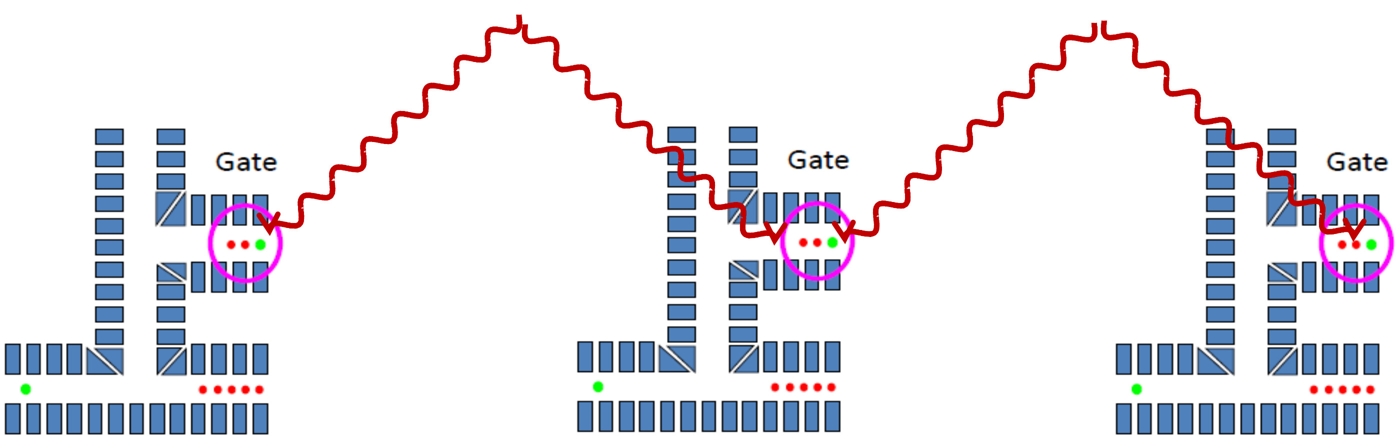


Technological challenge – large numbers of electrodes, many control regions

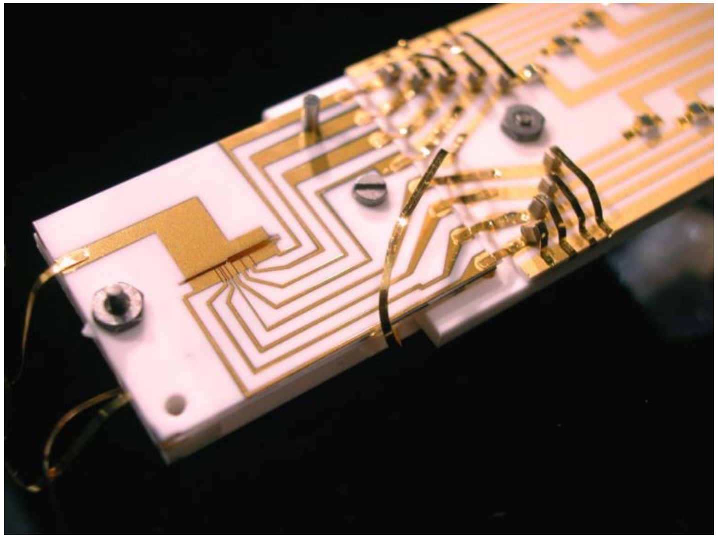
Distributing entanglement: probabilistic

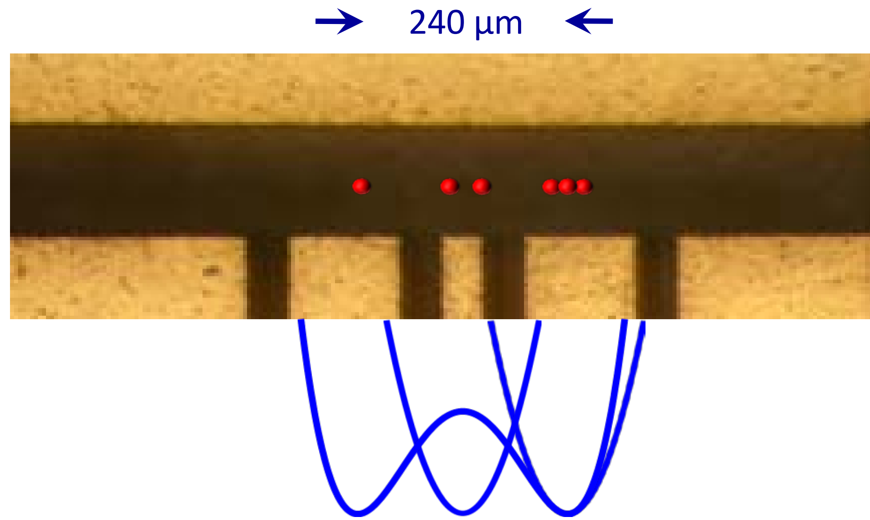


Entangled ions separated by **1m** (Moehring et al. Nature 449, 68 (2008))



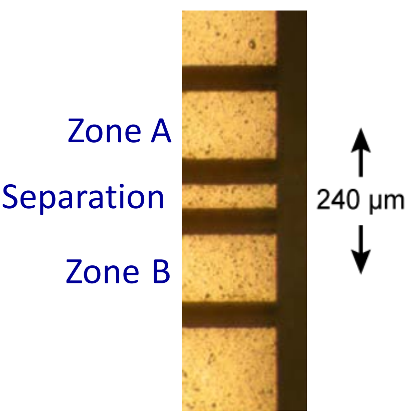
Transport with ions



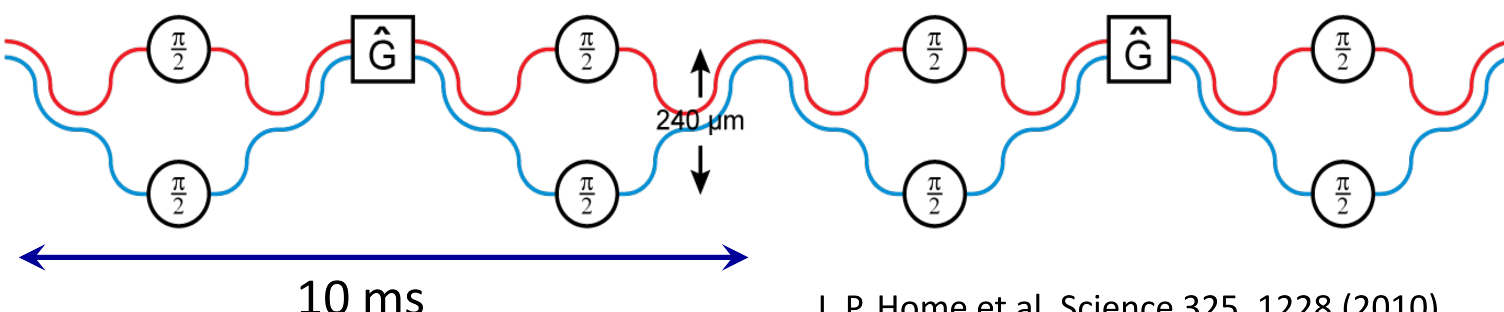


Move: 20 us, Separate 340 us, 0.5 quanta/separation

Internal quantum states of ions unaffected by transport
Motional states are affected – can be re-initialised



Total transport distance = 1 mm

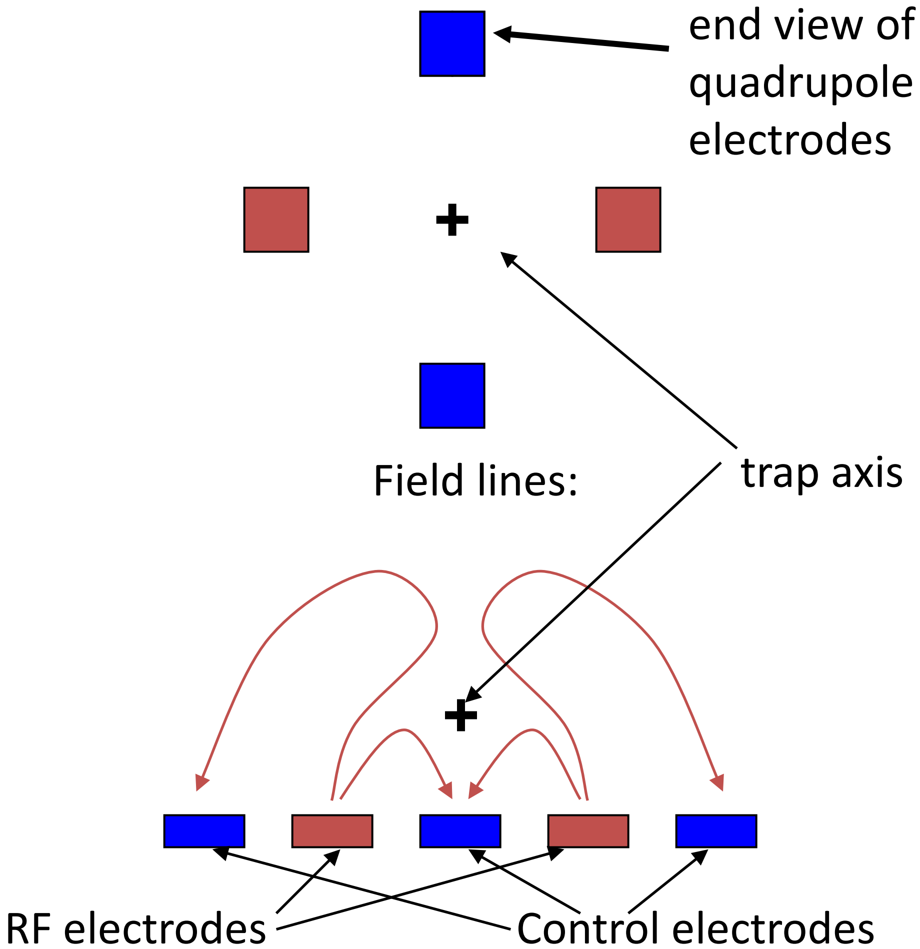


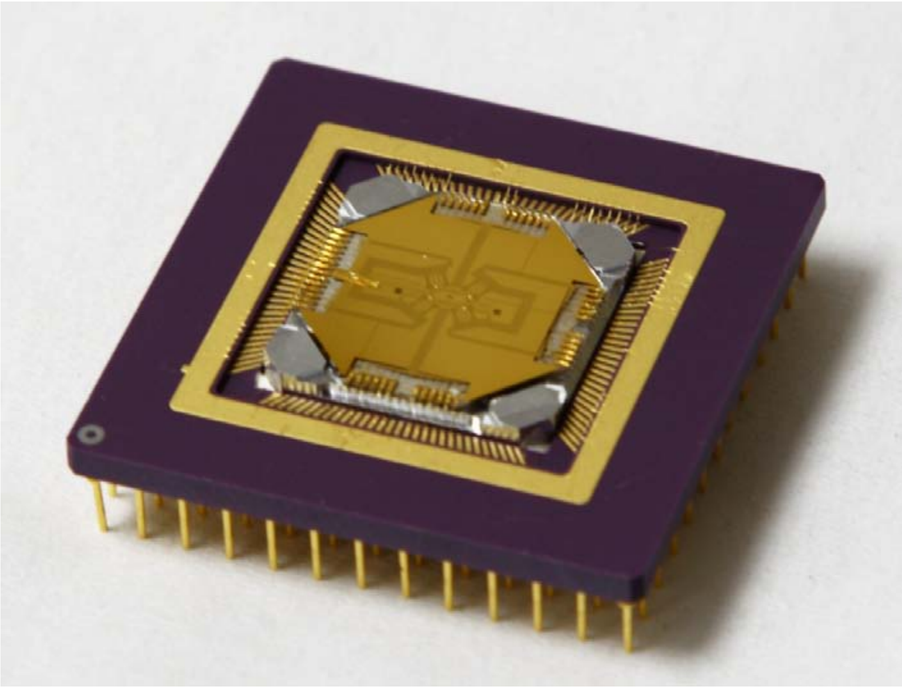
J .P. Home et al. Science 325, 1228 (2010)

Trapping ions on a chip

For microfabrication purposes, desirable to deposit trap structures on a surface

(Chiaverini *et al.*, Quant. Inf. & Computation (2005), Seidelin et al. PRL 96, 253003 (2006))

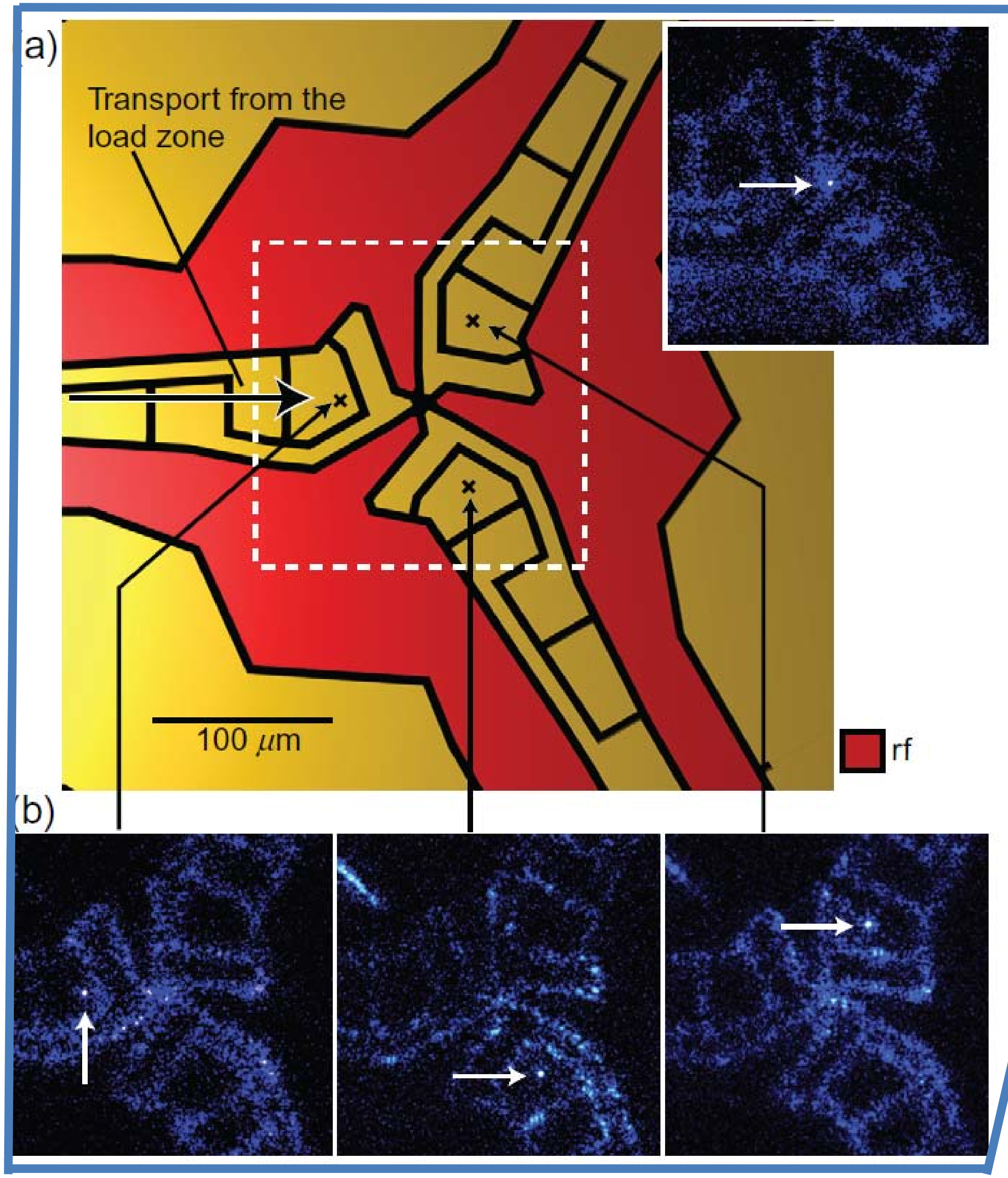


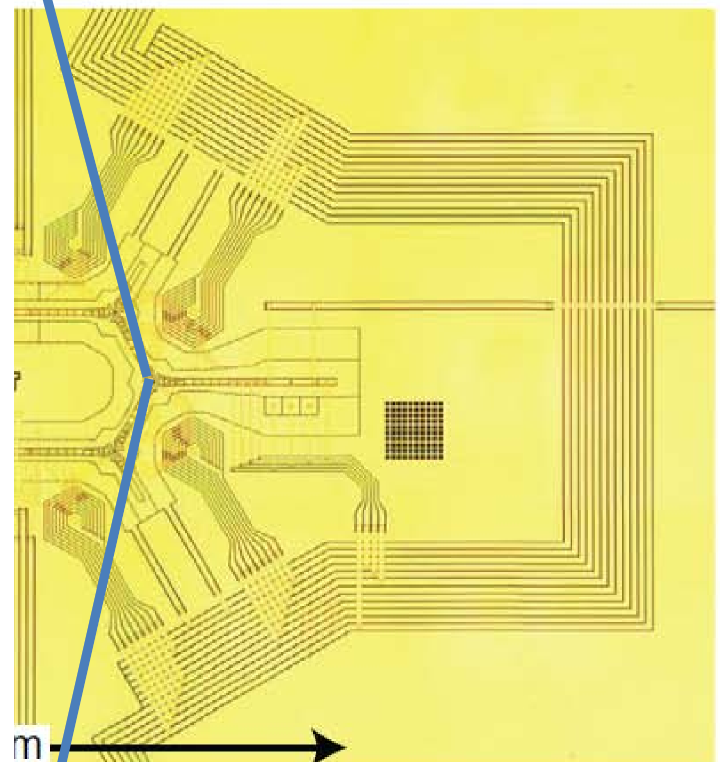


Challenges: shallow trap depth (100 meV)
charging of electrodes

Opportunities: high gradients

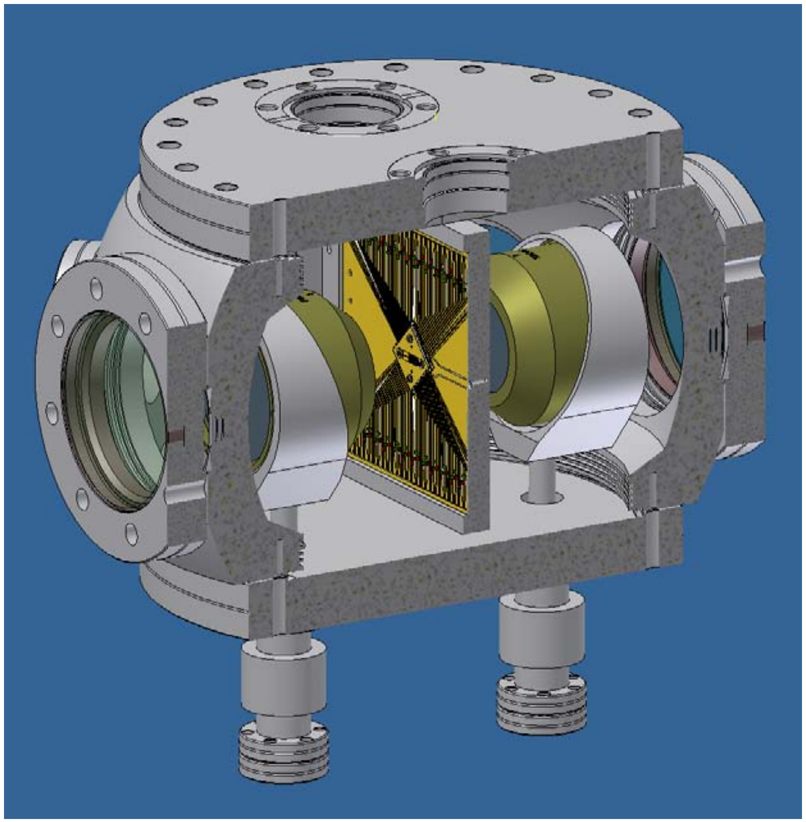
omplicated) chip

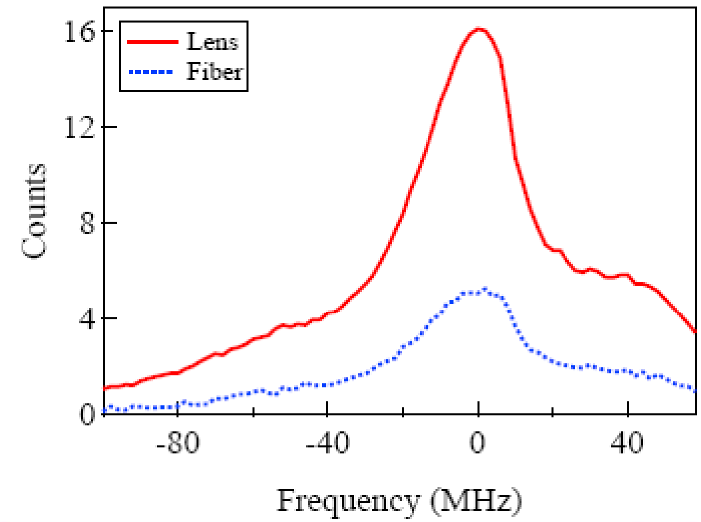


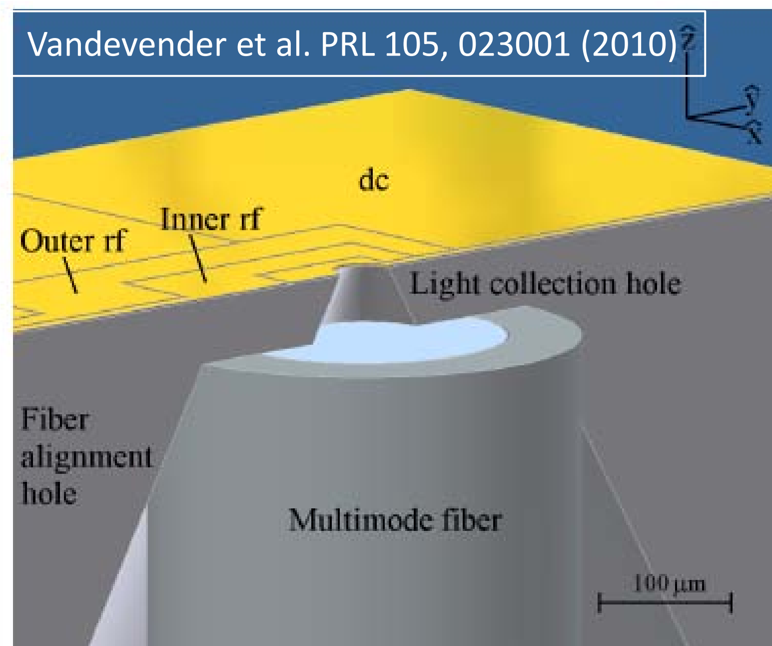


J. Amini et al. New. J. Phys 12, 033031 (2010)

Integrated components 1



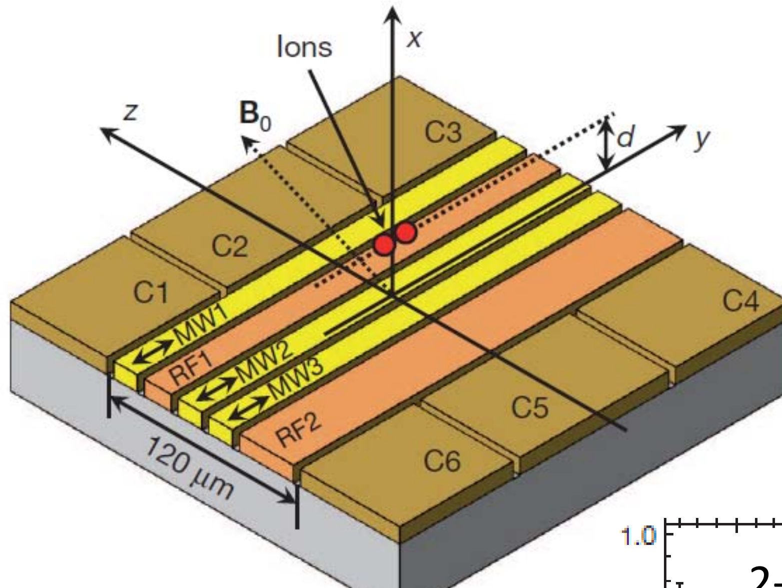




Vandevender et al. PRL 105, 023001 (2010)

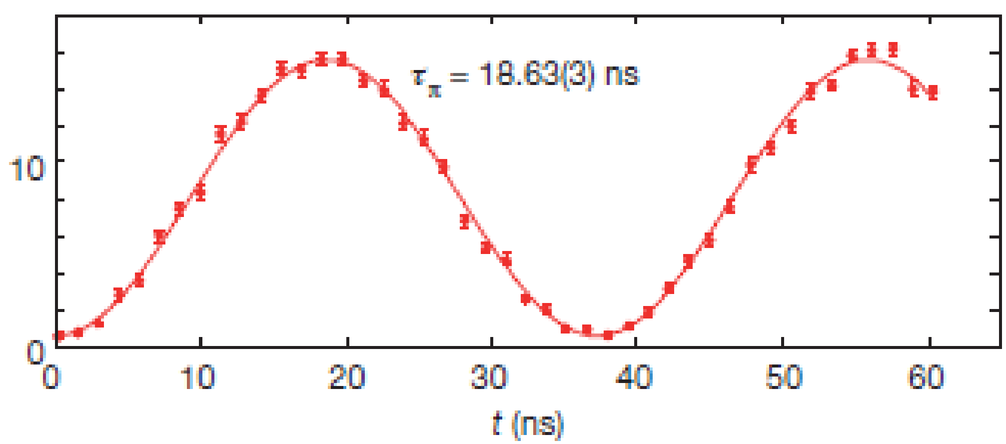
Integrated components

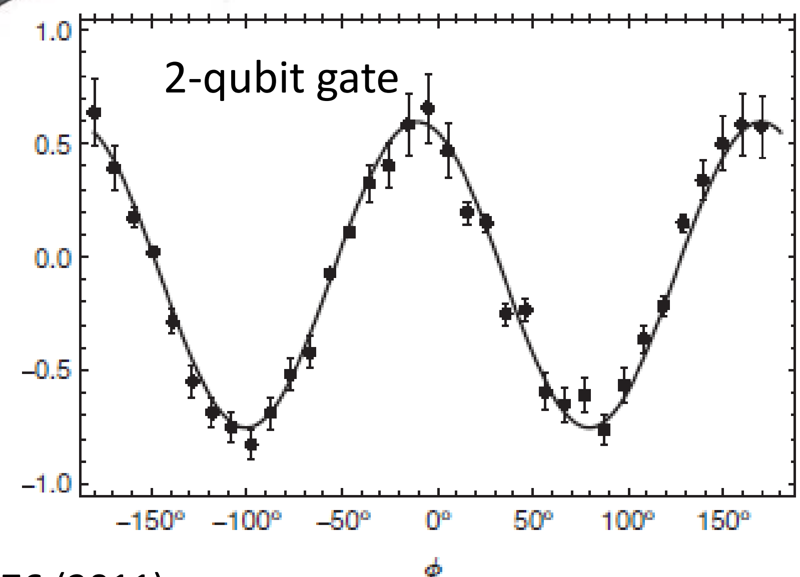
eg. Quantum control using microwaves – removes the need for high-power lasers



Gradients – produce state-dependent potentials through Zeeman shifts

Single-qubit gate



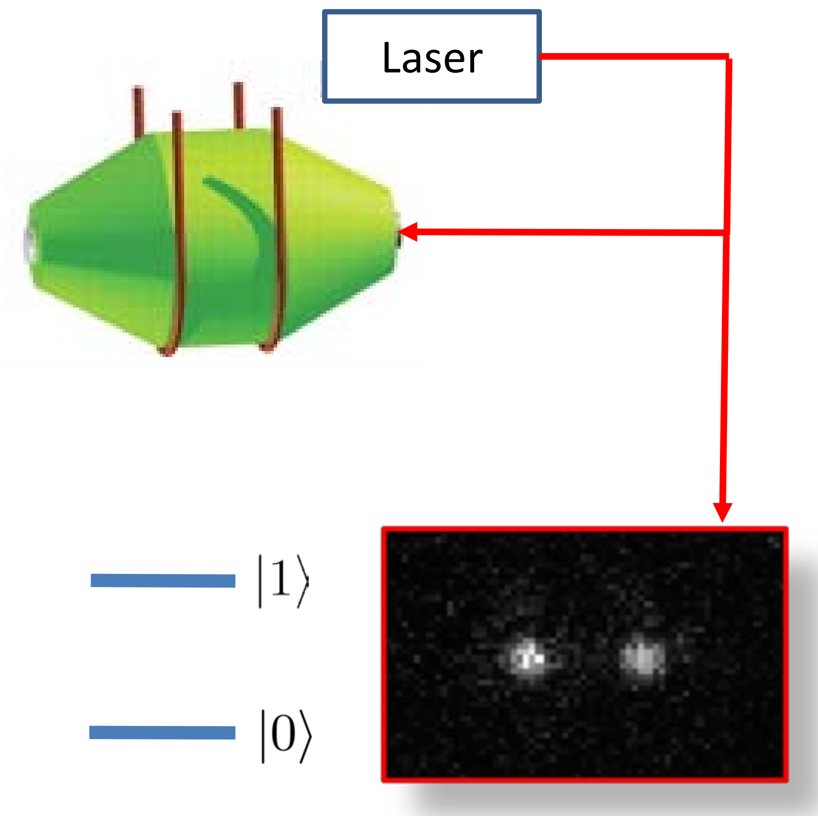


C. Ospelkaus et al. Nature 182, 476 (2011)

Trapped-ions and optical clocks

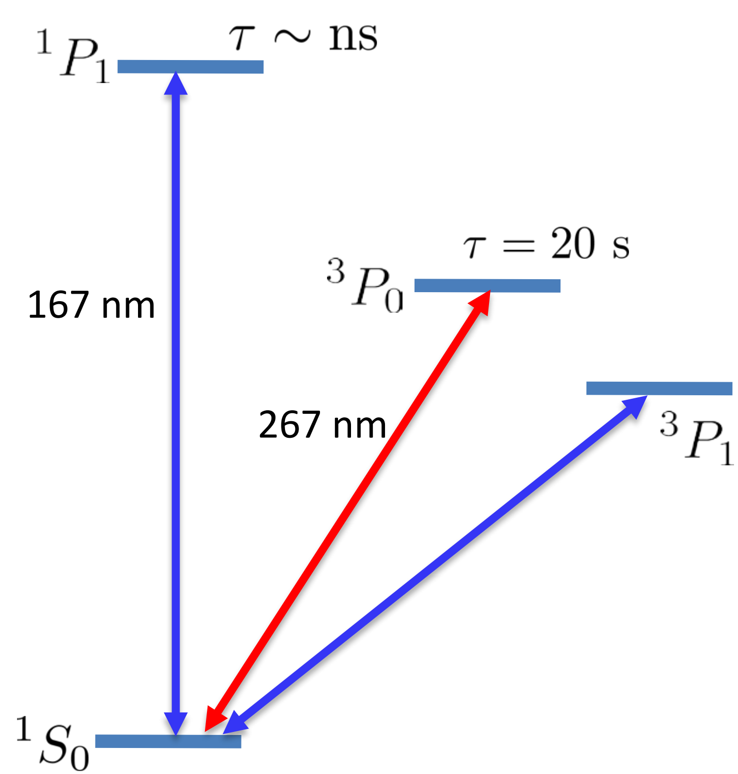
e.g. Rosenband et al., Science 319, 1808 (2008)

Frequency standards



Require very stable ion transition

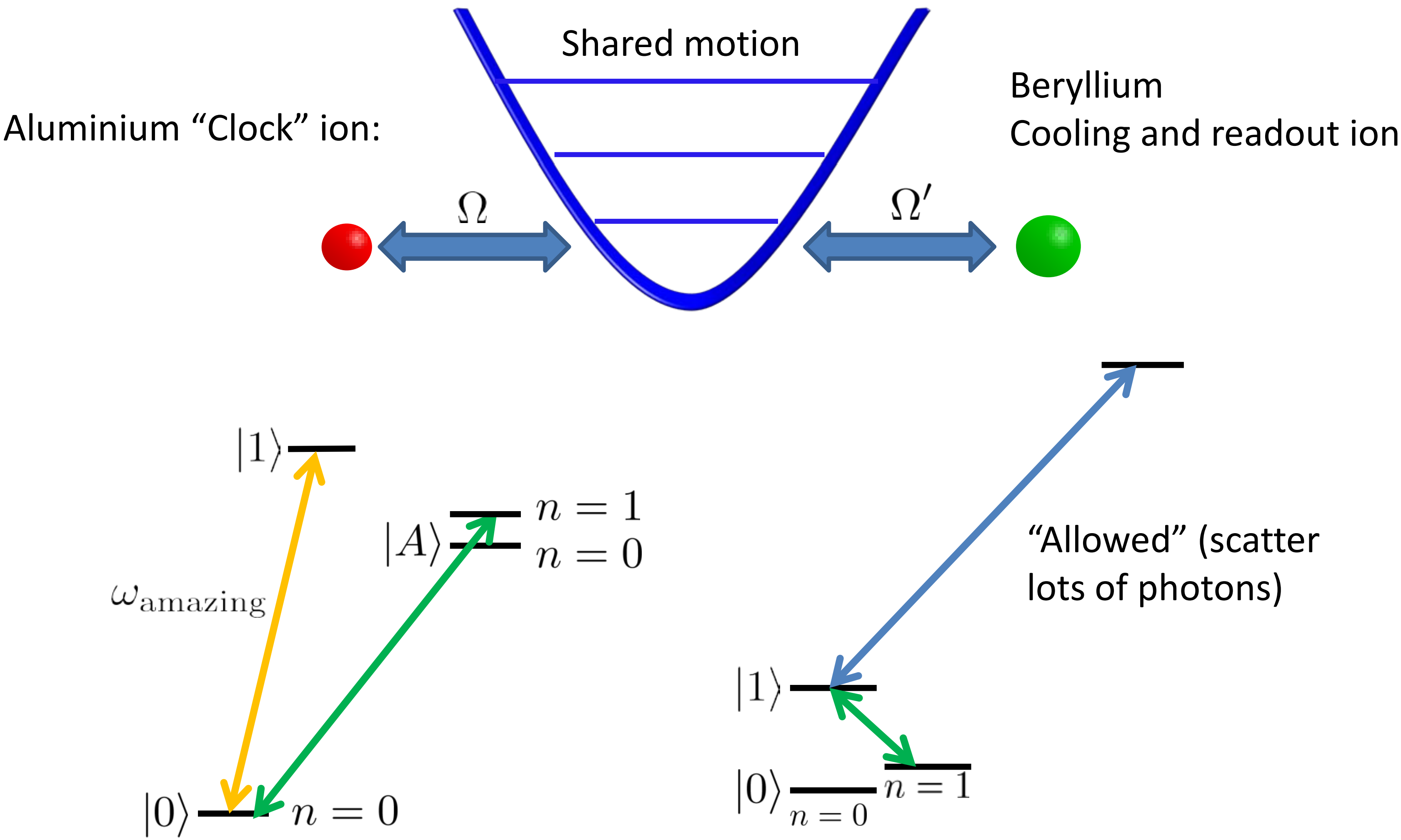
Aluminium ion



Has a very stable transition

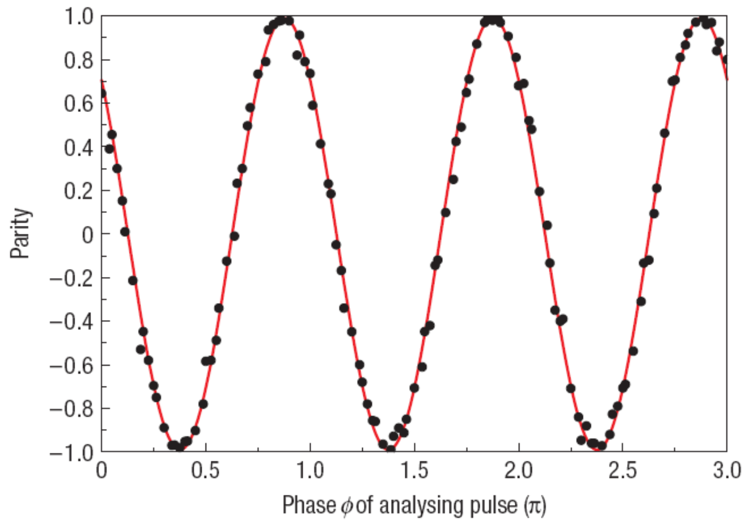
BUT 167 nm is vacuum UV

Atomic clocks – quantum logic readout



Most accurate and precise frequency standards – 8e-18 fractional uncertainty
(Chou et al. PRL *104*, 070802 (2010))

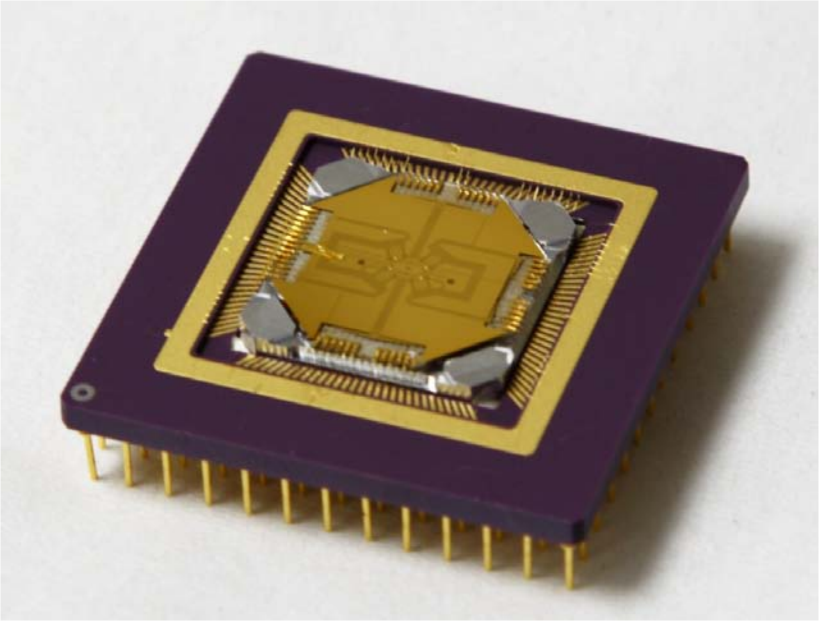
Trapped-ion summary



Have achieved quantum control of up to N ions

Have demonstrated all basic components required to create large scale entangled states

Algorithms & gates include Dense-coding, error-correction, Toffoli, Teleportation, Entanglement purification Entanglment swapping



Working on:

Higher precision

New manipulation methods

Scaling to many ions