

Rydberg atoms

Tobias Thiele

References

- T. Gallagher: Rydberg atoms

Content

- Part 1: Rydberg atoms
- Part 2: A typical beam experiment

Introduction – What is „Rydberg“?

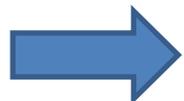
- Rydberg atoms are (any) atoms in state with high principal quantum number n .

Introduction – What is „Rydberg“?

- Rydberg atoms are (any) atoms in state with high principal quantum number n .
- Rydberg atoms are (any) atoms with exaggerated properties

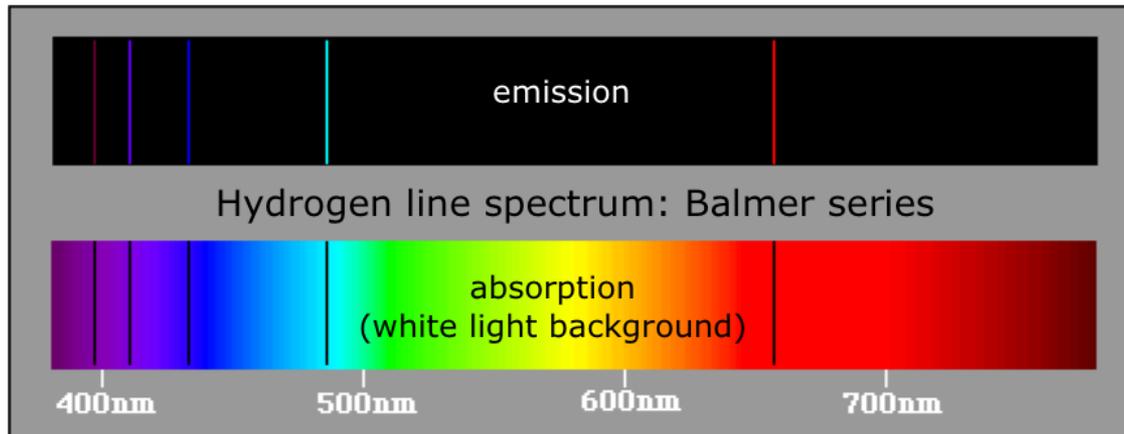
Introduction – What is „Rydberg“?

- Rydberg atoms are (any) atoms in state with high principal quantum number n .
- Rydberg atoms are (any) atoms with exaggerated properties

 *equivalent!*

Introduction – How was it found?

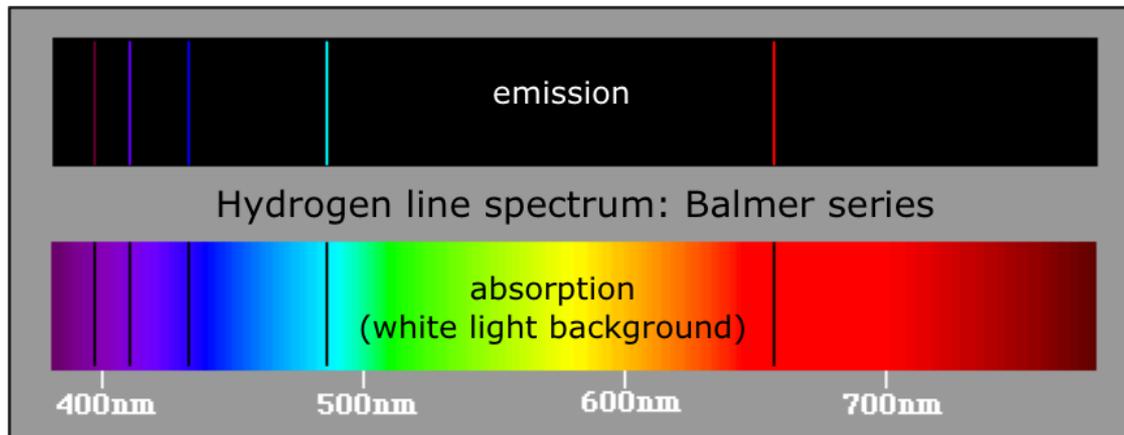
- In 1885: Balmer series:
 - Visible absorption wavelengths of H:



$$\lambda = \frac{bn^2}{n^2 - 4}$$

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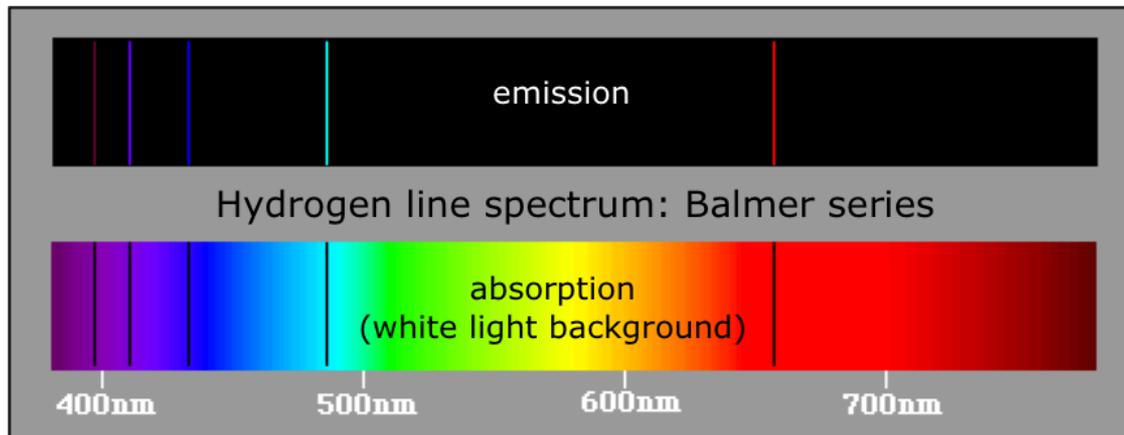


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- Other series discovered by Lyman, Brackett, Paschen, ...

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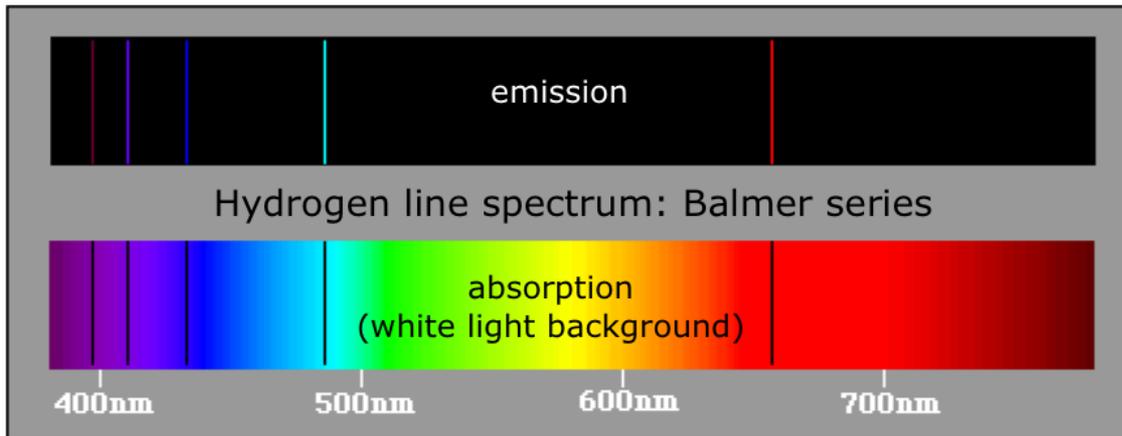
$$\lambda = \frac{bn^2}{n^2 - 4}$$

- Other series discovered by Lyman, Brackett, Paschen, ...

- Summarized by Johannes Rydberg: $\tilde{\nu} = \tilde{\nu}_{\infty} - \frac{Ry}{n^2}$

Introduction – Generalization

- In 1885: Balmer series:
 - Visible absorption wavelengths of H:



$$\lambda = \frac{bn^2}{n^2 - 4}$$

- Other series discovered by Lyman, Brackett, Paschen, ...

- Quantum Defect was found for other atoms: $\tilde{\nu} = \tilde{\nu}_\infty - \frac{Ry}{(n - \delta_l)^2}$

Introduction – Rydberg formula?

- Energy follows Rydberg formula:

$$E = E_{\infty} - \frac{hRy}{(n - \delta_l)^2}$$



Introduction – Rydberg formula?

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$$E = E_{\infty} - \frac{\overset{\approx 13.6 \text{ eV}}{hRy}}{(n - \delta_l)^2}$$

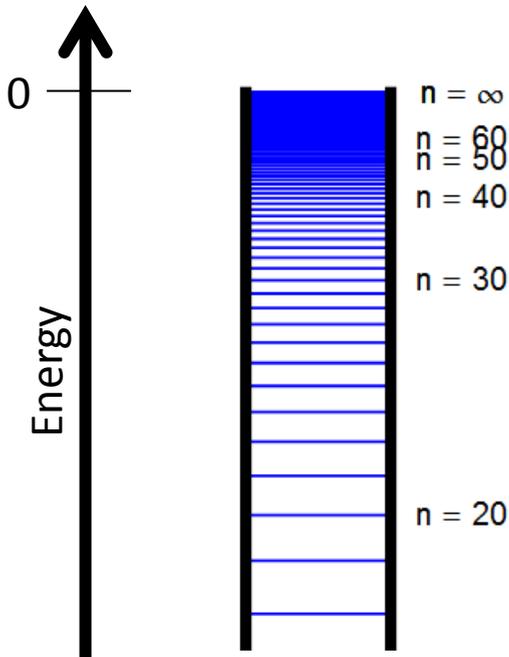


Introduction – Hydrogen?

- Energy follows Rydberg formula:

$$E = E_{\infty} - \frac{hRy}{(n - \delta_l)^2} = -\frac{hRy}{n^2}$$

Annotations: E_{∞} is boxed with a blue '0' and an arrow pointing to it. hRy is circled in blue, with a box containing $\approx 13.6 \text{ eV}$ above it. δ_l is boxed with a blue '0' and an arrow pointing to it.

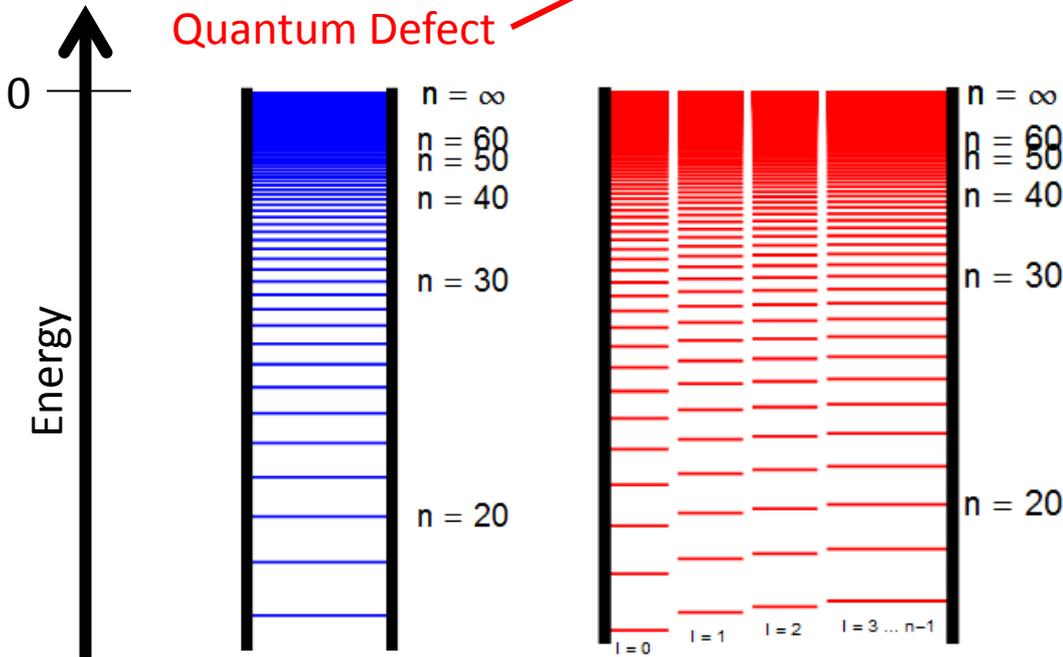


Quantum Defect?

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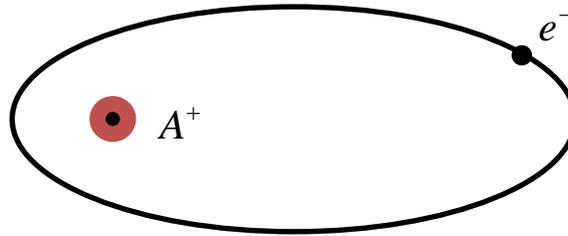
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Quantum Defect



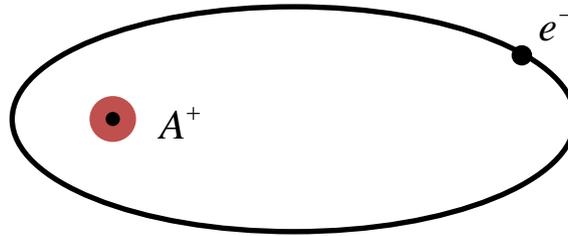
Rydberg Atom Theory

- Rydberg Atom



Rydberg Atom Theory

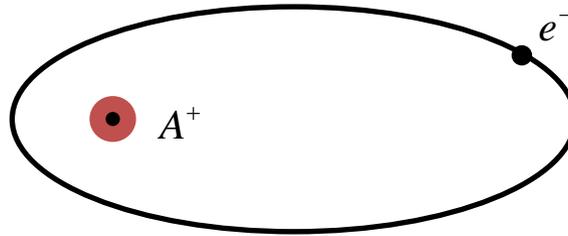
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- Almost like Hydrogen
 - Core with one positive charge
 - One electron

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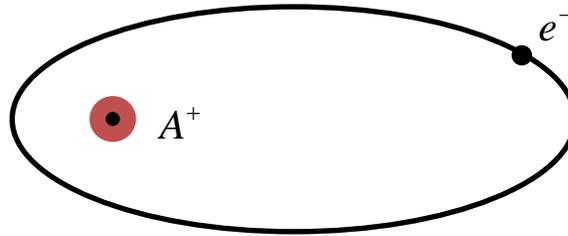
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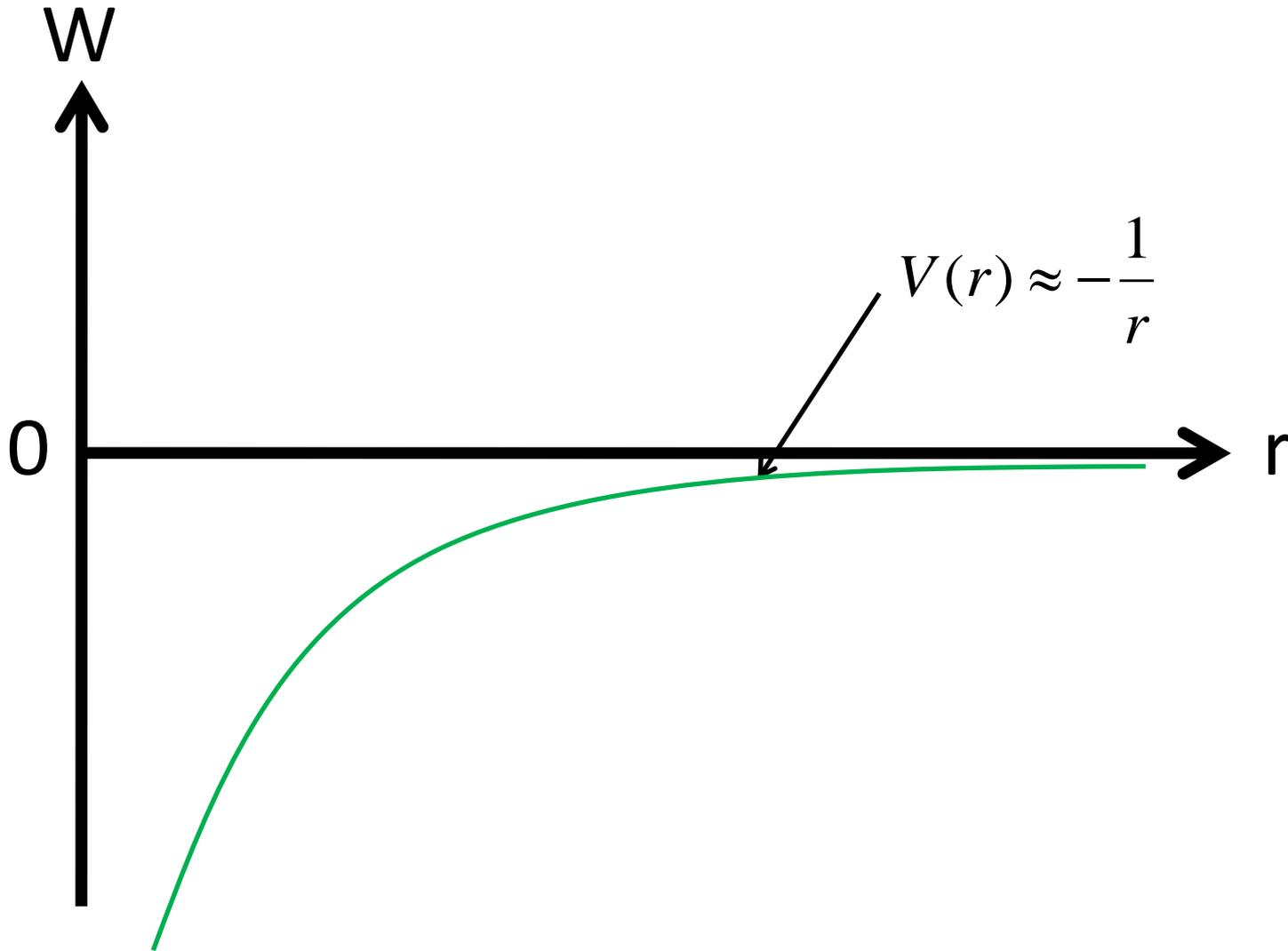
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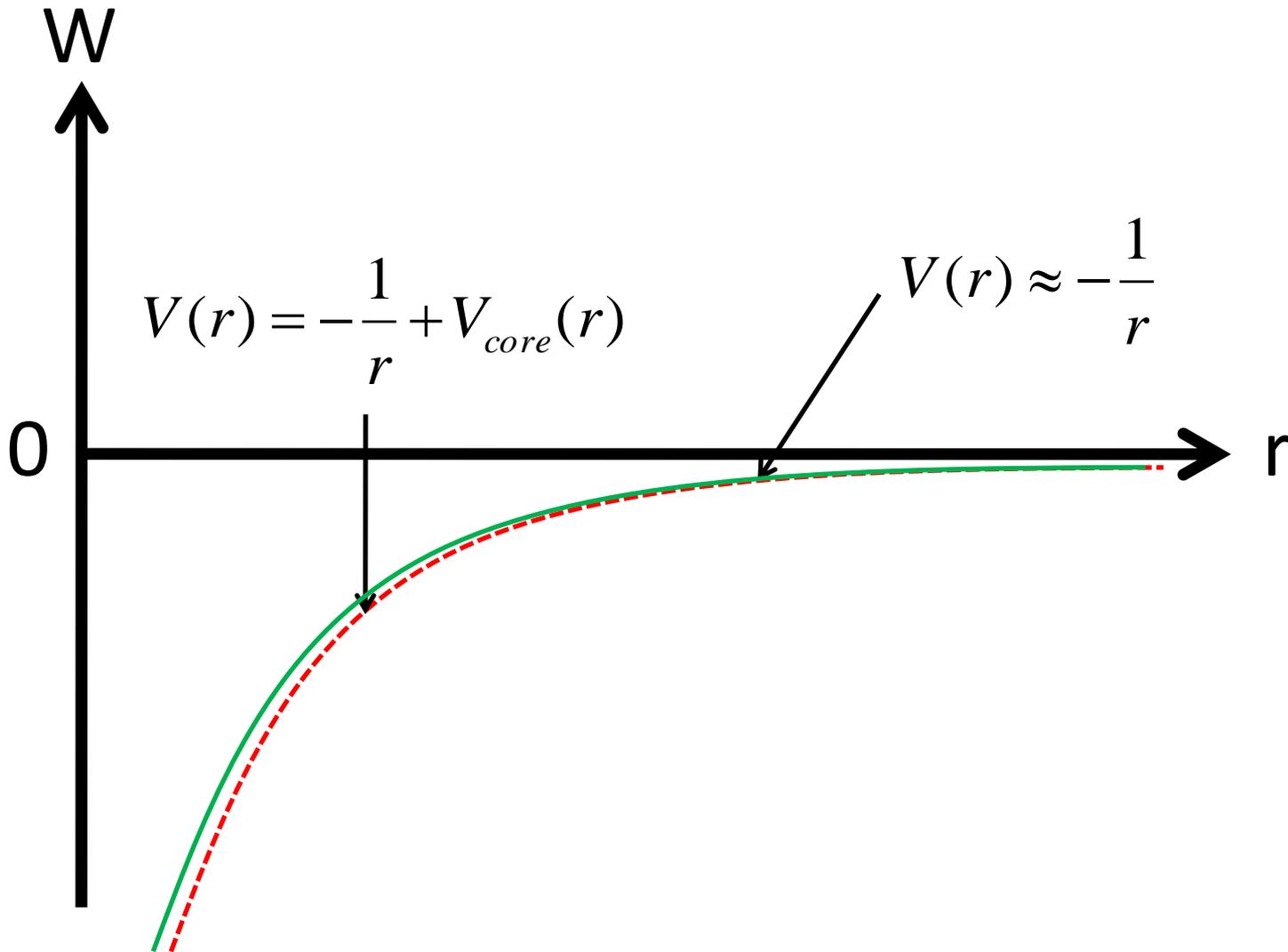


- Almost like Hydrogen
 - Core with one positive charge
 - One electron
- What is the difference?
 - No difference in angular momentum states

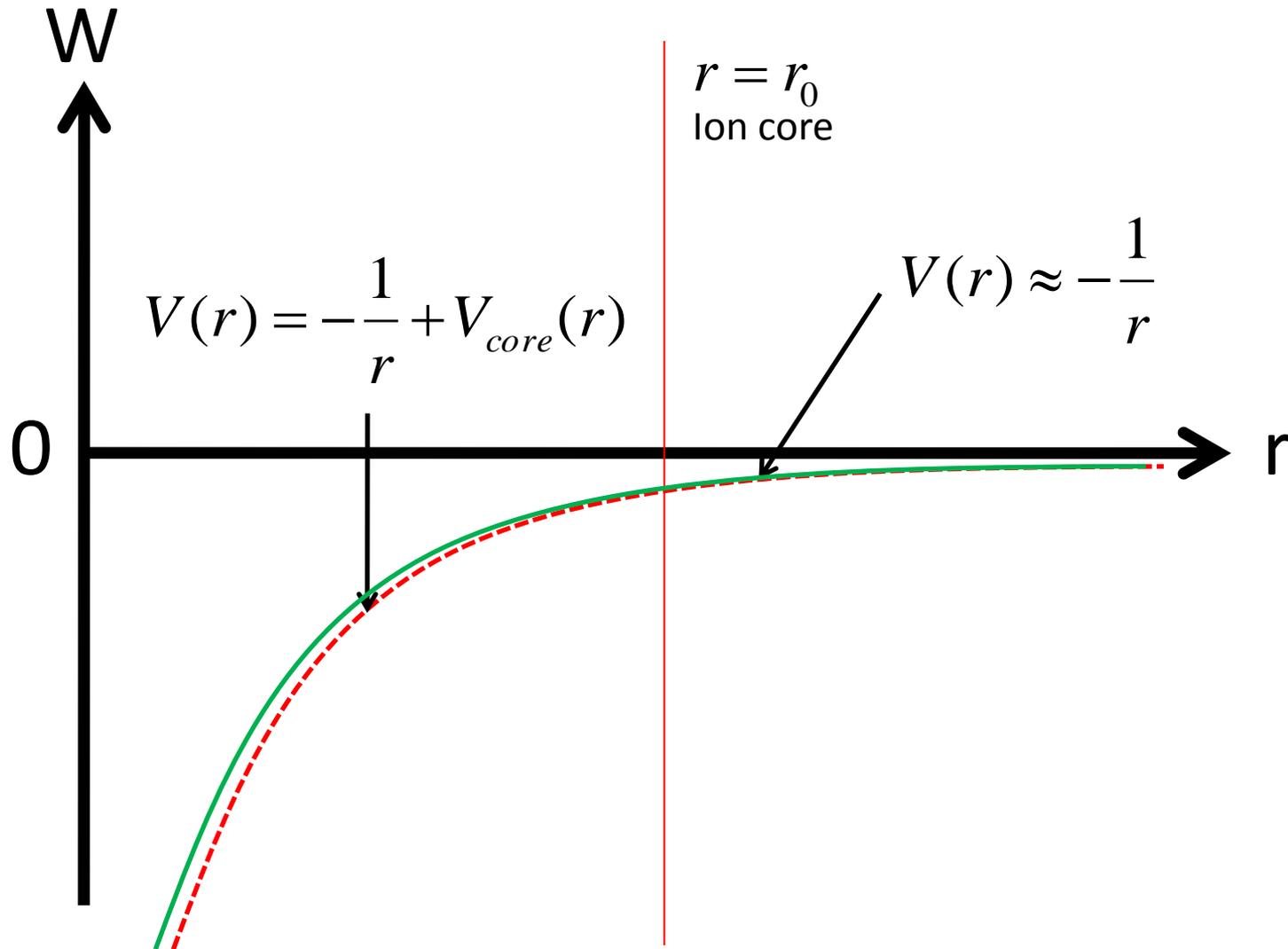
Radial parts-Interesting regions



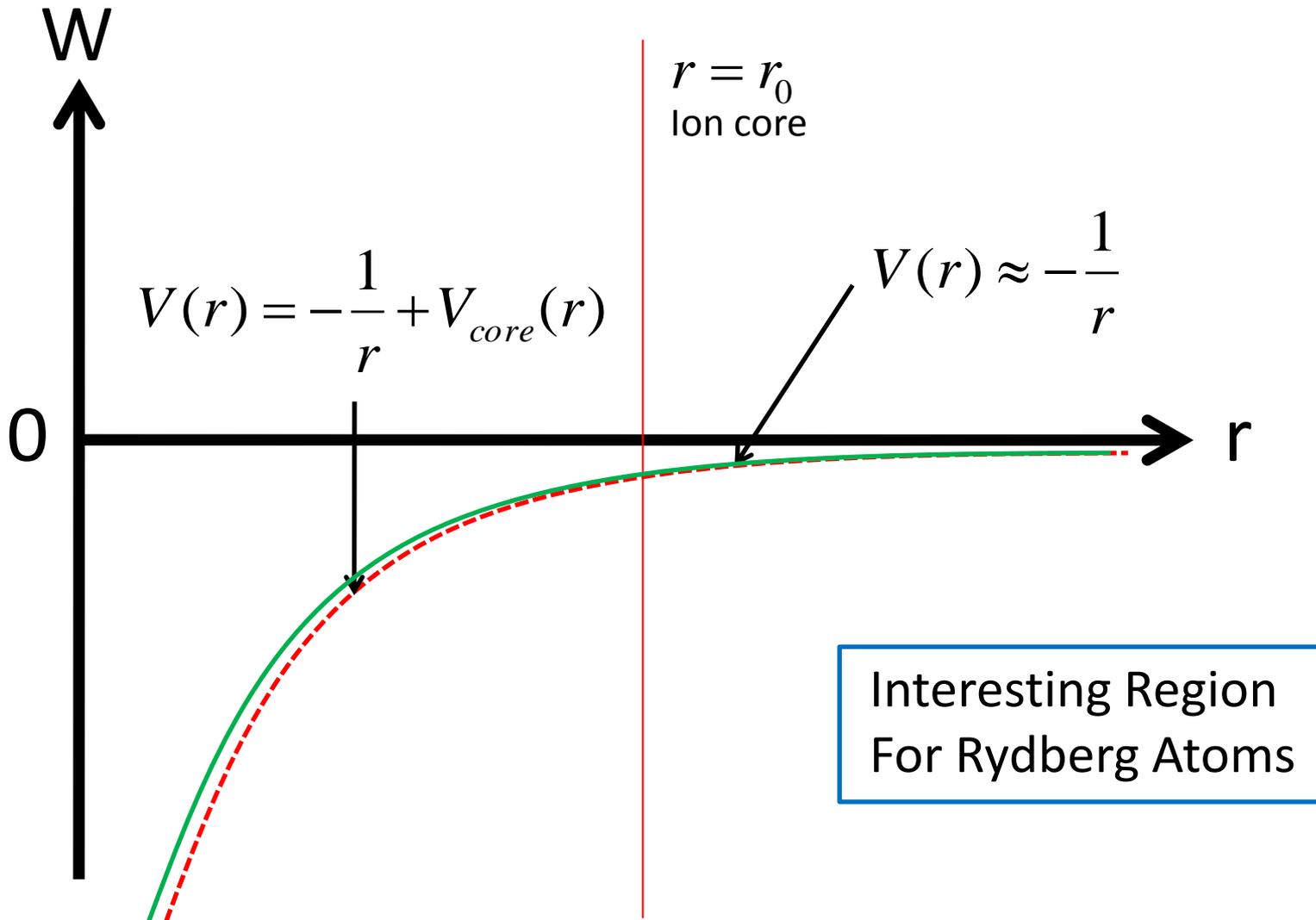
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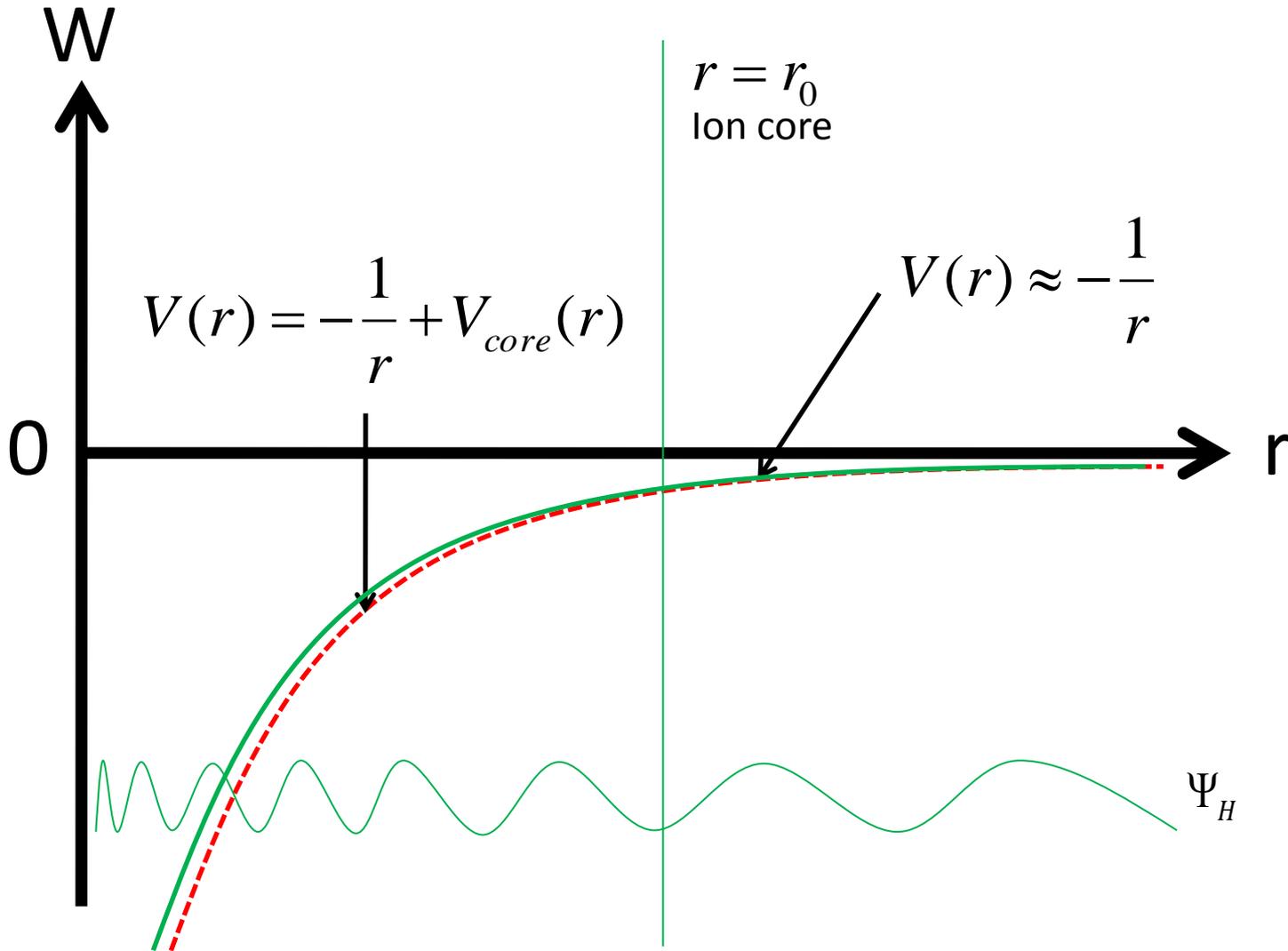
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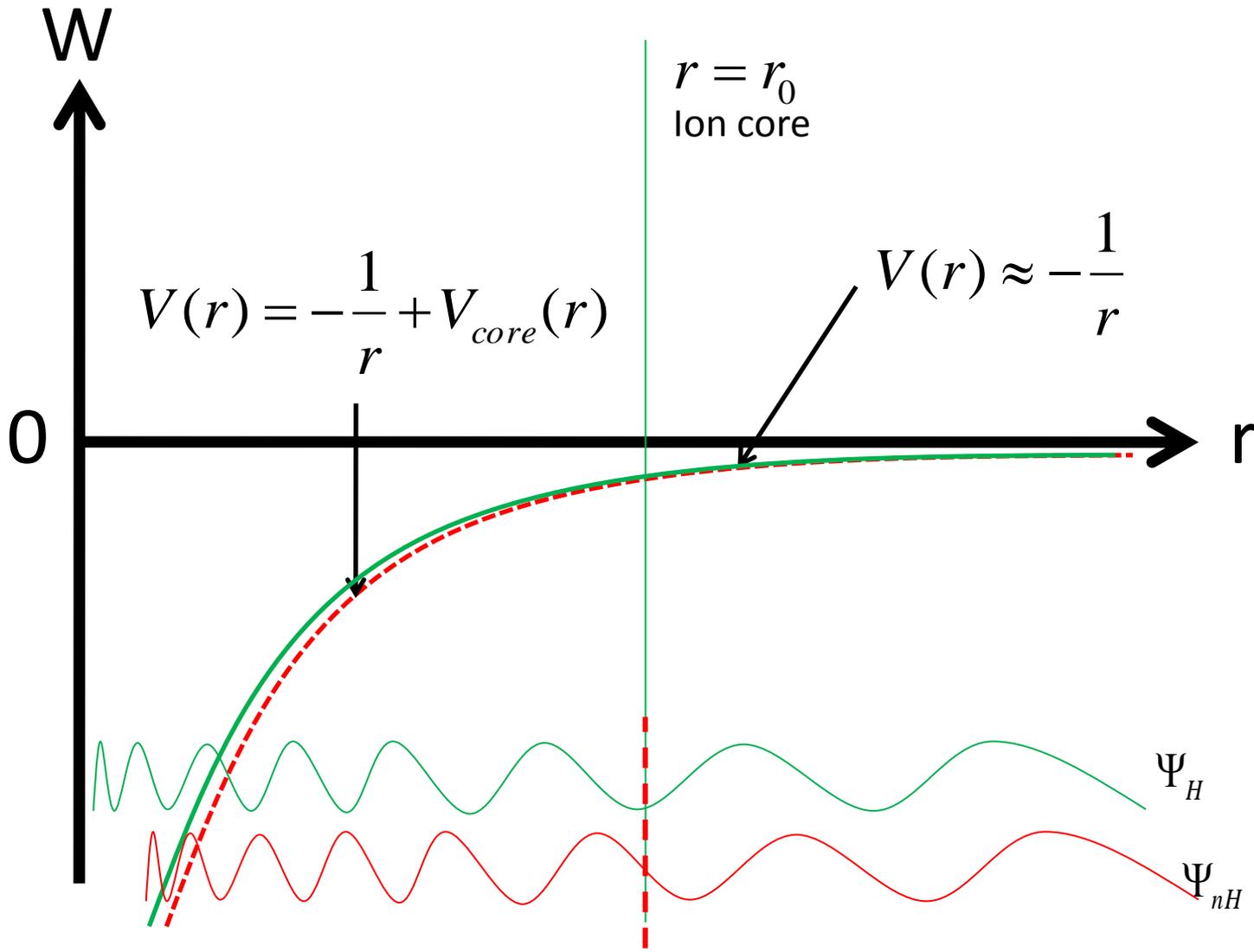
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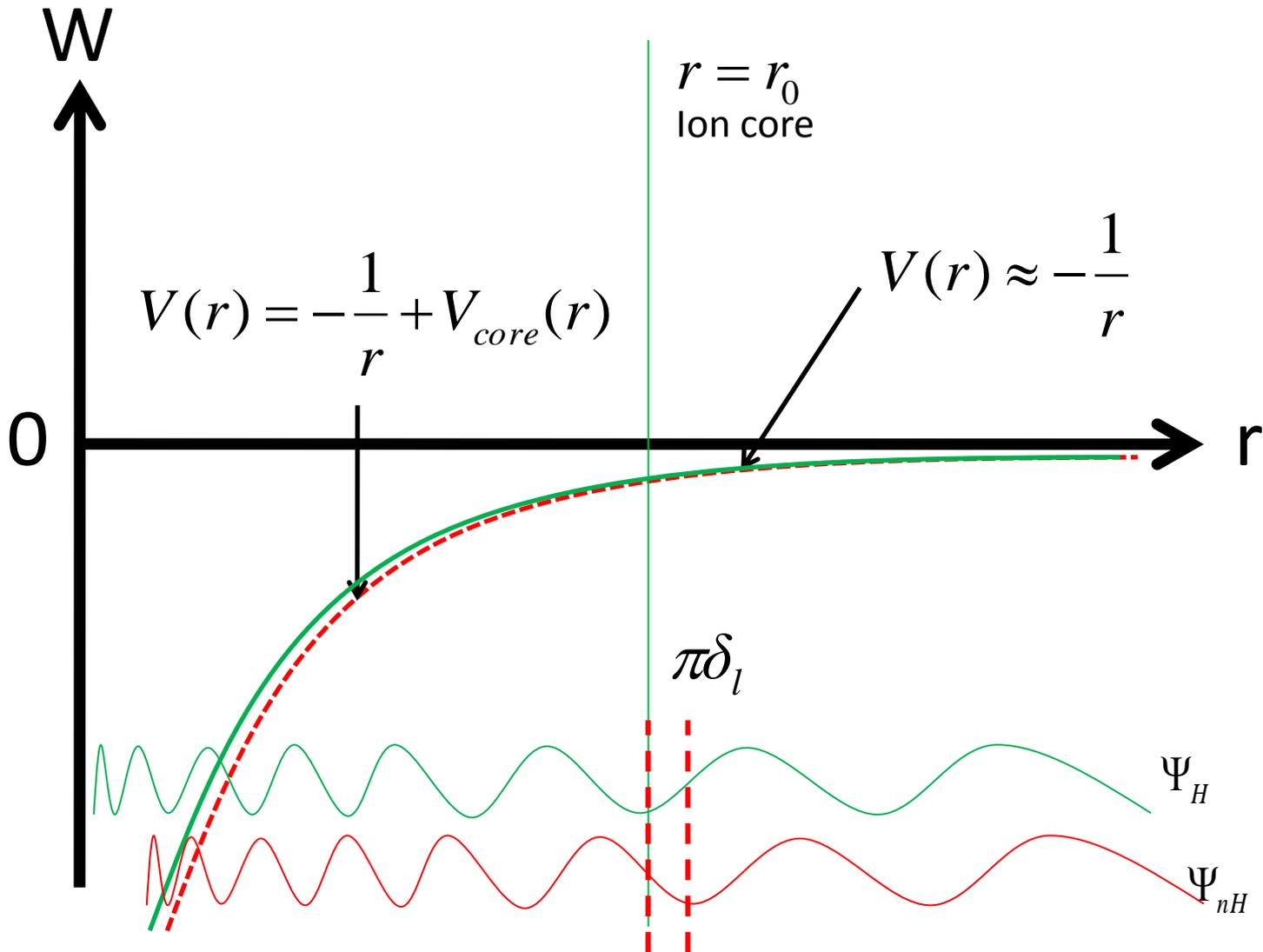
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(Helium) Energy Structure

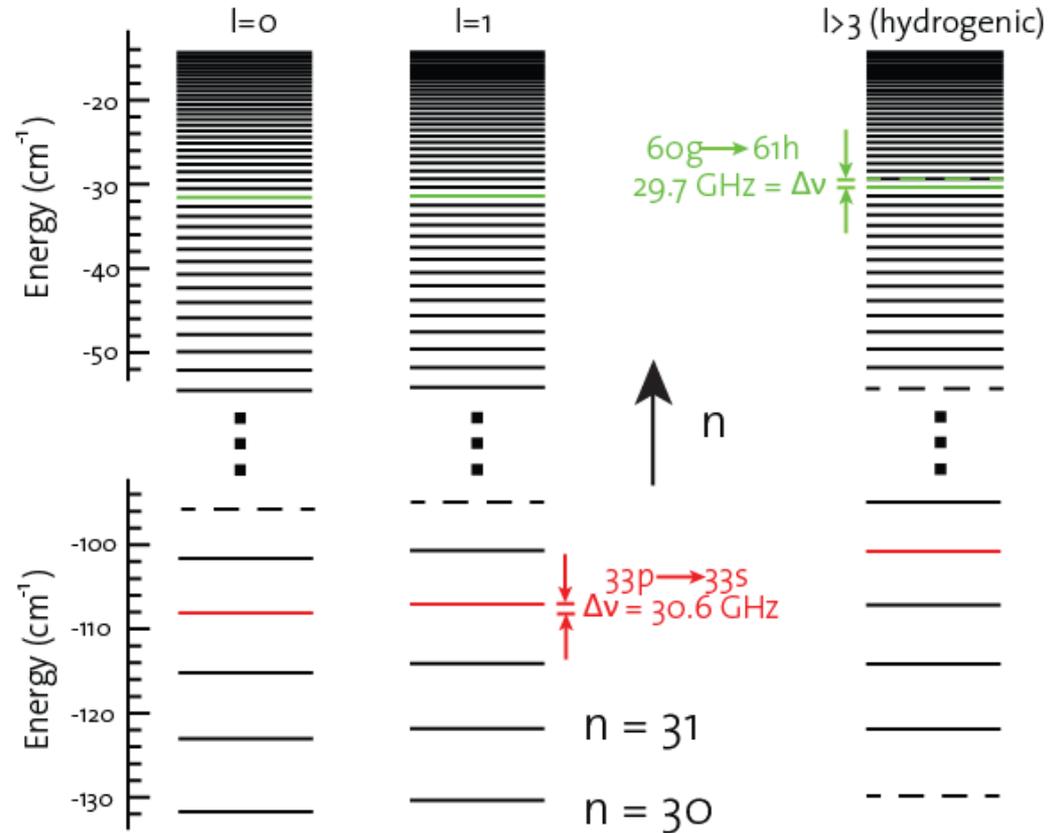
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- δ_l usually measured
 - Only large for low l (s,p,d,f)

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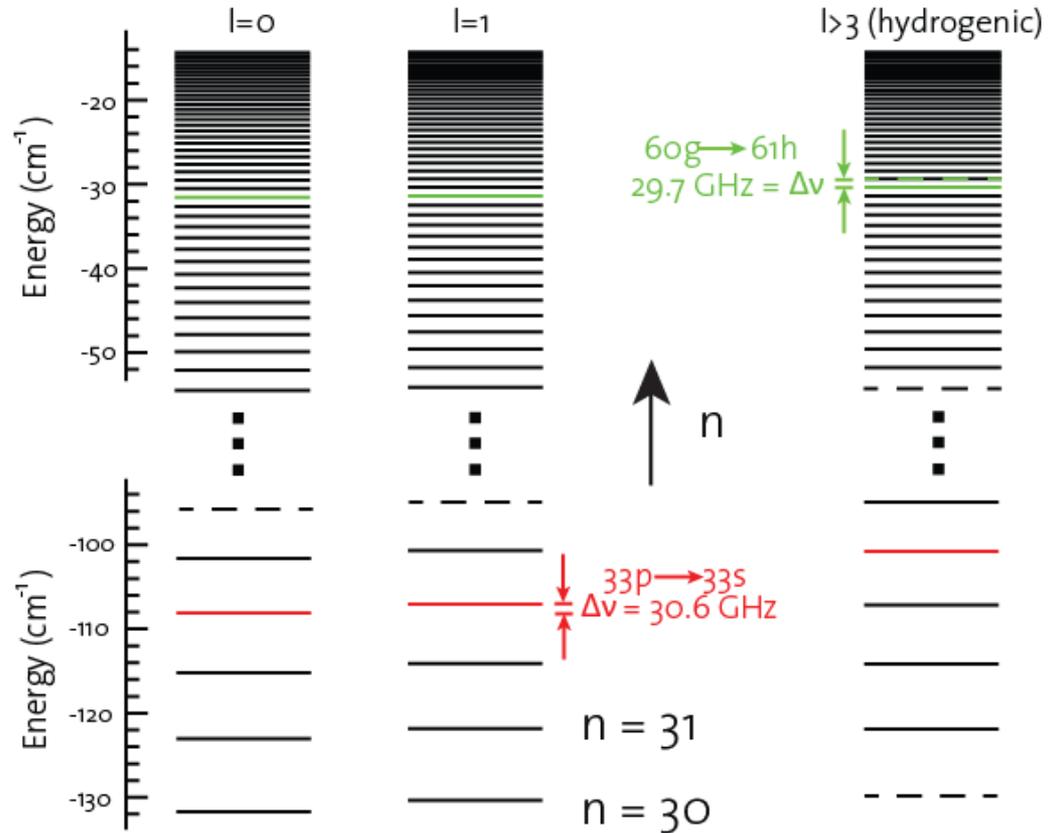
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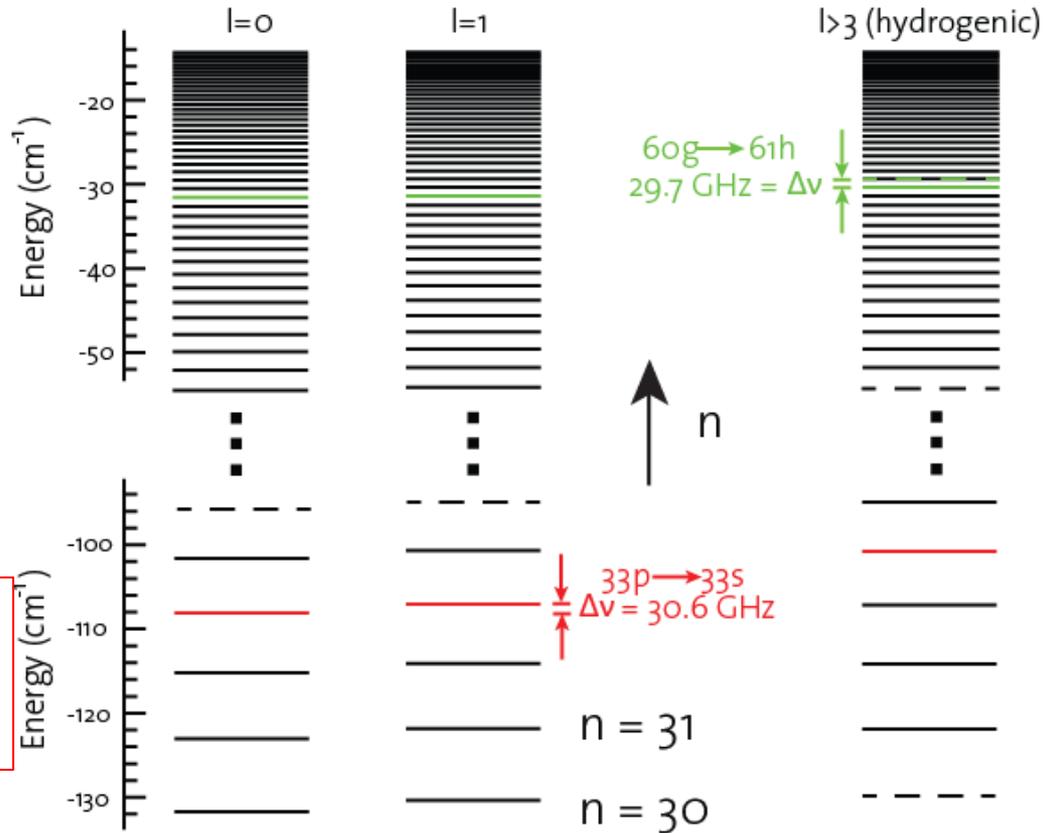
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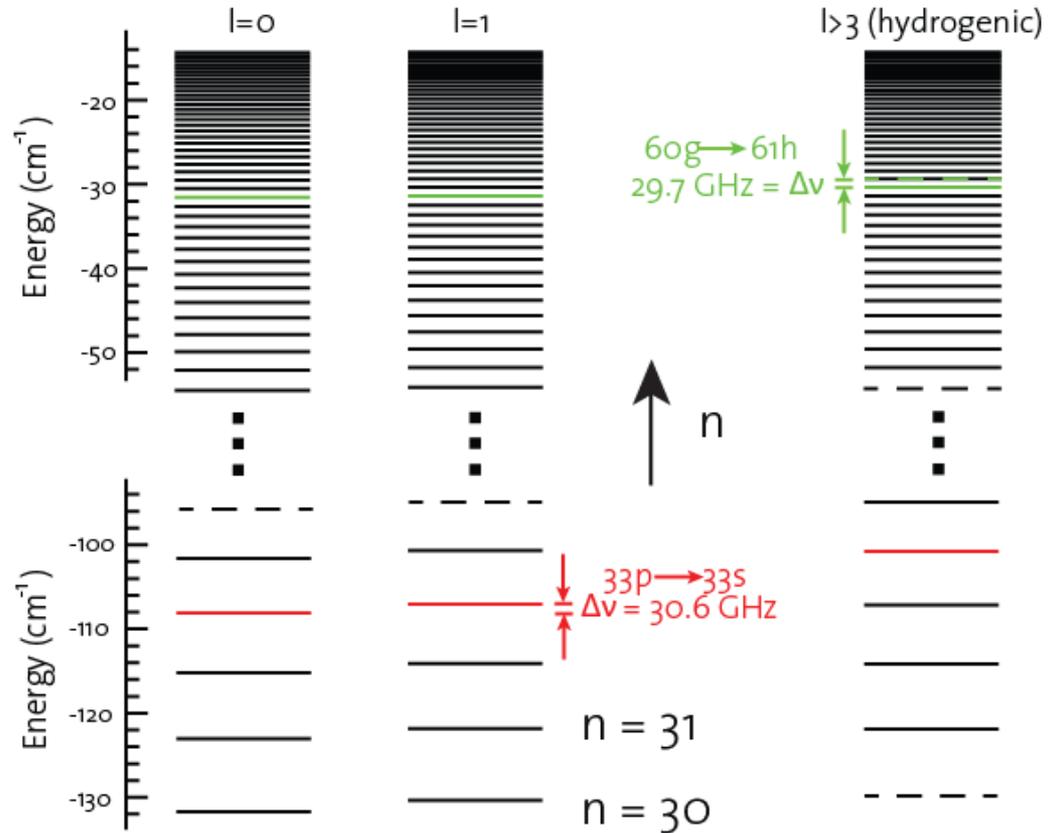
Excentric orbits penetrate into core.
 Large deviation from Coulomb.
 Large phase shift \rightarrow large quantum defect

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$$\frac{dW}{dn} = \frac{1}{(n - \delta_l)^3}$$



Electric Dipole Moment

- Electron most of the time far away from core

$W = -\frac{1}{2(n - \delta_l)^2}$	$\frac{dW}{dn} = \frac{1}{(n - \delta_l)^3}$
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- We find electric Dipole Moment
 - $\langle \Psi_f | \vec{d} | \Psi_i \rangle \propto \langle r \rangle \langle l \pm 1 | \cos(\theta) | l \rangle \propto n^2$

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 - $\langle \Psi_f | \vec{d} | \Psi_i \rangle \propto \langle r \rangle \langle l \pm 1 | \cos(\theta) | l \rangle \propto n^2$
- Cross Section: $\sigma \propto \langle r \rangle^2 \propto n^4$

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Stark Effect $H\Psi = (H_0 + \vec{d}\vec{F})\Psi = E\Psi$

- For non-Hydrogenic Atom (e.g. Helium)
 - „Exact“ solution by numeric diagonalization of

$$\langle \Psi_f | H | \Psi_i \rangle = \langle \Psi_f | H_0 | \Psi_i \rangle + \langle \Psi_f | \vec{d} | \Psi_i \rangle \vec{F}$$

in undisturbed (standard) basis (\tilde{n}, l, m)

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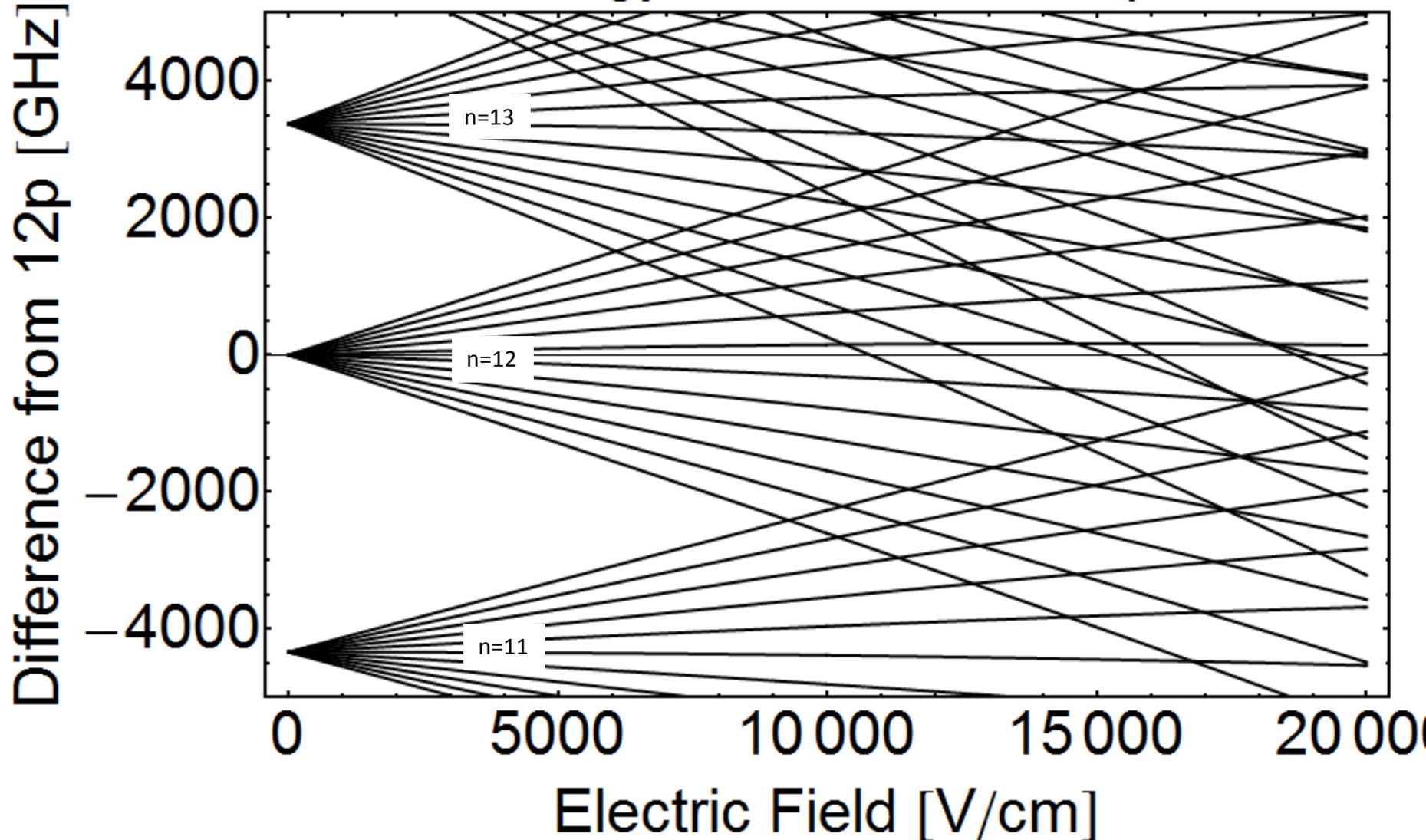
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Numerov

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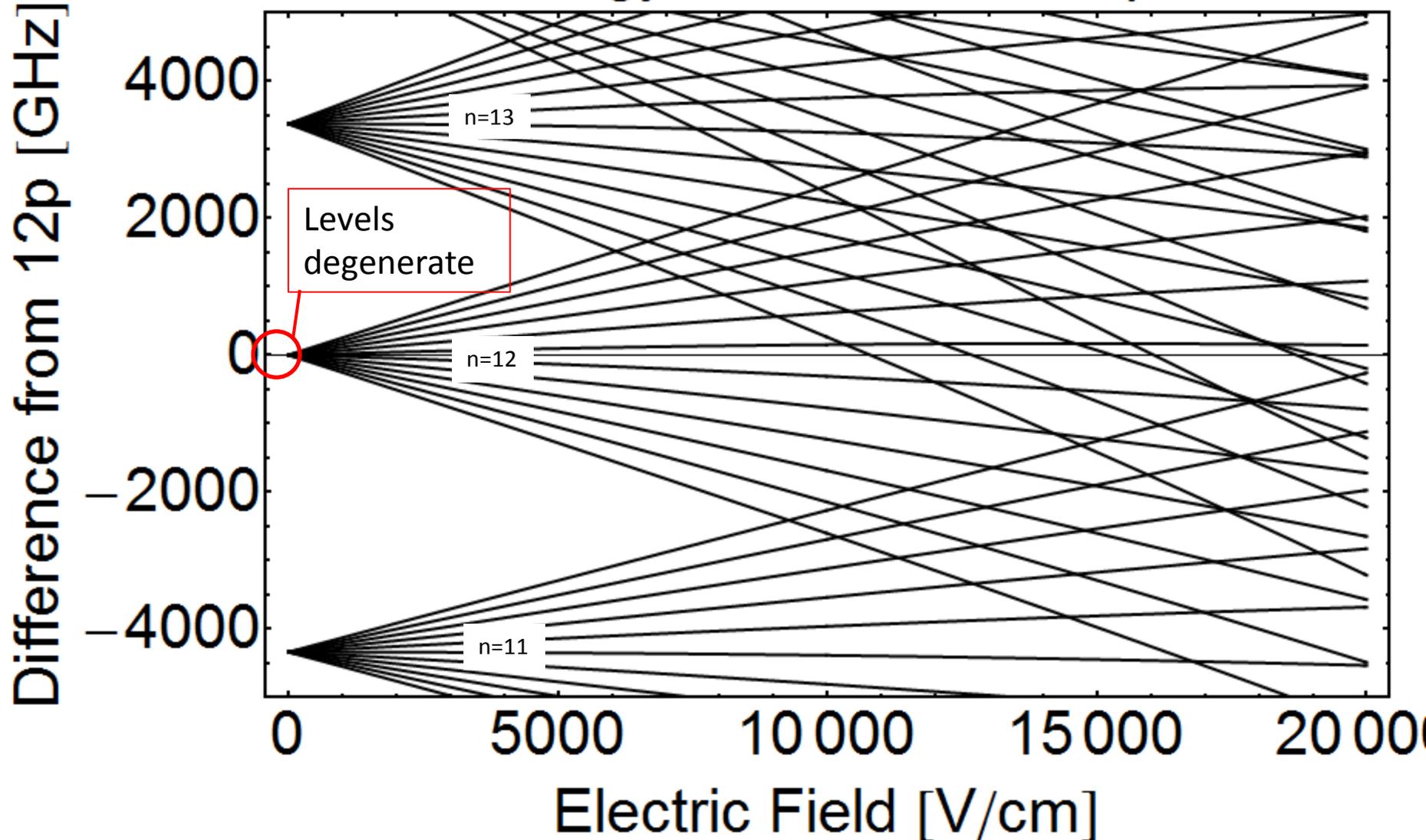
Stark Map Hydrogen

Energy levels around 12p



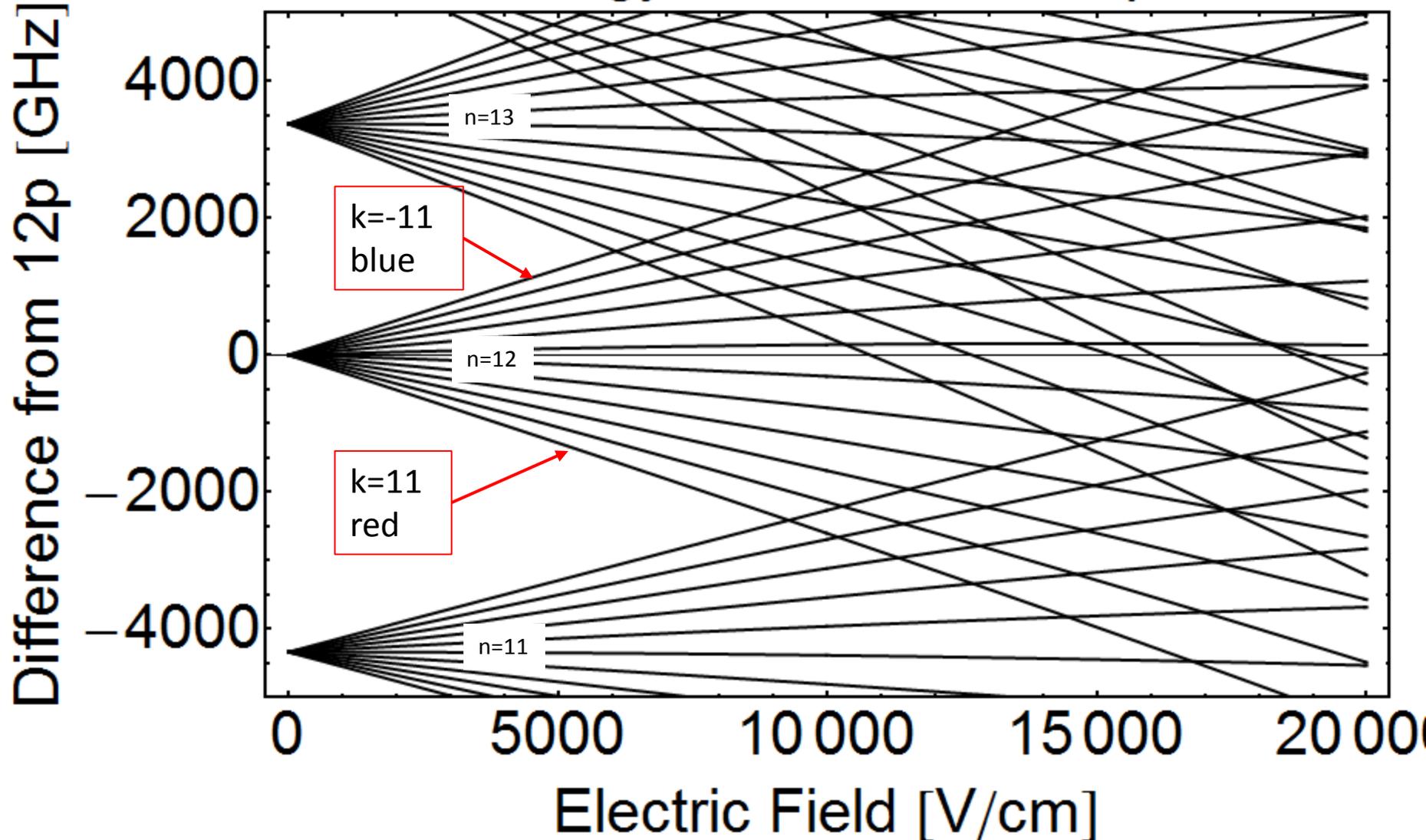
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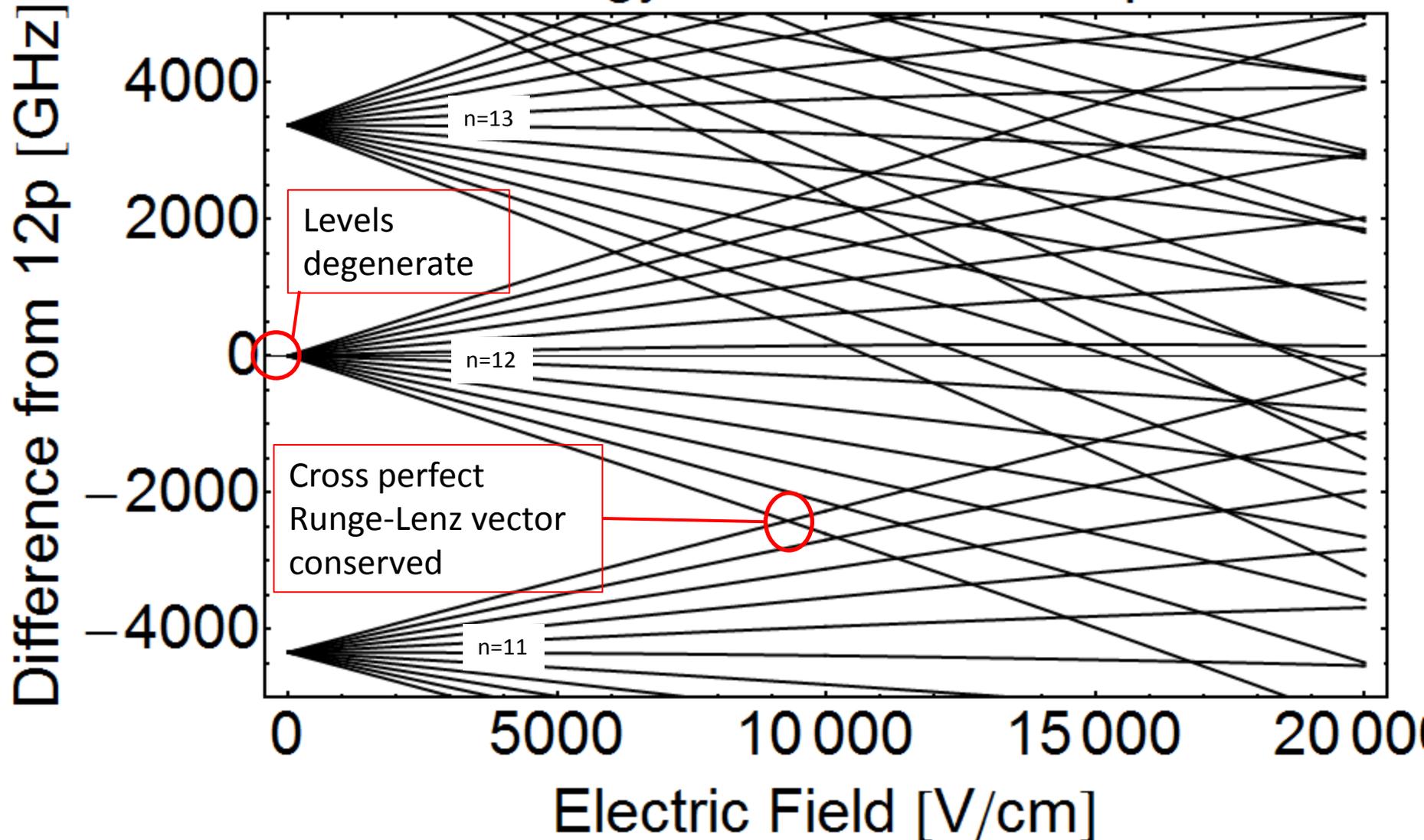
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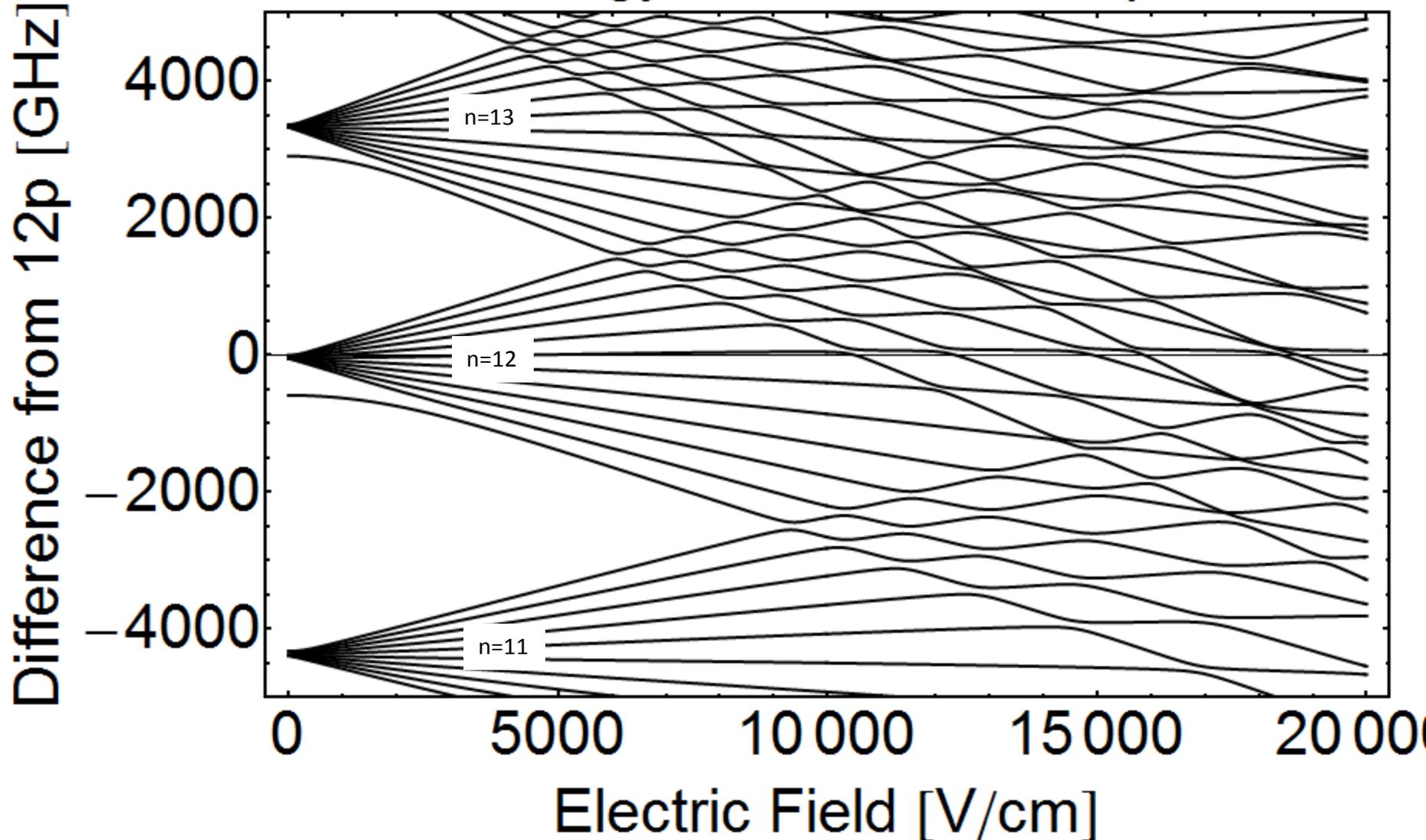
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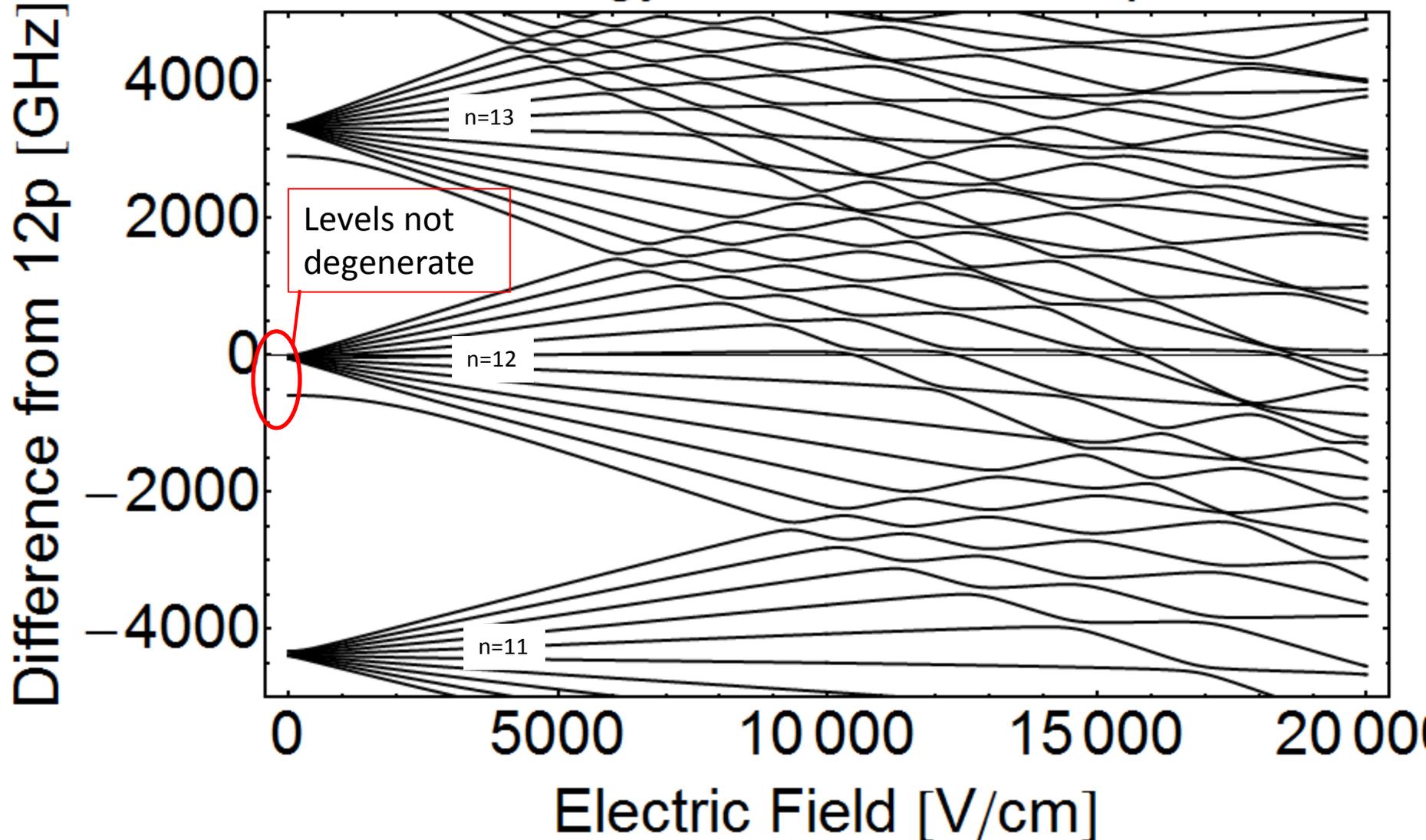
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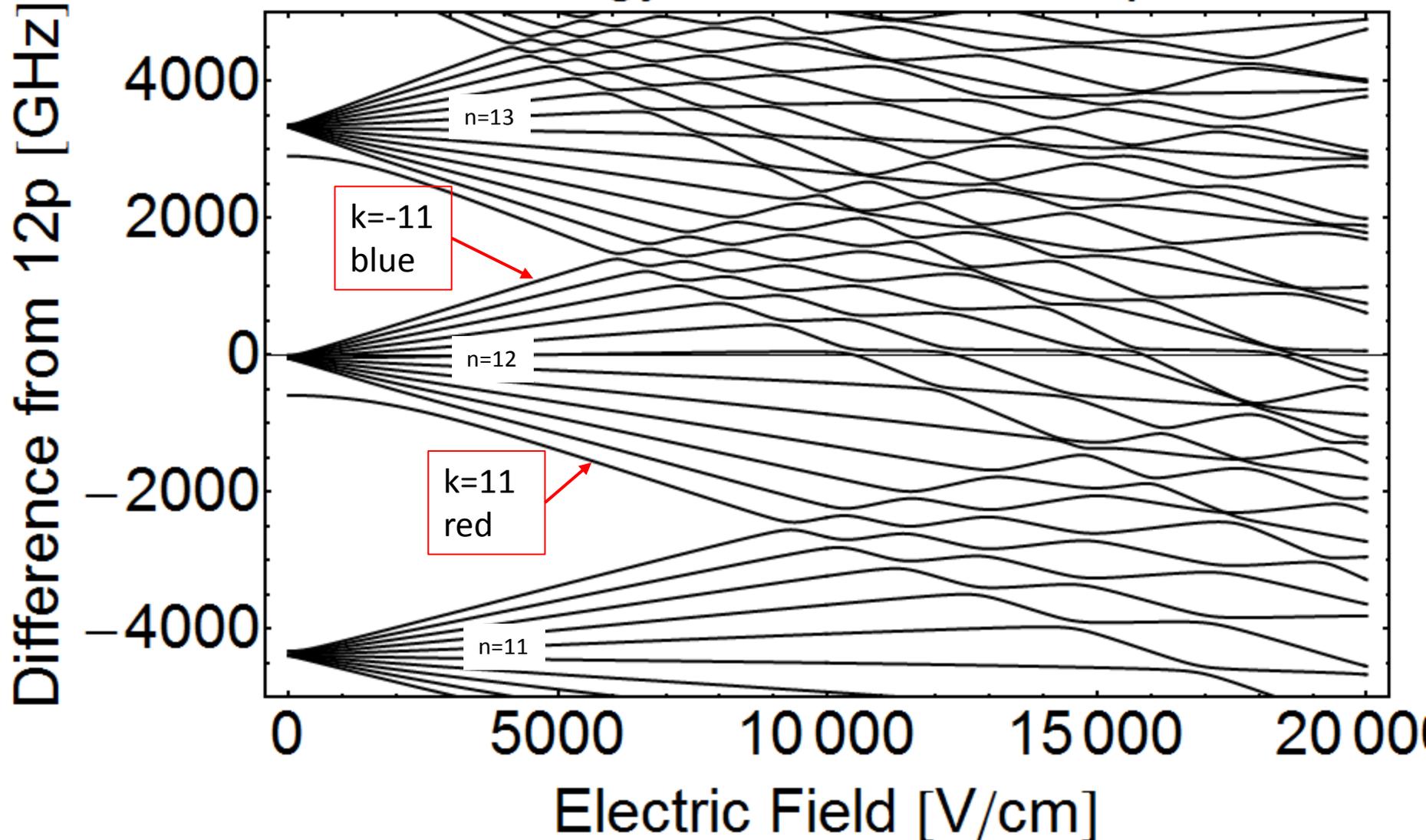
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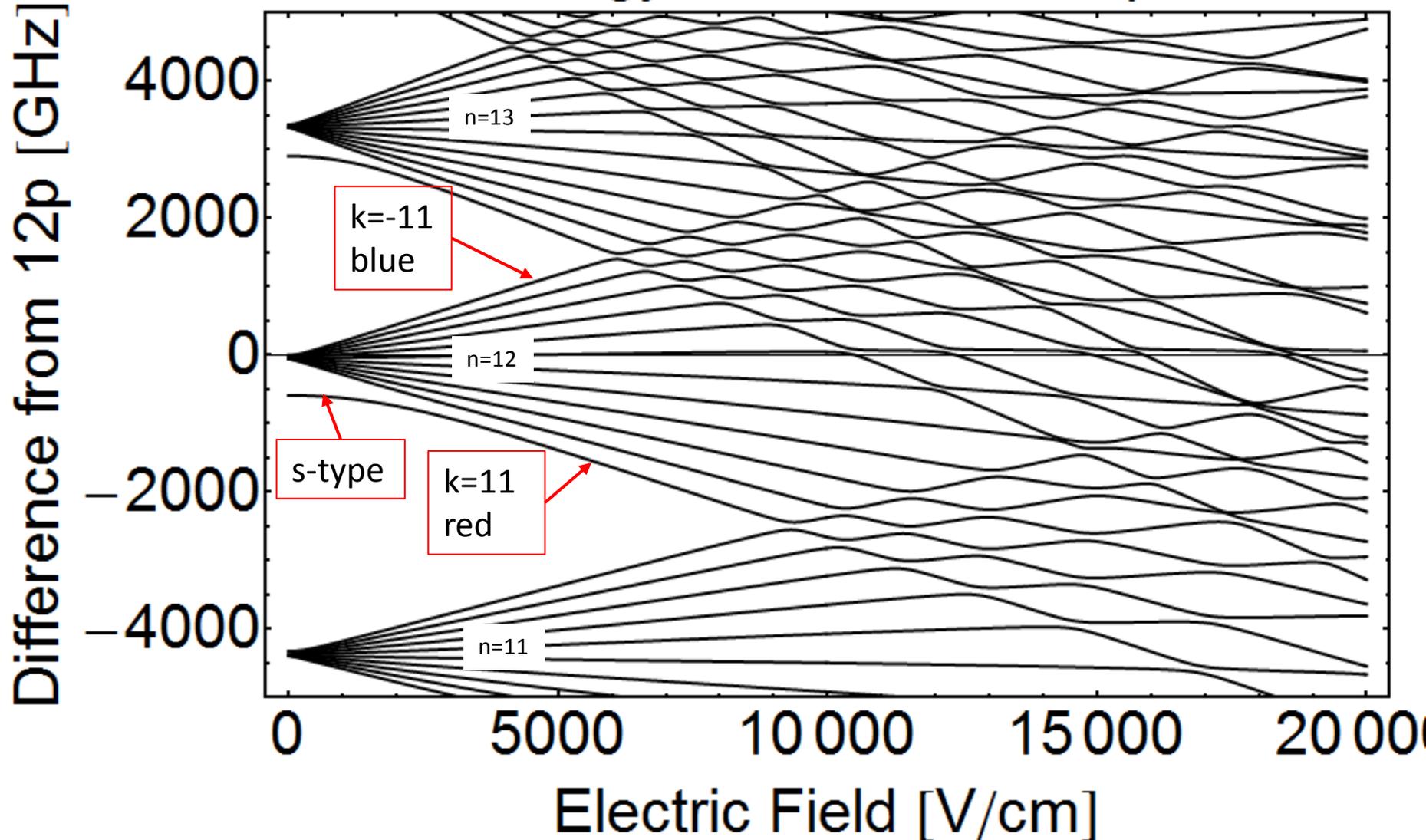
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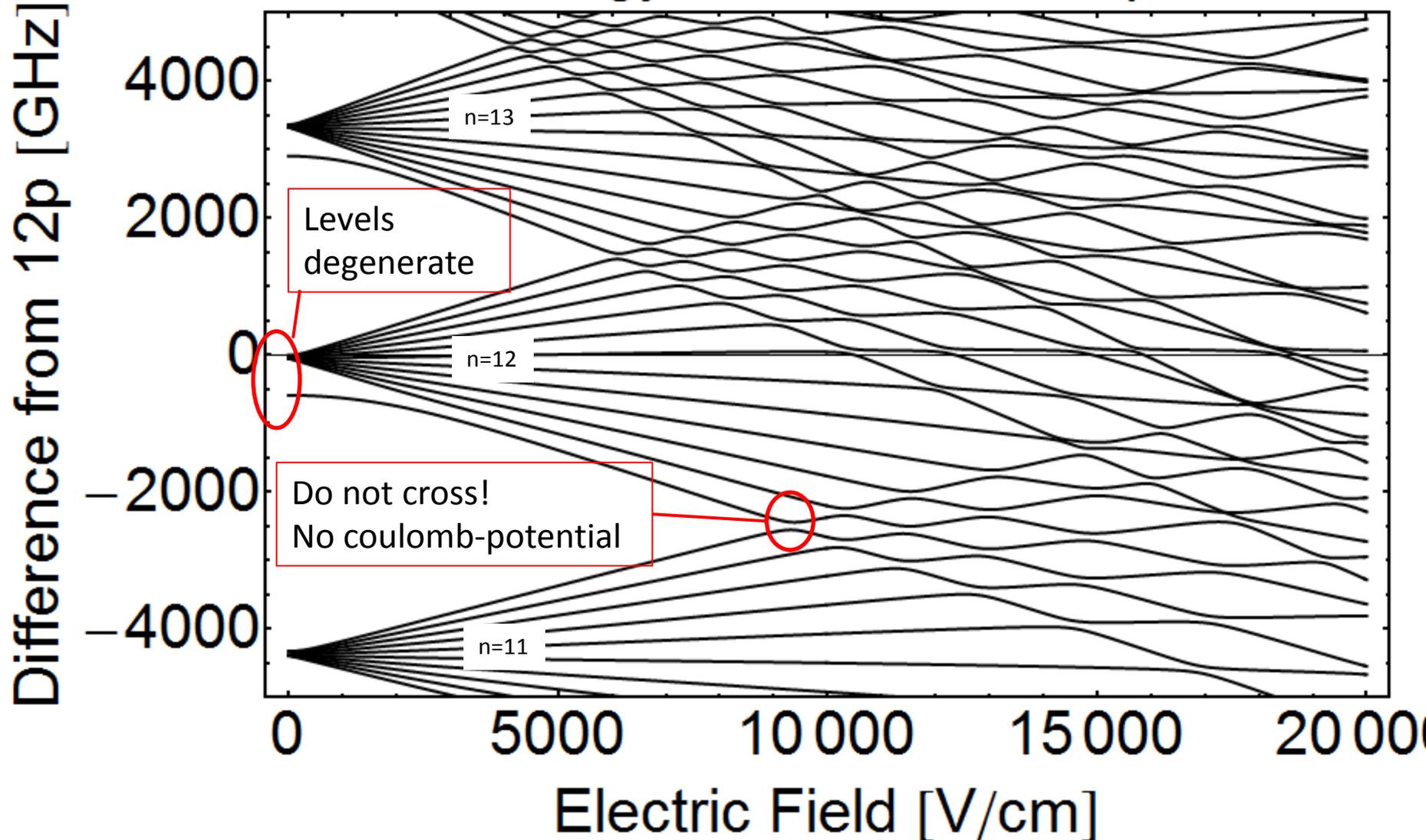
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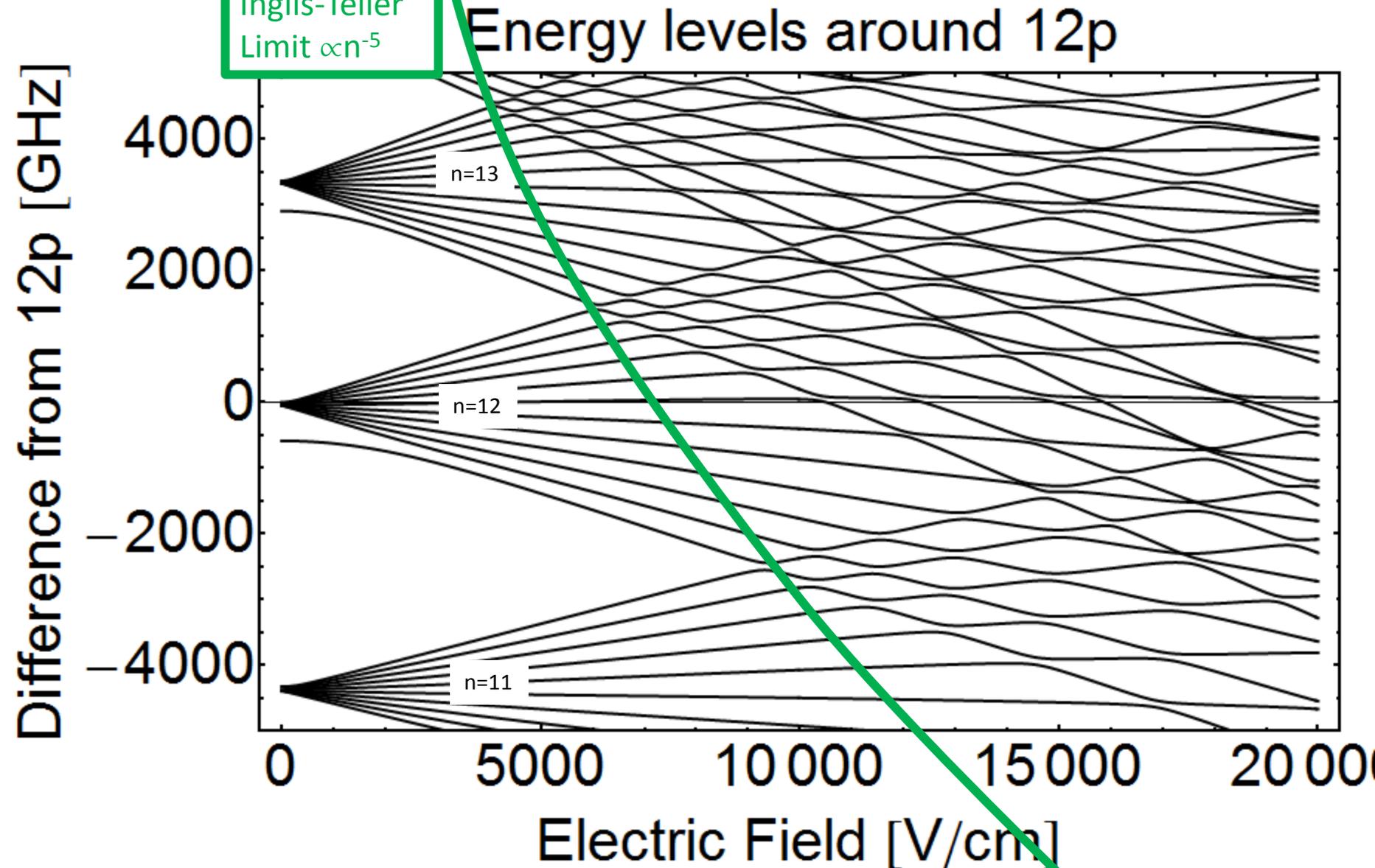
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Hydrogen Atom in an electric Field

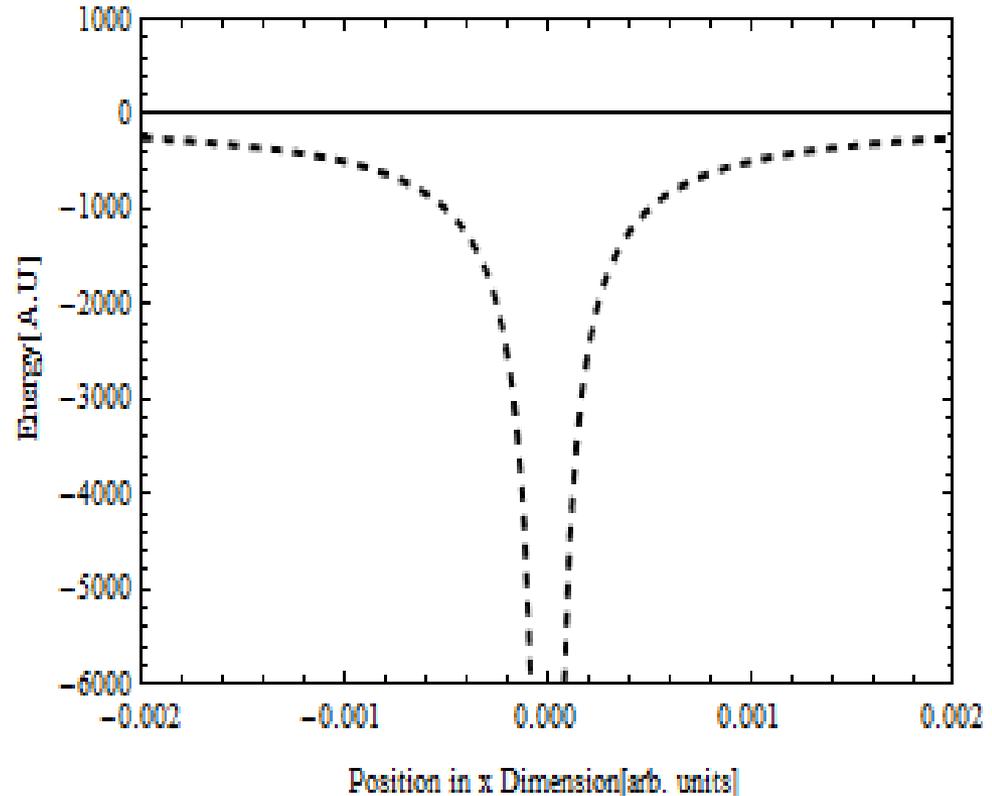
- Rydberg Atoms very sensitive to electric fields
 - Solve: $H\Psi = (H_0 + \vec{d}\vec{F})\Psi = E\Psi$ in parabolic coordinates

- Energy-Field dependence: Perturbation-Theory

$$W = -\frac{1}{2n^2} - \frac{3}{2} F \underbrace{(n_1 - n_2)}_k n + \frac{F}{16} n^4 \left(17n^2 - 3 \underbrace{(n_1 - n_2)^2}_k - 9m^2 + 19 \right) + O(n^5)$$

Hydrogen Atom in an electric Field

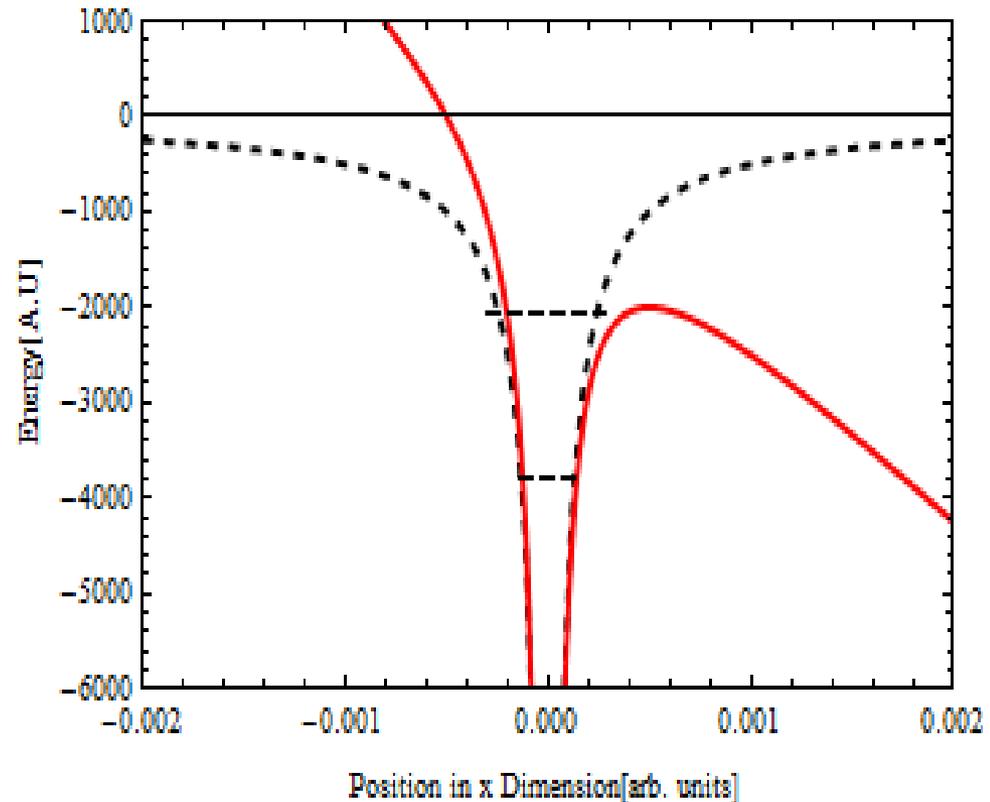
- When do Rydberg atoms ionize?
 - No field applied



$W = -\frac{1}{2(n-\delta_l)^2}$	$\frac{dW}{dn} = \frac{1}{(n-\delta_l)^3}$	$\langle \vec{d} \rangle \approx a_0 n^2$	$\sigma \propto n^4$	$F_{IT} \propto n^{-5}$
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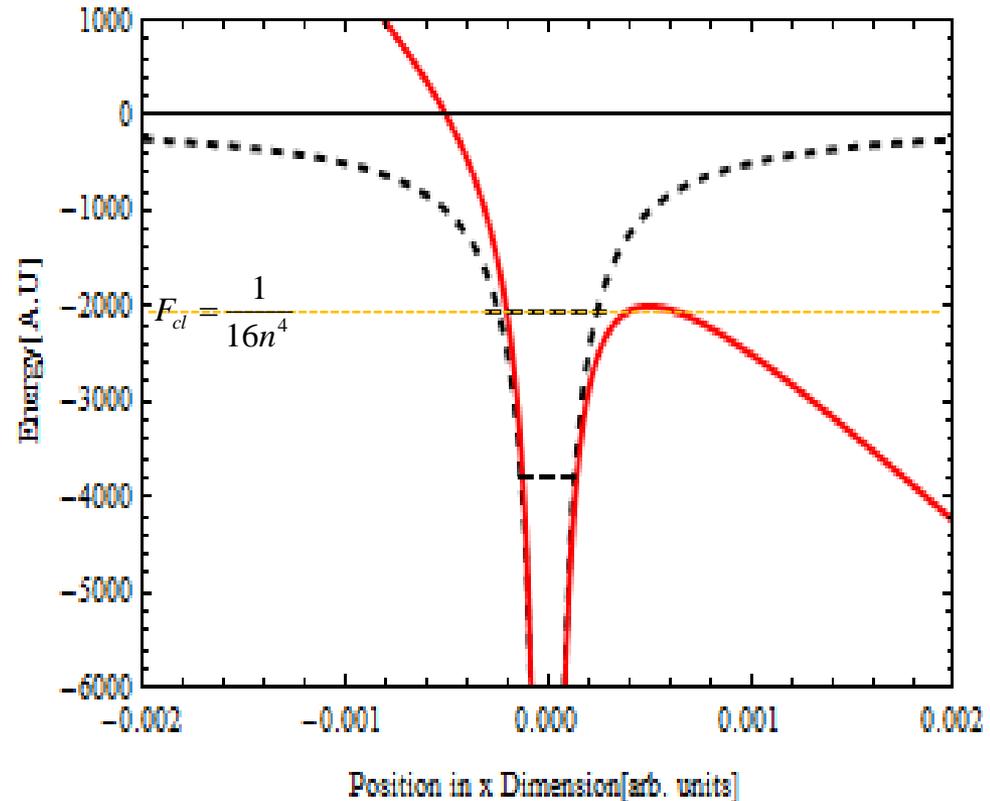
Hydrogen Atom in an electric Field

- When do Rydberg atoms ionize?

- No field applied
- Electric Field applied
- Classical ionization:

$$V = -\frac{1}{r} + Fz$$

$$\Rightarrow F_{cl} = \frac{W^2}{4} = \frac{1}{16n^4}$$



$W = -\frac{1}{2(n-\delta_l)^2}$	$\frac{dW}{dn} = \frac{1}{(n-\delta_l)^3}$	$\langle \vec{d} \rangle \approx a_0 n^2$	$\sigma \propto n^4$	$F_{IT} \propto n^{-5}$	$F_{cl} \propto \frac{1}{16} n^{-4}$
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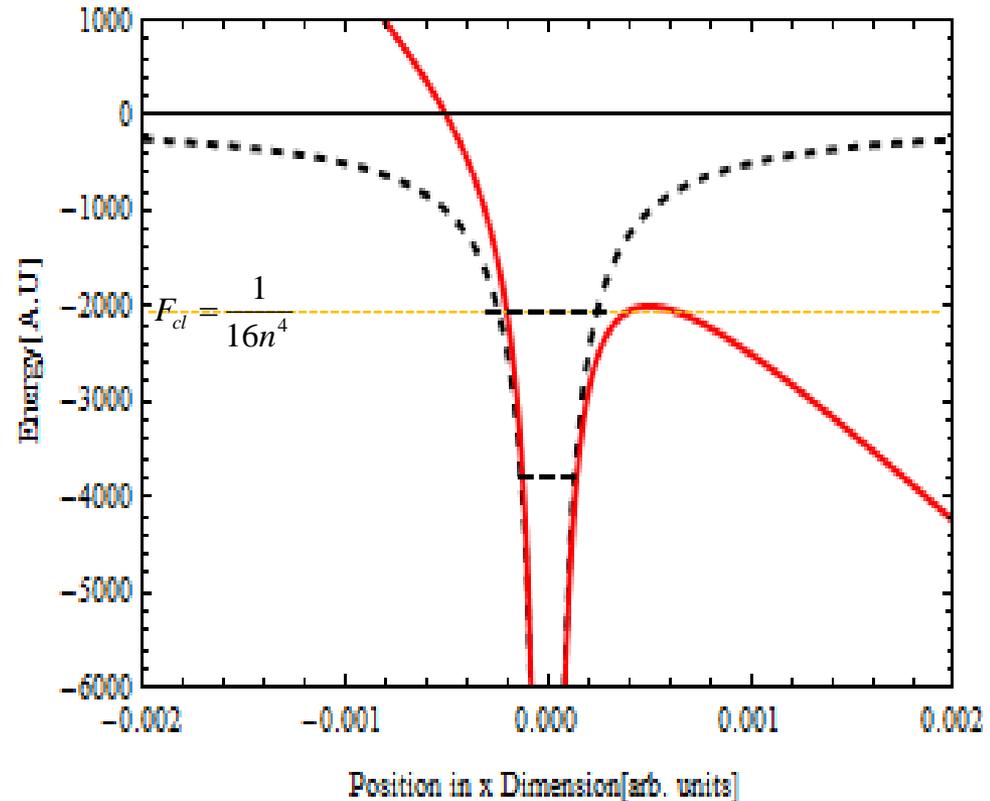
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- Valid only for

- Non-H atoms if F is increased slowly



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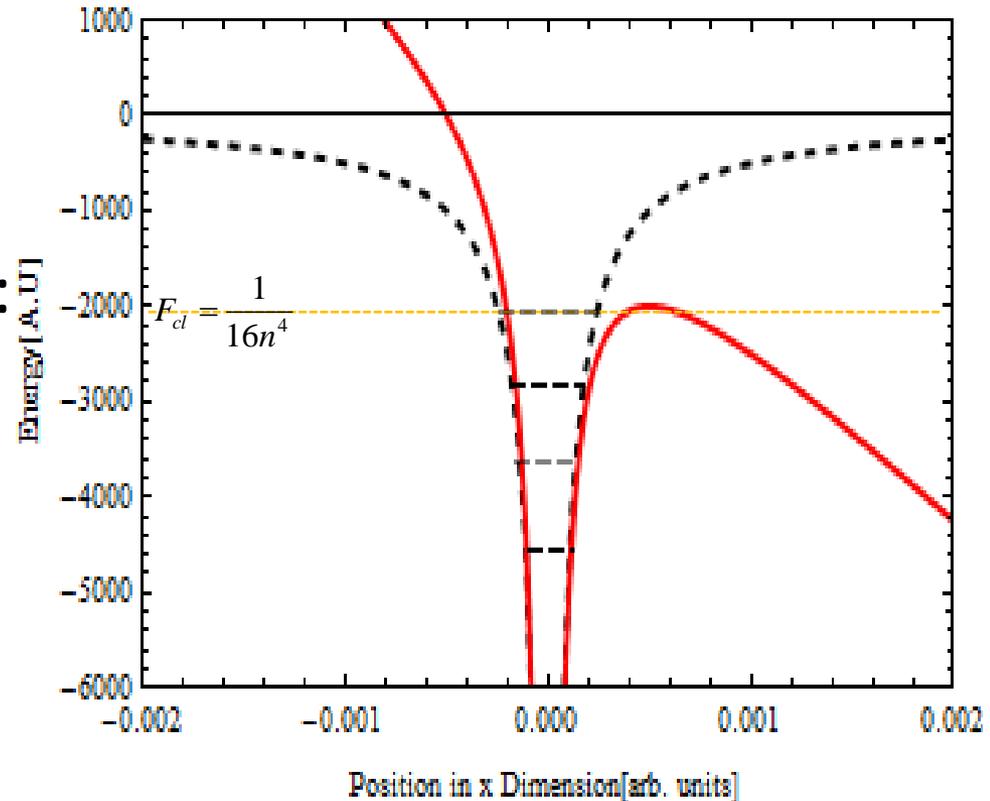
Hydrogen Atom in an electric Field

- When do Rydberg atoms ionize?

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- Electric Field applied
- Quasi-Classical ioniz.:

$$V(\eta) = 2 \left(-\frac{Z_2}{\eta} + \frac{m^2 - 1}{4\eta^2} - \frac{F\eta}{4} \right)$$

$$\Rightarrow F = \frac{W^2}{4Z_2} \Big|_{m=0}$$



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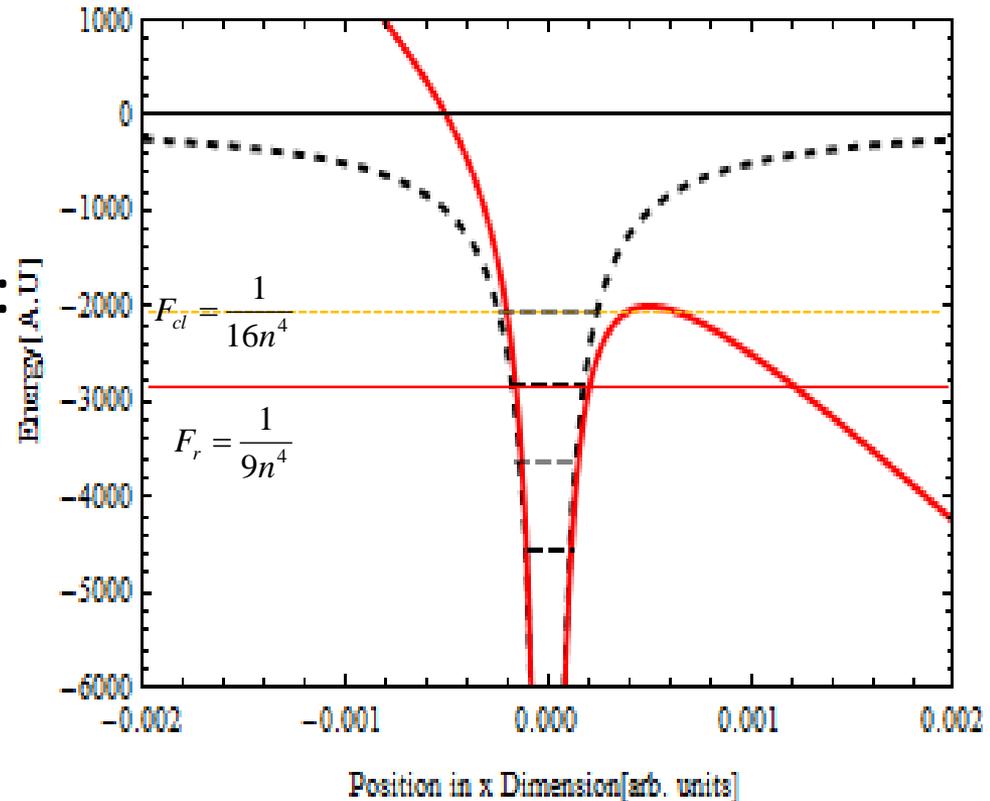
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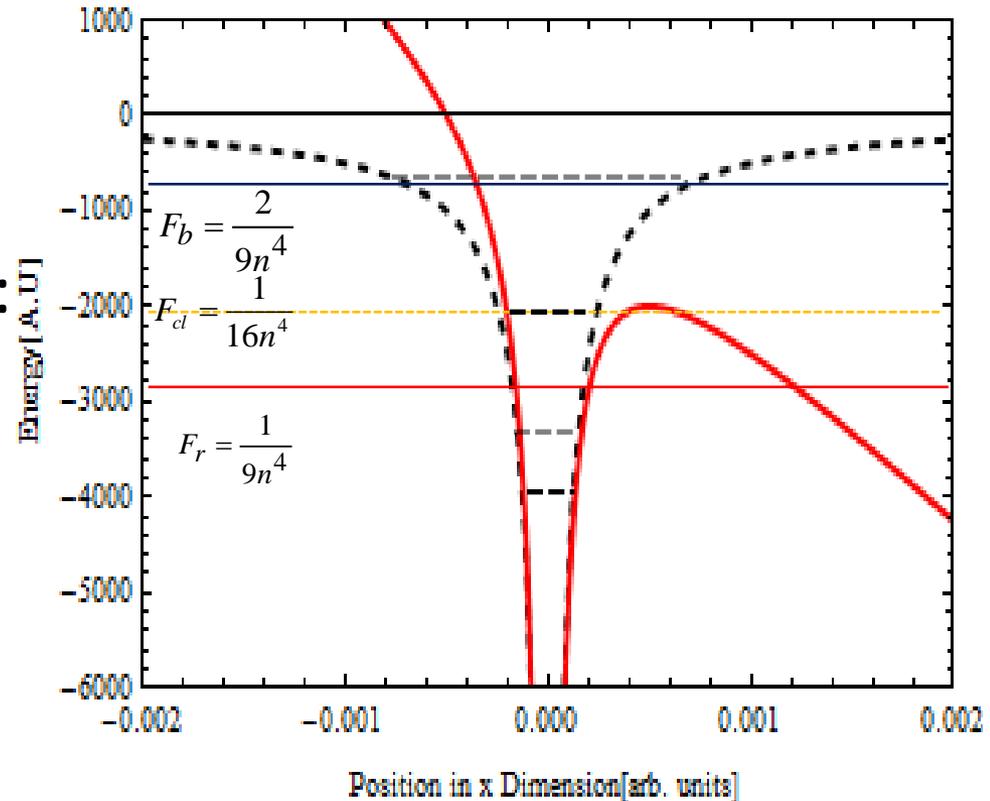
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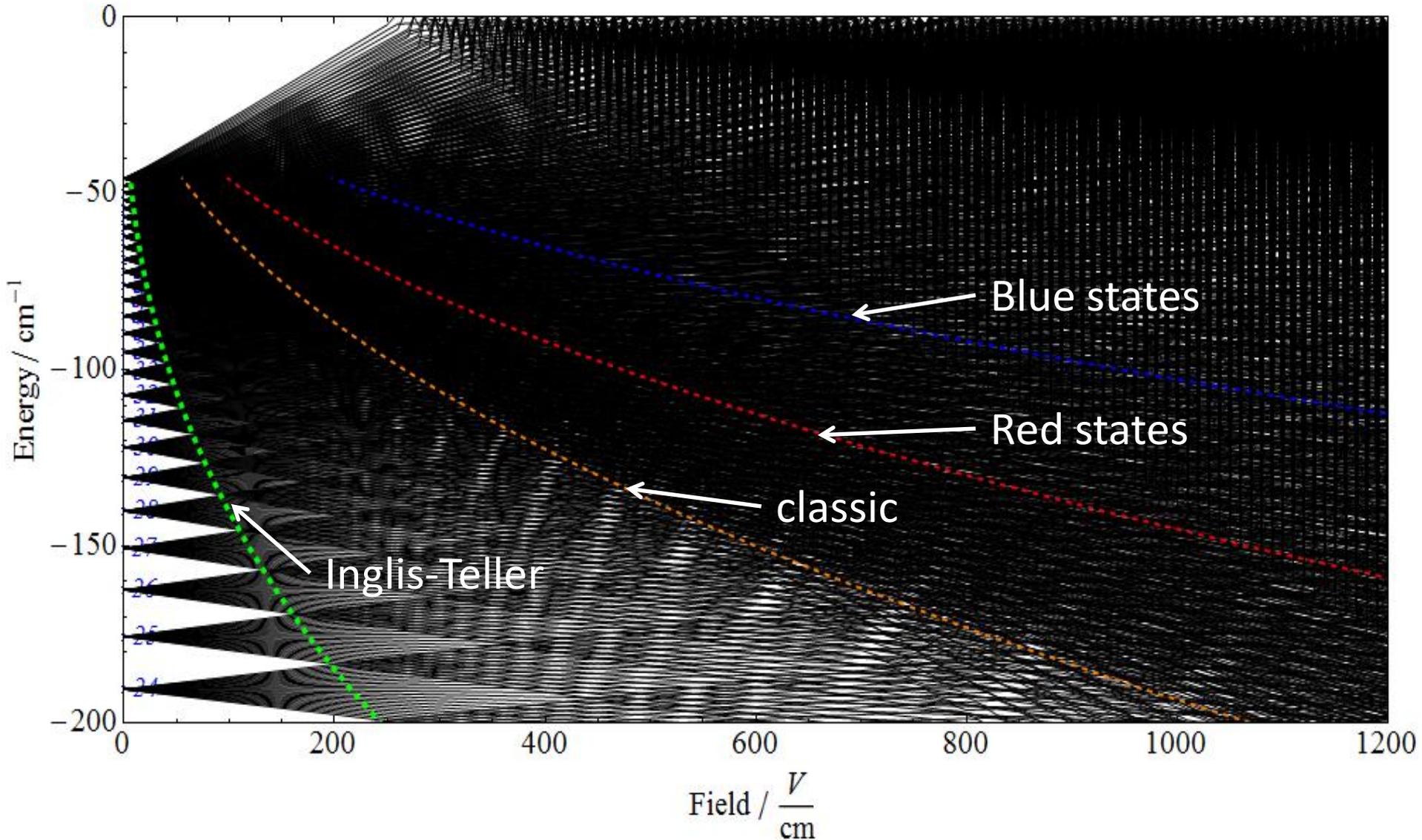
$$V(\eta) = 2 \left(-\frac{Z_2}{\eta} + \frac{m^2 - 1}{4\eta^2} - \frac{F\eta}{4} \right)$$

$$\Rightarrow F = \frac{W^2}{4Z_2} \left| \begin{array}{l} = \frac{1}{9n^4} \quad \text{red} \\ \approx \frac{2}{9n^4} \quad \text{blue} \end{array} \right.$$



$W = -\frac{1}{2(n-\delta_l)^2}$	$\frac{dW}{dn} = \frac{1}{(n-\delta_l)^3}$	$\langle \vec{d} \rangle \approx a_0 n^2$	$\sigma \propto n^4$	$F_{IT} \propto n^{-5}$	$F_{cl} \propto \frac{1}{16} n^{-4}$
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Hydrogen Atom in an electric Field



Lifetime

- From Fermis golden rule
 - Einstein A coefficient for two states $n, l \rightarrow n', l'$

$$A_{n',l',n,l} = \frac{4e^2 \omega_{n',l',n,l}^3}{3\hbar c^3} \frac{\max(l, l')}{2l+1} |\langle n'l' | r | nl \rangle|^2$$

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For $l \approx 0: \propto n^{-3/2}$
Overlap of WF

$W = -\frac{1}{2(n-\delta_l)^2}$	$\frac{dW}{dn} = \frac{1}{(n-\delta_l)^3}$	$\langle \vec{d} \rangle \approx a_0 n^2$	$\sigma \propto n^4$	$F_{IT} \propto n^{-5}$	$F_{cl} \propto \frac{1}{16} n^{-4}$	$\tau_{n,0} \propto n^3$
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Lifetime

- From Fermis golden rule

- Einstein A coefficient for two states $n, l \rightarrow n', l'$

$$A_{n',l',n,l} = \frac{4e^2 \omega_{n',l',n,l}^3}{3\hbar c^3} \frac{\max(l, l')}{2l+1} |\langle n'l' | r | nl \rangle|^2$$

- Lifetime $\tau_{n,l} = \left(\sum_{n',l' < n,l} A_{n',l',n,l} \right)^{-1}$

For $l \approx n$: $\propto n^2$
Overlap of WF

For $l \approx n$: $\propto n^{-3}$
Overlap of WF

$W = -\frac{1}{2(n-\delta_l)^2}$	$\frac{dW}{dn} = \frac{1}{(n-\delta_l)^3}$	$\langle \vec{d} \rangle \approx a_0 n^2$	$\sigma \propto n^4$	$F_{IT} \propto n^{-5}$	$F_{cl} \propto \frac{1}{16} n^{-4}$	$\tau_{n,l} \propto n^3, n^5$
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Lifetime

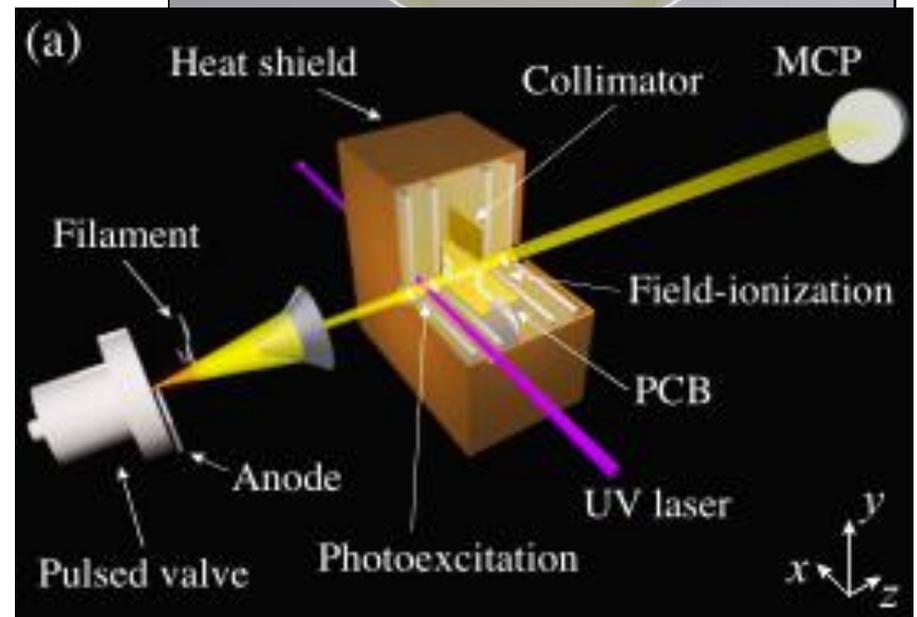
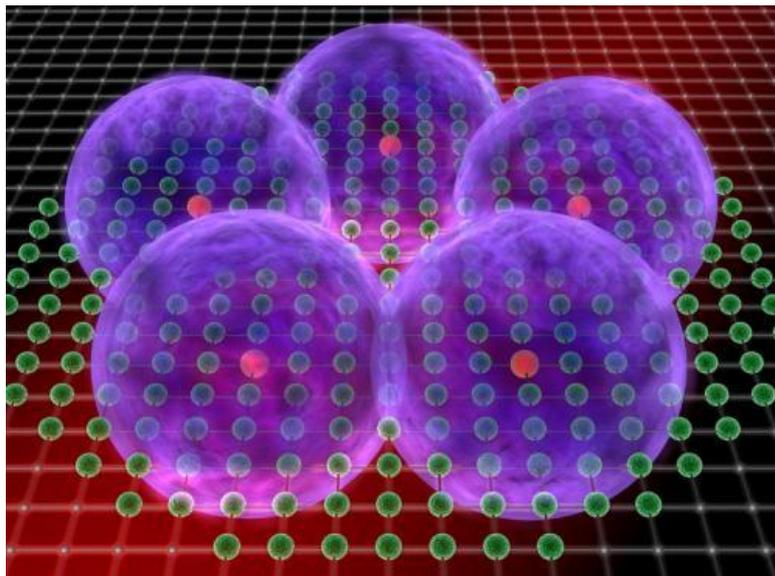
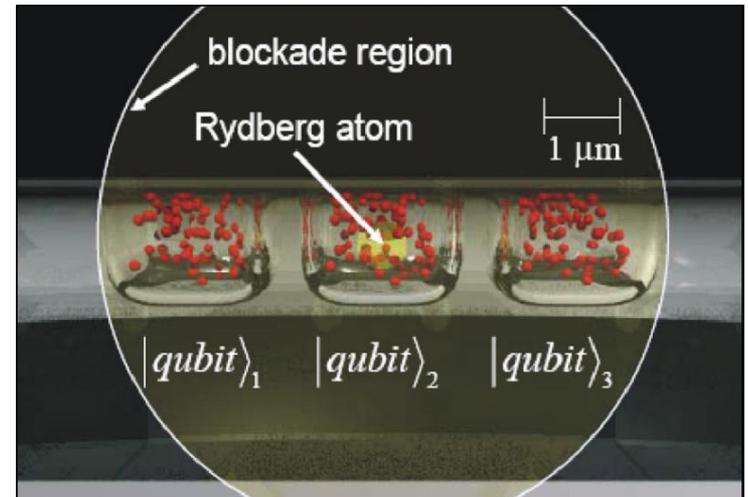
$$A_{n',l',n,l} = \frac{4e^2 \omega_{n',l',n,l}^3}{3\hbar c^3} \frac{\max(l,l')}{2l+1} \left| \langle n'l' | r | nl \rangle \right|^2 \quad \tau_{n,l} = \left(\sum_{n',l' < n,l} A_{n',l',n,l} \right)^{-1}$$

State	Stark State 60 p (n',l') small	Circular state 60 l=59 m=59 (n',l') \approx ($n \pm 1, l \pm 1$)	Statistical mixture
Scaling	n^3 (overlap of $\psi \propto n^{-3/2}$)	n^5 $\langle r \rangle \propto n^2$	$n^{4.5}$
Lifetime	7.2 μ s	70 ms	\approx ms

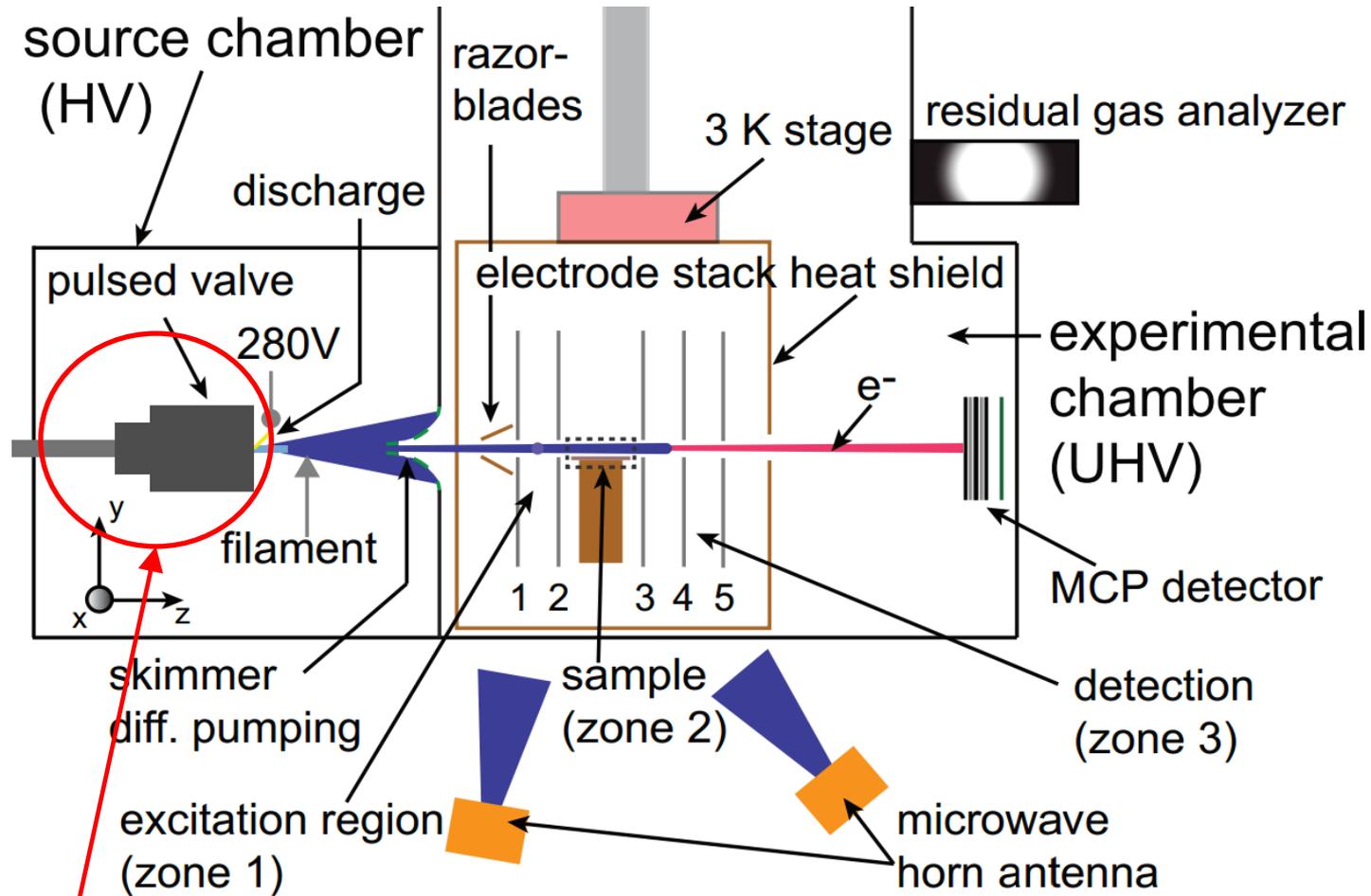
$W = -\frac{1}{2(n-\delta_l)^2}$	$\frac{dW}{dn} = \frac{1}{(n-\delta_l)^3}$	$\langle \vec{d} \rangle \approx a_0 n^2$	$\sigma \propto n^4$	$F_{IT} \propto n^{-5}$	$F_{cl} \propto \frac{1}{16} n^{-4}$	$\tau_{n,l} \propto n^{-3}, n^{-5}$
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Part 2- Generation of Rydberg atoms

- Typical Experiments:
 - Beam experiments
 - (ultra) cold atoms
 - Vapor cells

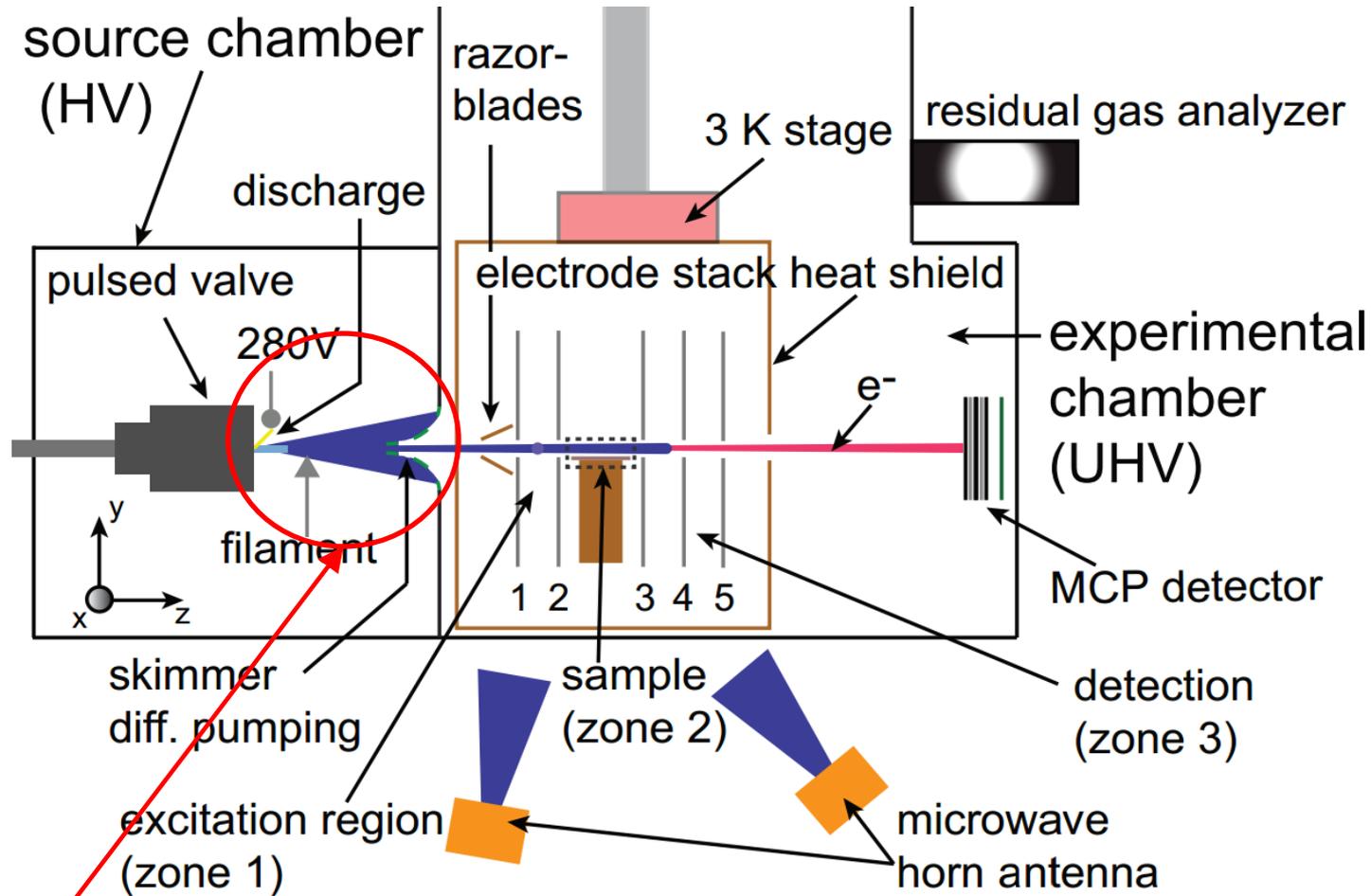


ETH physics Rydberg experiment



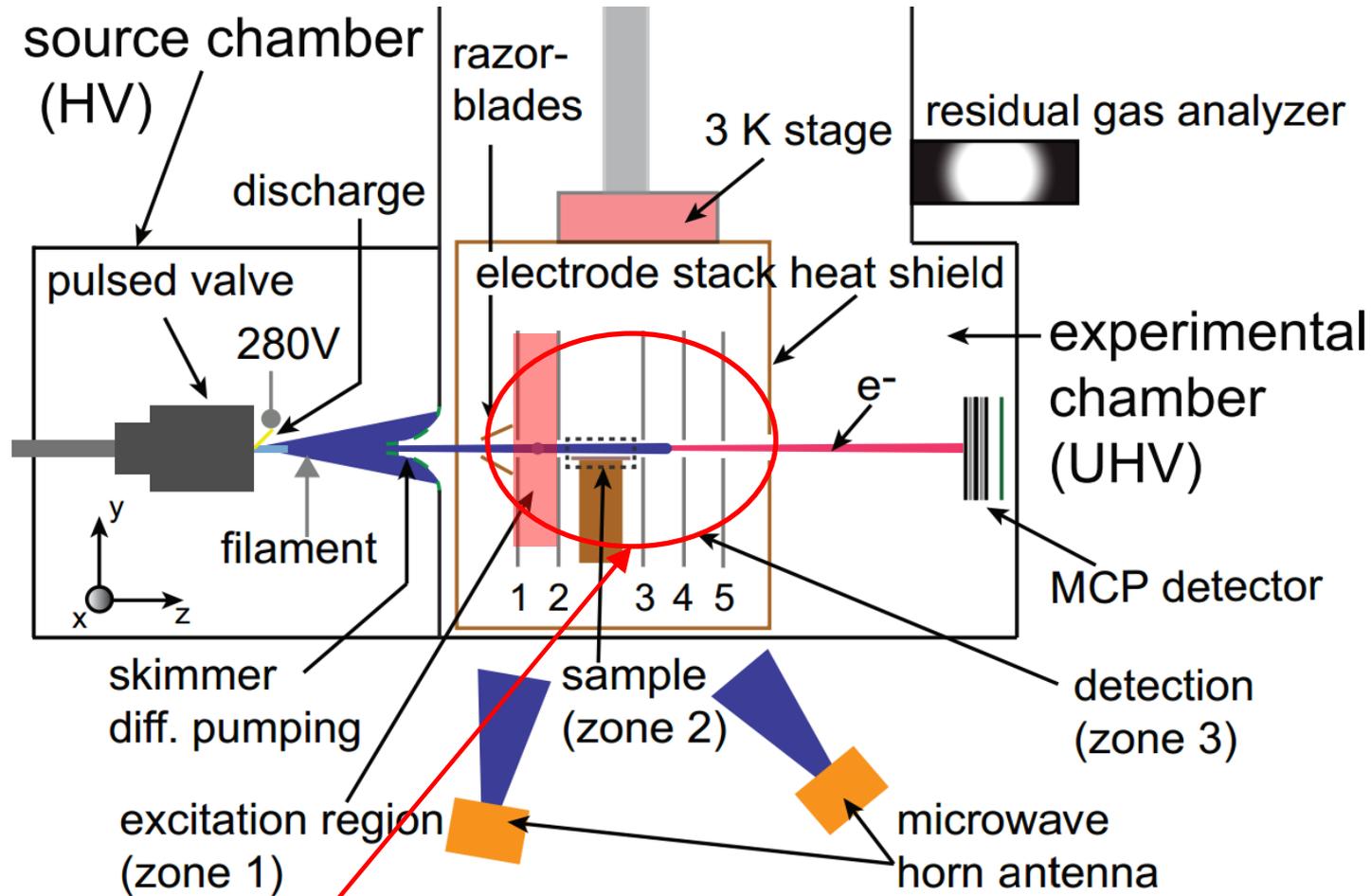
Creation of a cold supersonic beam of Helium.
Speed: 1700m/s, pulsed: 25Hz, temperature atoms=100mK

ETH physics Rydberg experiment



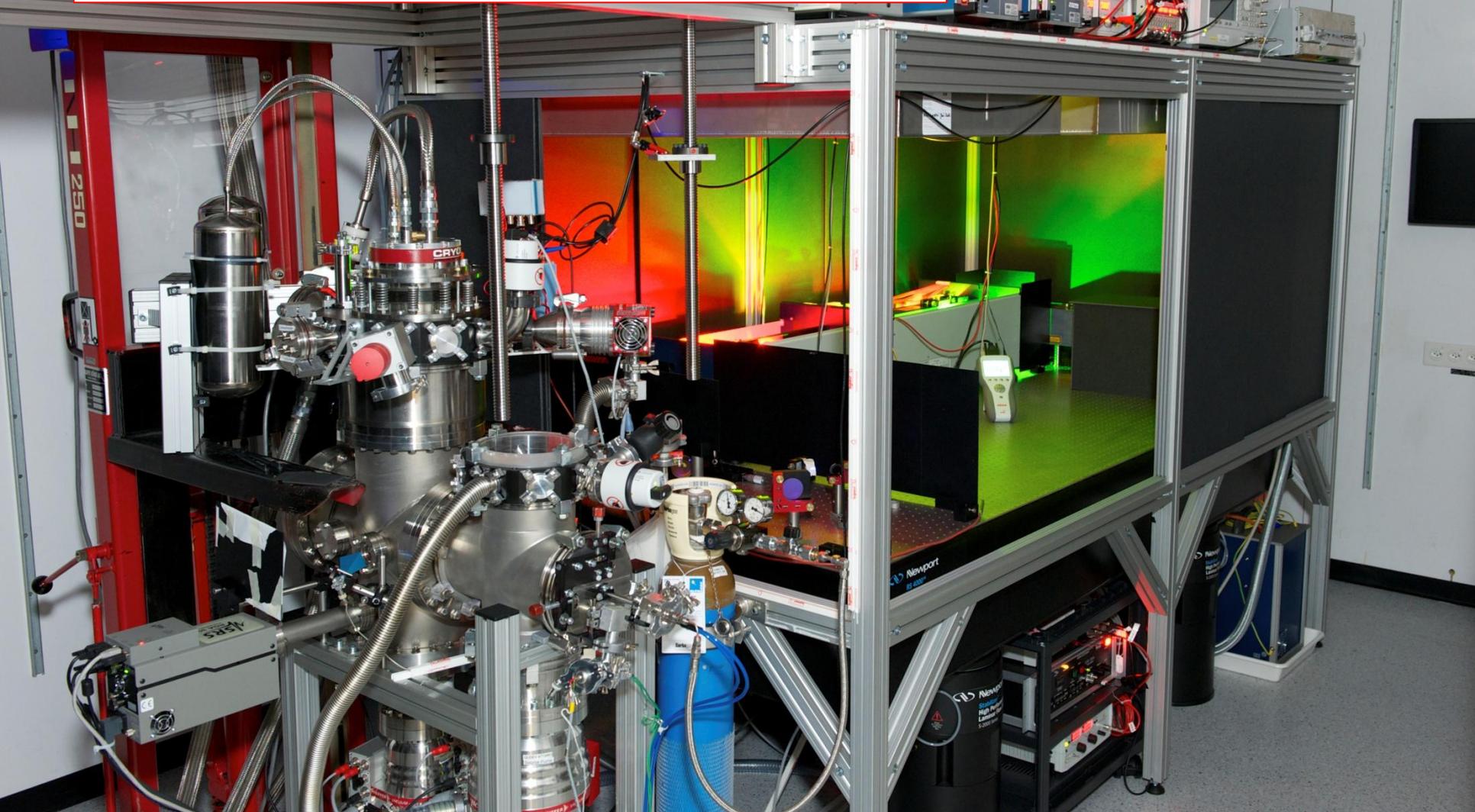
Excite electrons to the 2s-state, (to overcome very strong binding energy in the xuv range) by means of a discharge – like a lightning.

ETH physics Rydberg experiment



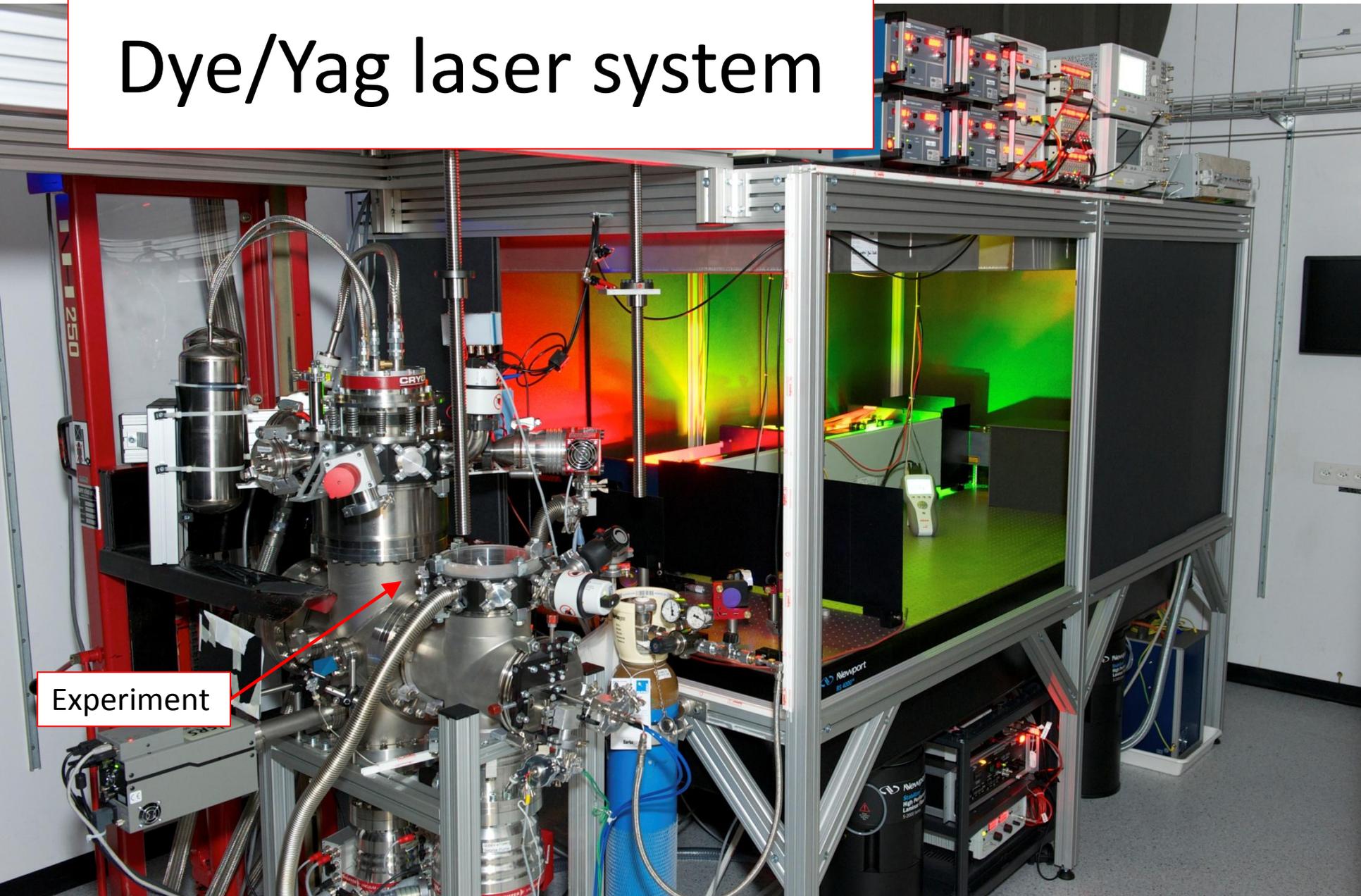
Actual experiment consists of 5 electrodes. Between the first 2 the atoms get excited to Rydberg states up to $n=\infty$ with a dye laser.

Dye/Yag laser system



Dye/Yag laser system

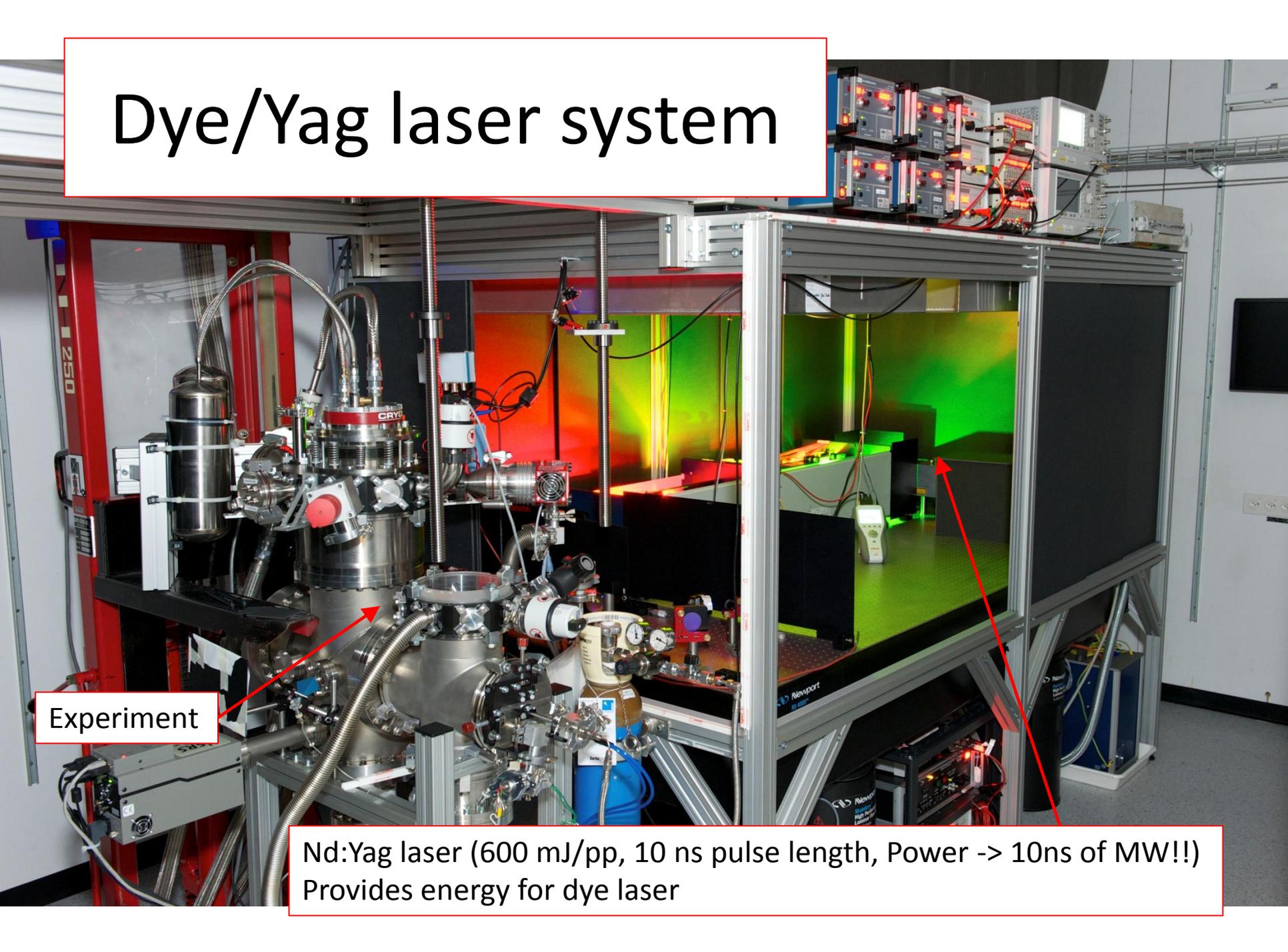
Experiment



Dye/Yag laser system

Experiment

Nd:Yag laser (600 mJ/pp, 10 ns pulse length, Power \rightarrow 10ns of MW!!)
Provides energy for dye laser

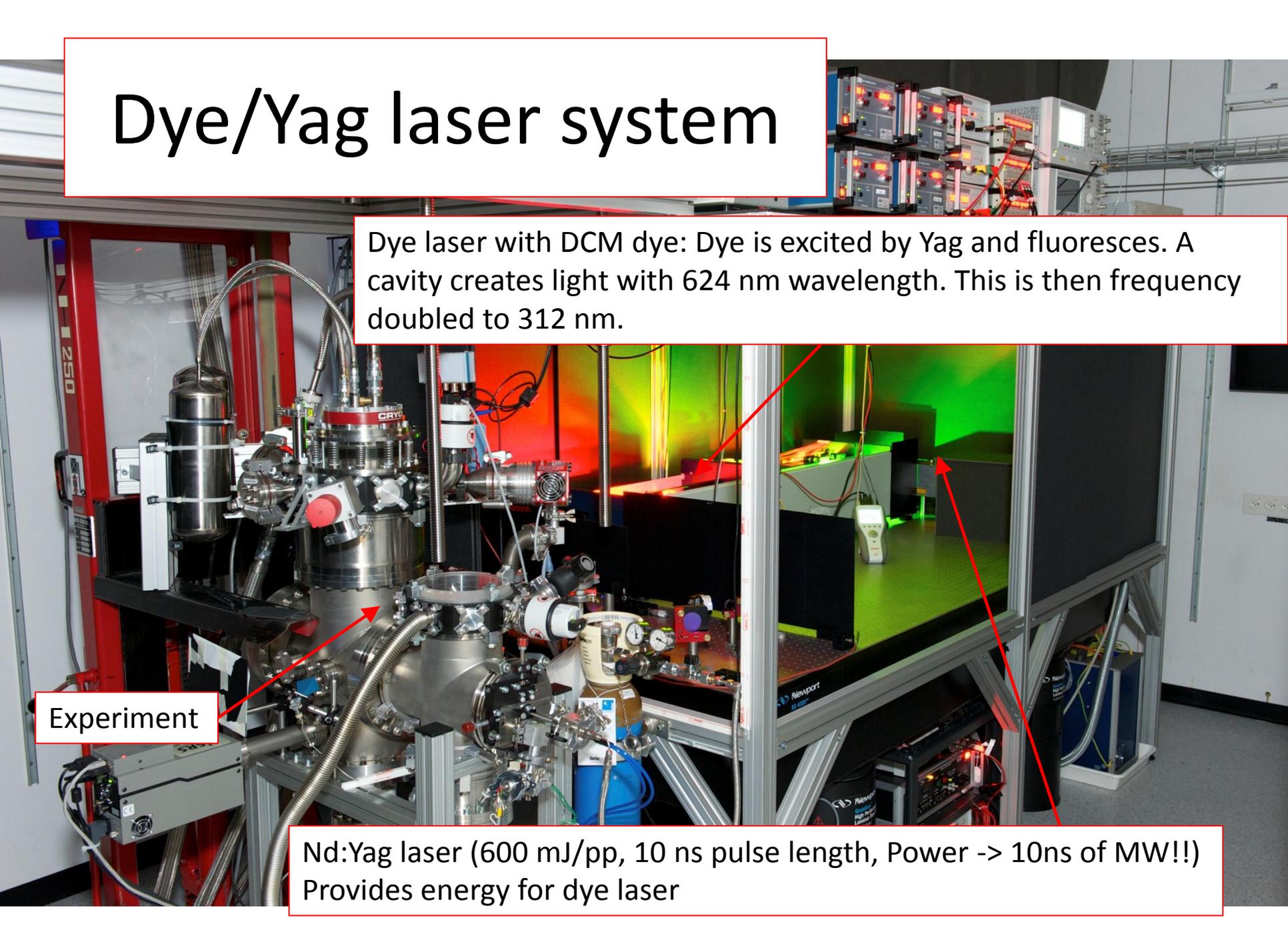


Dye/Yag laser system

Dye laser with DCM dye: Dye is excited by Yag and fluoresces. A cavity creates light with 624 nm wavelength. This is then frequency doubled to 312 nm.

Experiment

Nd:Yag laser (600 mJ/pp, 10 ns pulse length, Power -> 10ns of MW!!)
Provides energy for dye laser



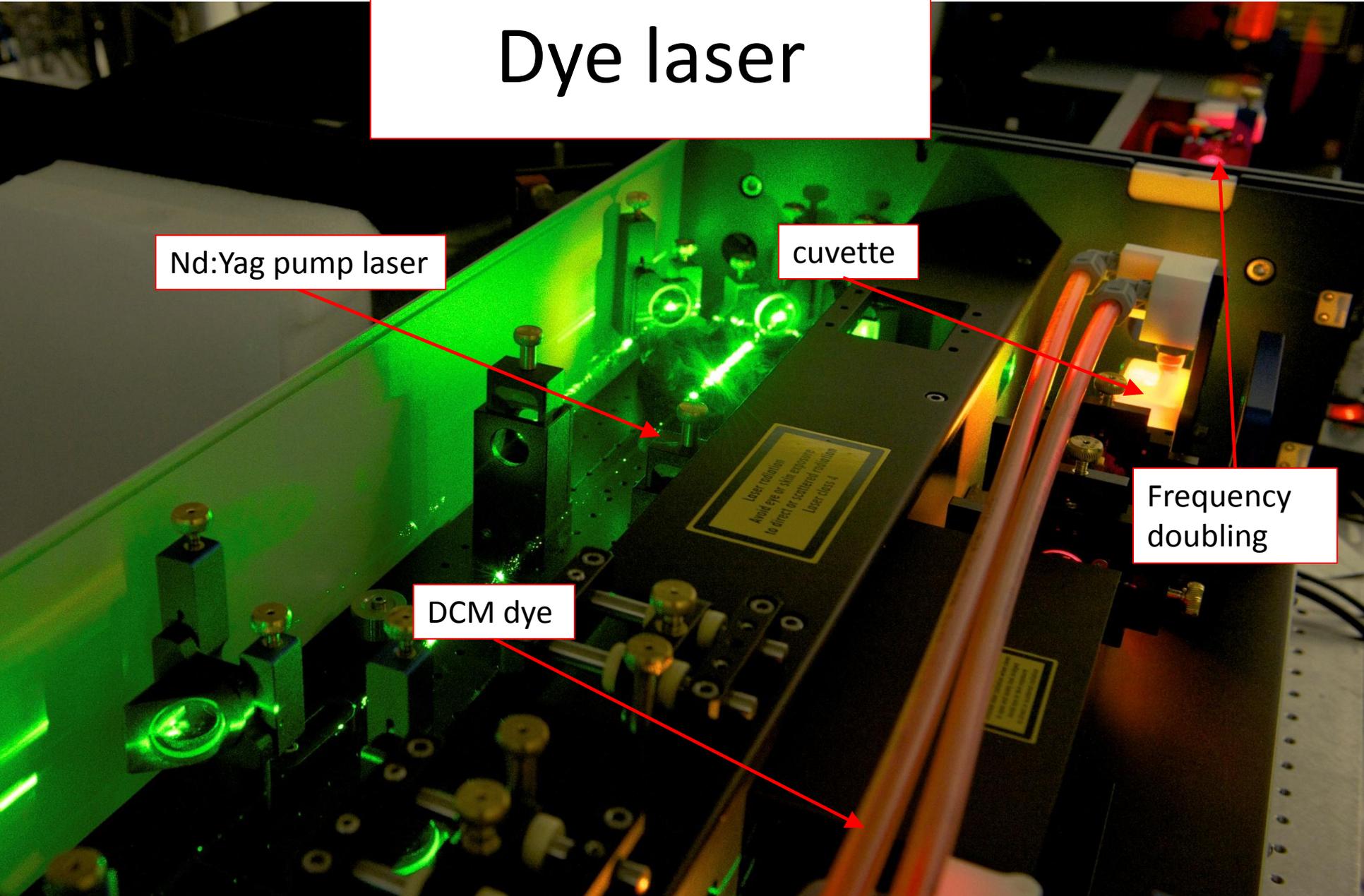
Dye laser

Nd:Yag pump laser

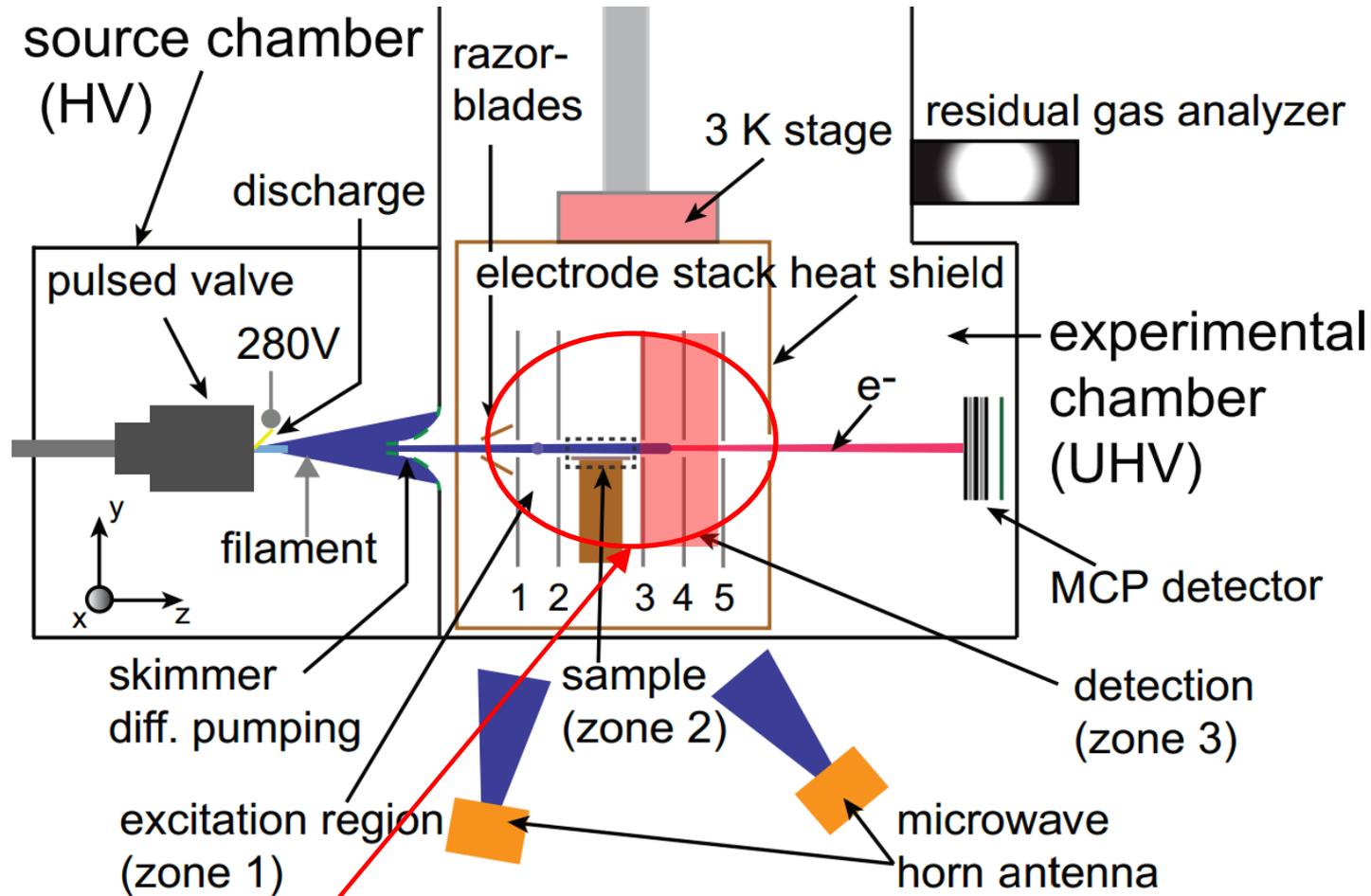
cuvette

DCM dye

Frequency doubling

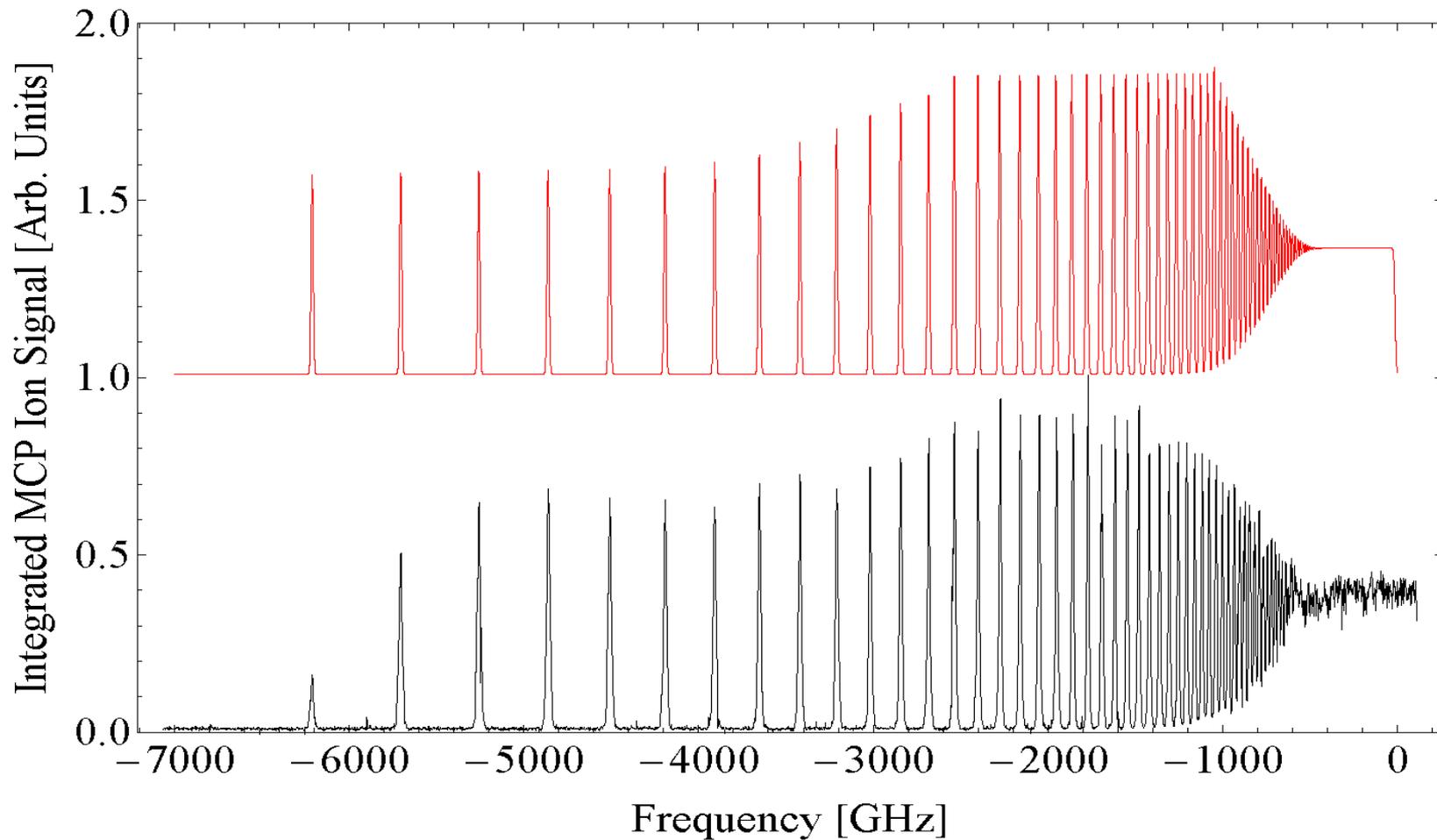


ETH physics Rydberg experiment

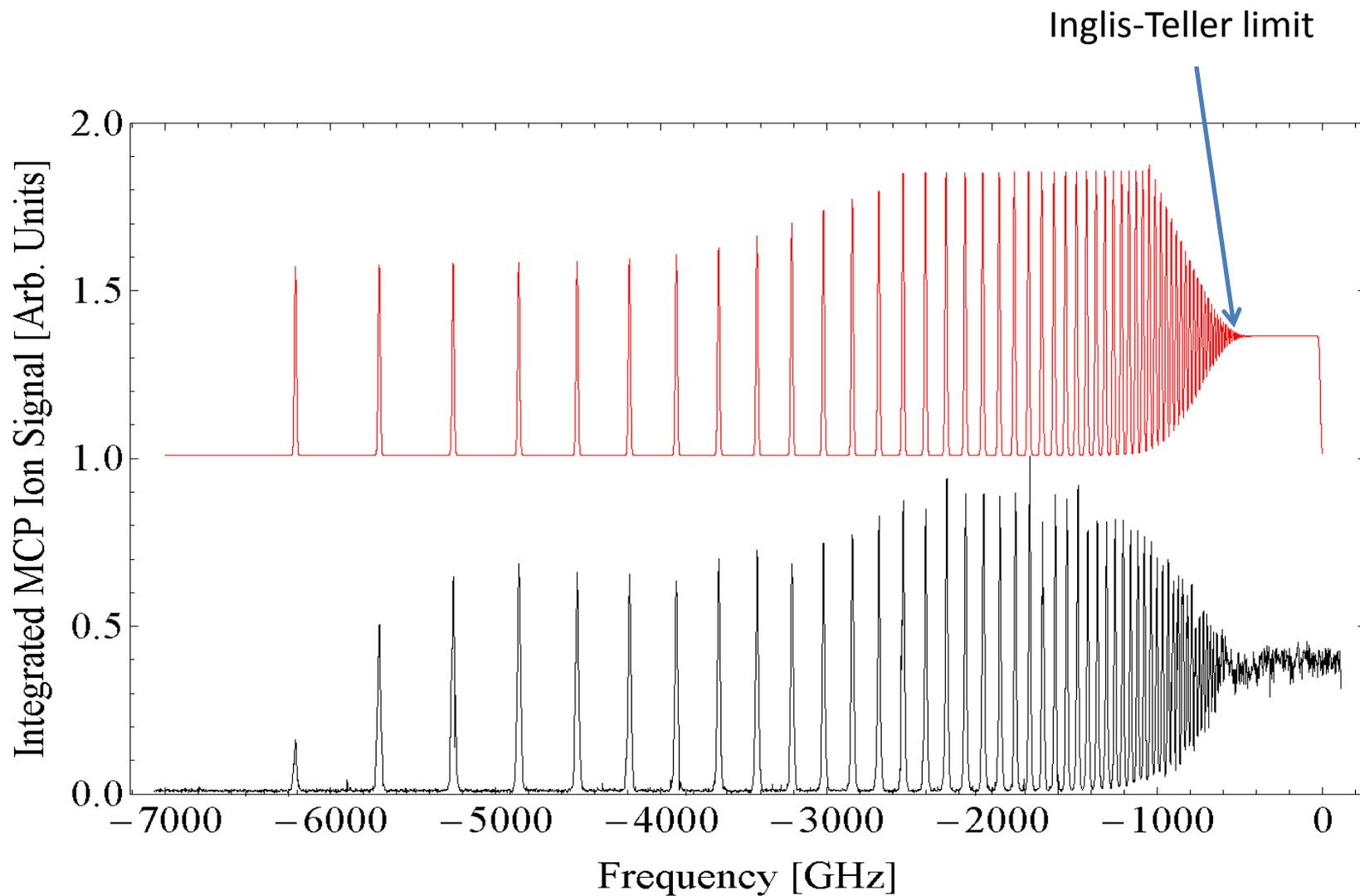


Detection: 1.2 kV/cm electric field applied in 10 ns. Rydberg atoms ionize and electrons are Detected at the MCP detector (single particle multiplier) .

Results TOF 100ns

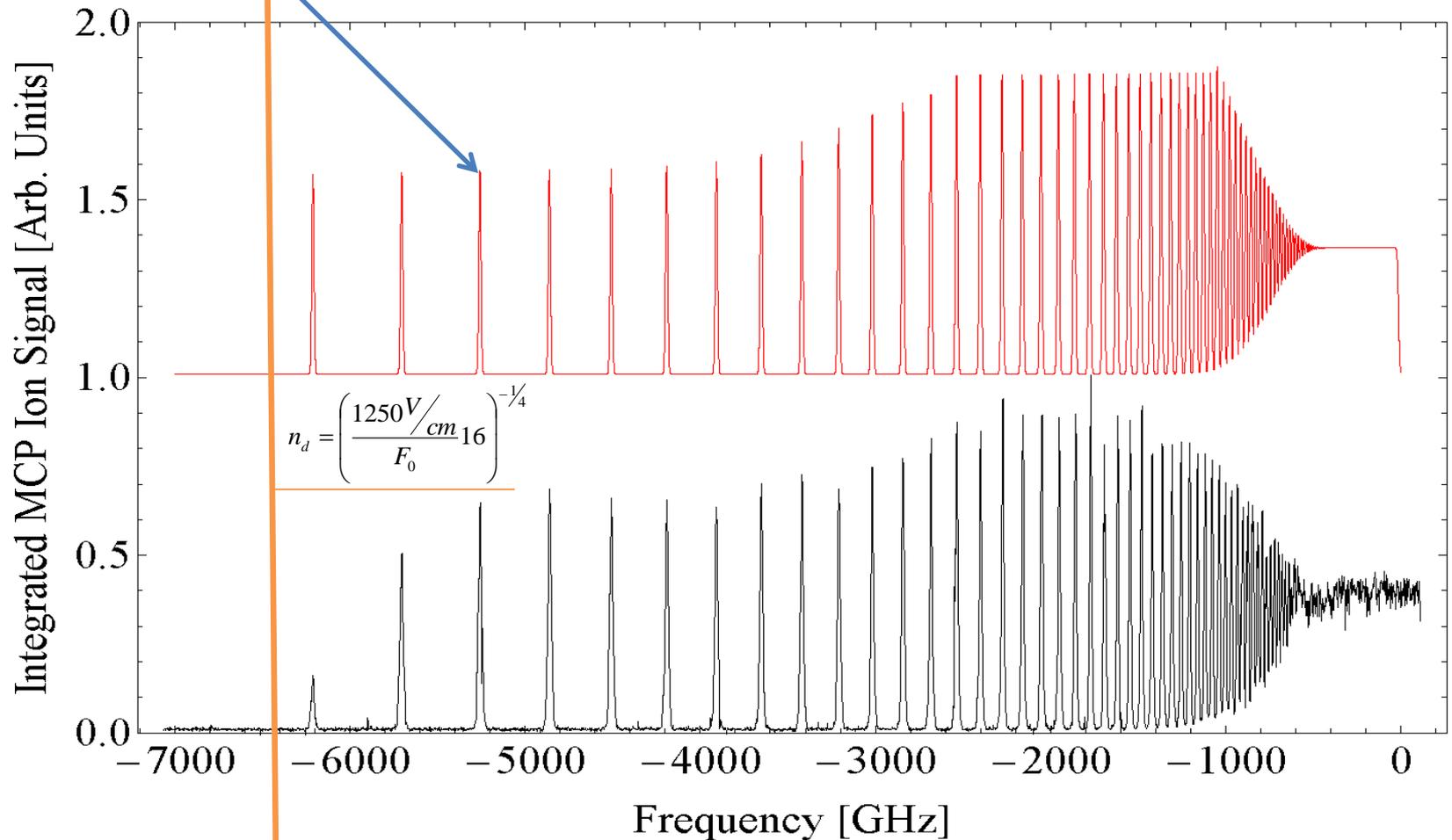


Results TOF 100ns

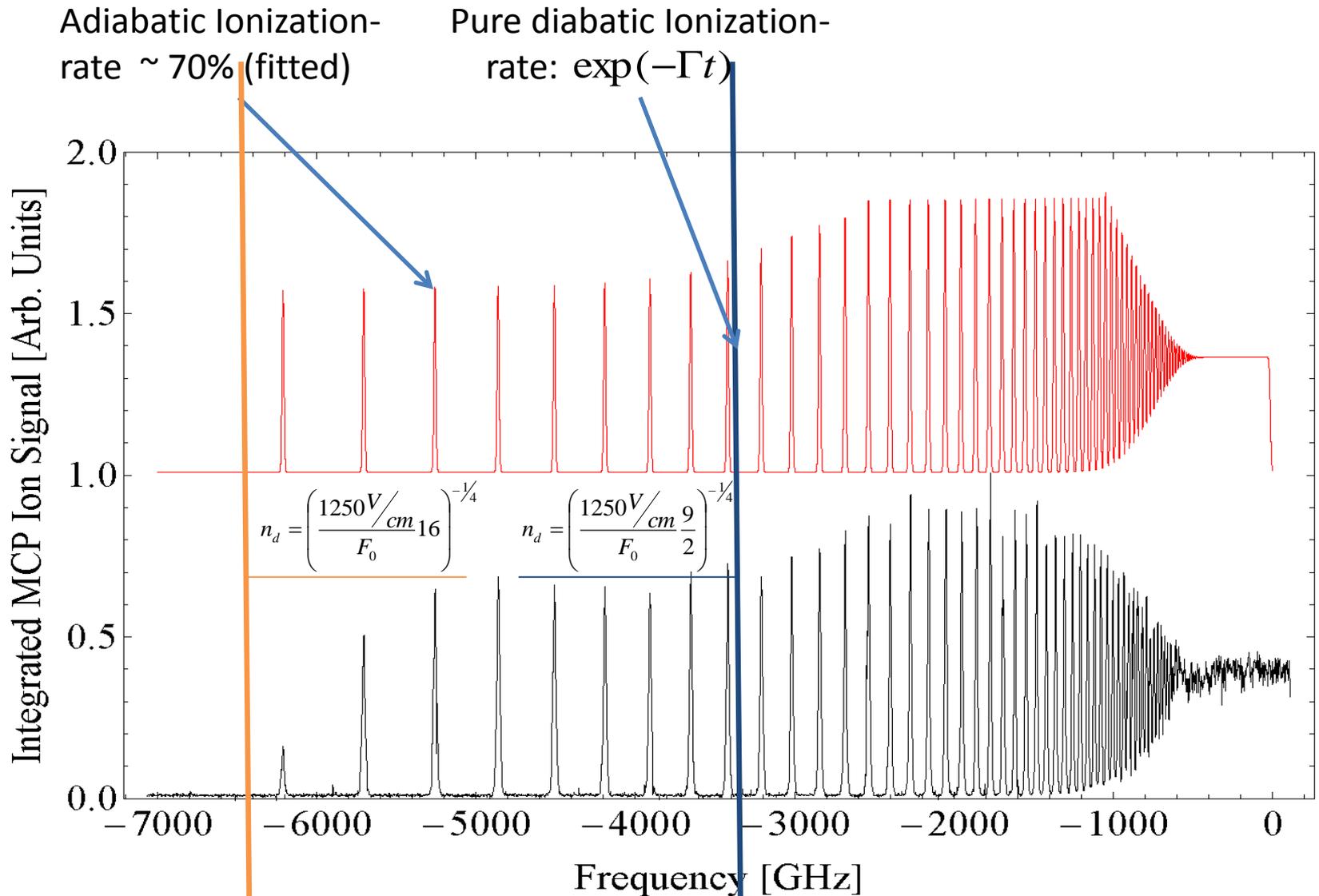


Results TOF 100ns

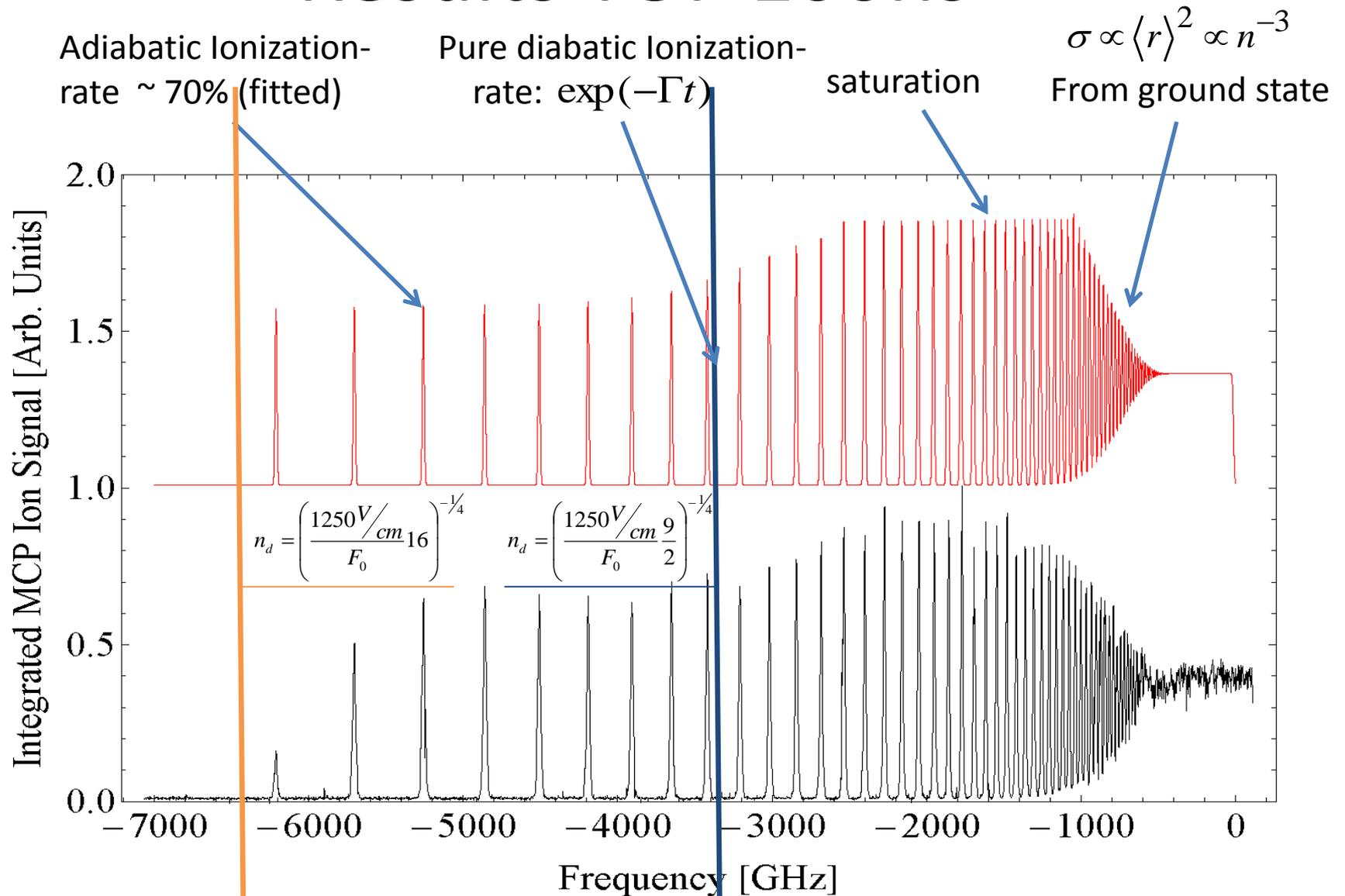
Adiabatic Ionization-
rate $\sim 70\%$ (fitted)



Results TOF 100ns



Results TOF 100ns



Results TOF $15\mu\text{s}$

$\sigma \propto \langle r \rangle^2 \propto n^{-3}$
From ground state

