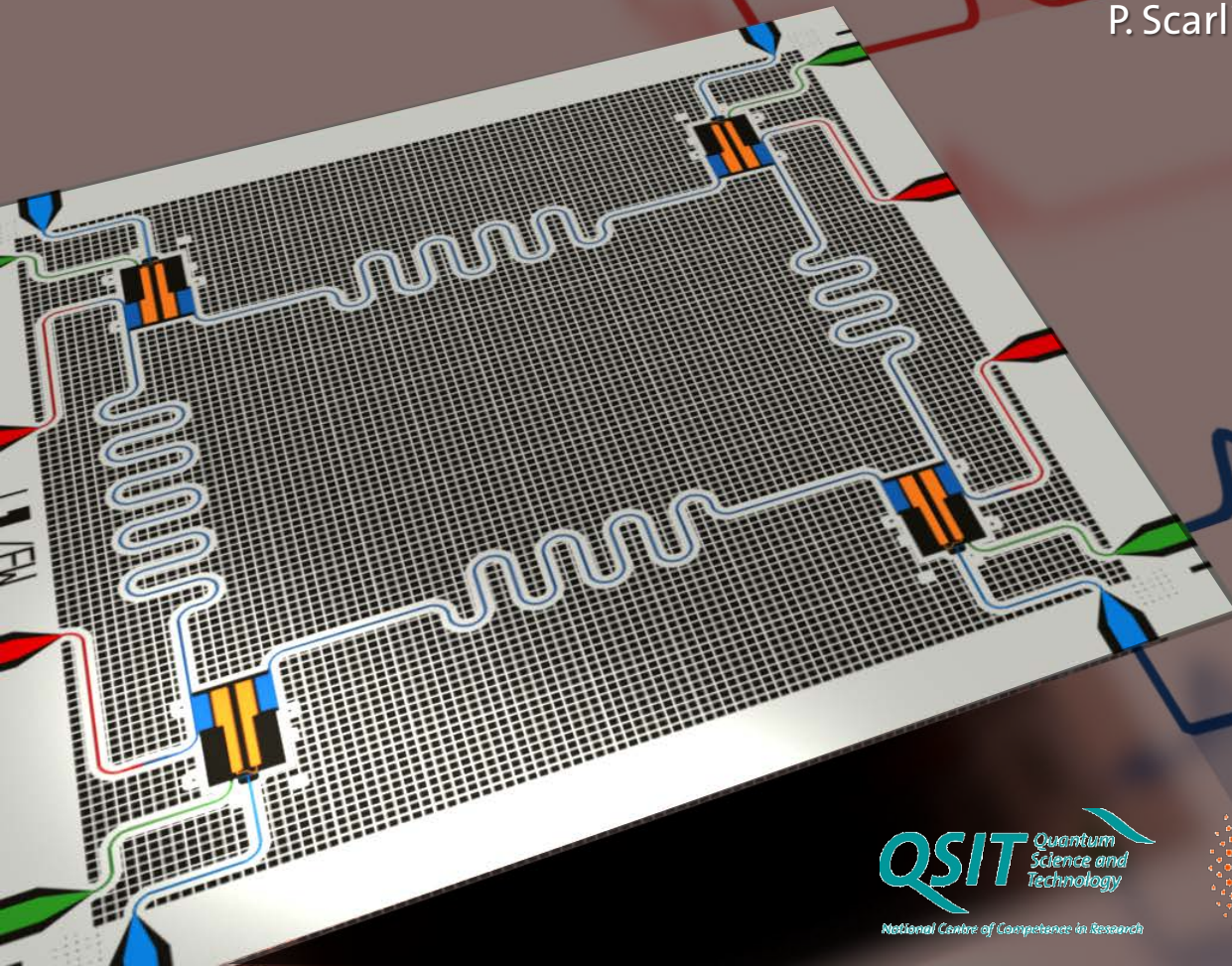


Exploring Quantum Physics with Superconducting Circuits

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CIRCUIT AND CAVITY
QUANTUM ELECTRODYNAMICS



Conventional Electronic Circuits

basic circuit elements:

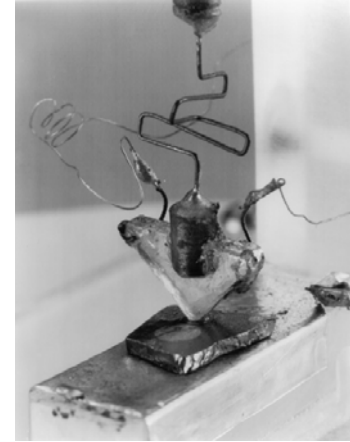


basis of modern
information and
communication
technology

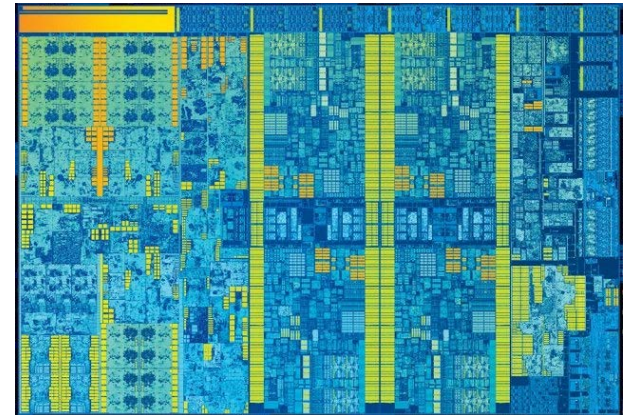
properties :

- classical physics
- no quantum mechanics
- no superposition or entanglement
- no quantization of fields

first transistor at Bell Labs (1947)



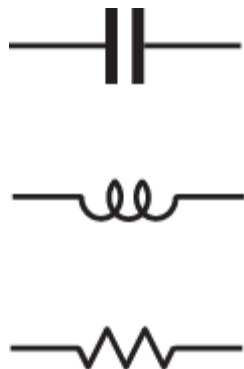
Intel Core i7-6700K Processor



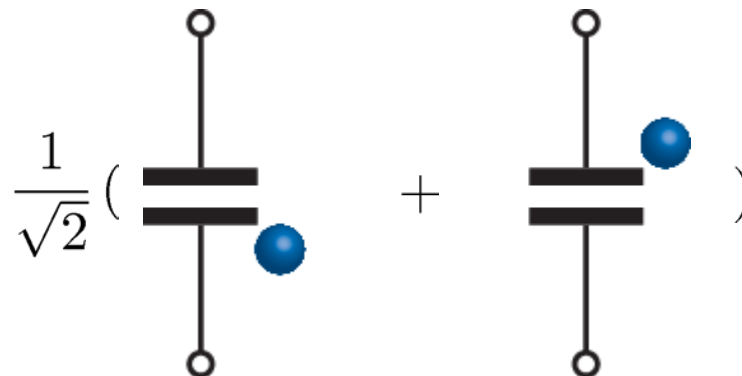
smallest feature size 14 nm
clock speed ~ 4.2 GHz
> $3 \cdot 10^9$ transistors
power consumption > 10 W

Classical and Quantum Electronic Circuit Elements

basic circuit elements:



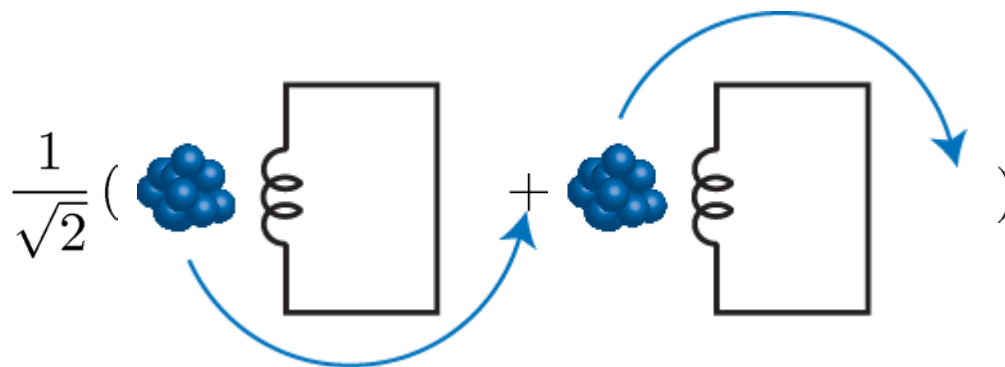
charge on a capacitor:



quantum superposition states of:

- charge q
- flux ϕ

current or magnetic flux in an inductor:

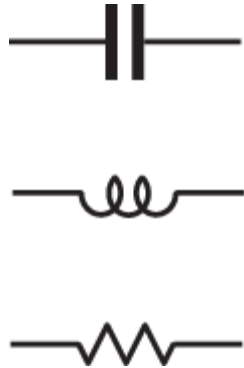


commutation relation (c.f. x, p):

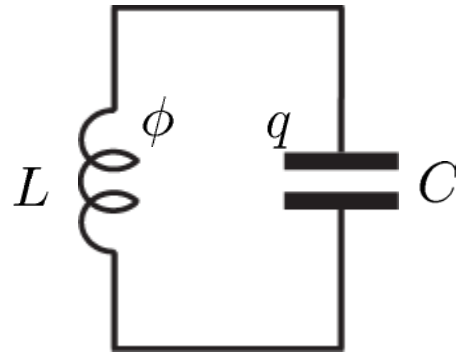
$$[\hat{\phi}, \hat{q}] = i\hbar$$

Constructing Linear Quantum Electronic Circuits

basic circuit elements:



harmonic LC oscillator:



$$\omega = \frac{1}{\sqrt{LC}} \sim 5 \text{ GHz}$$

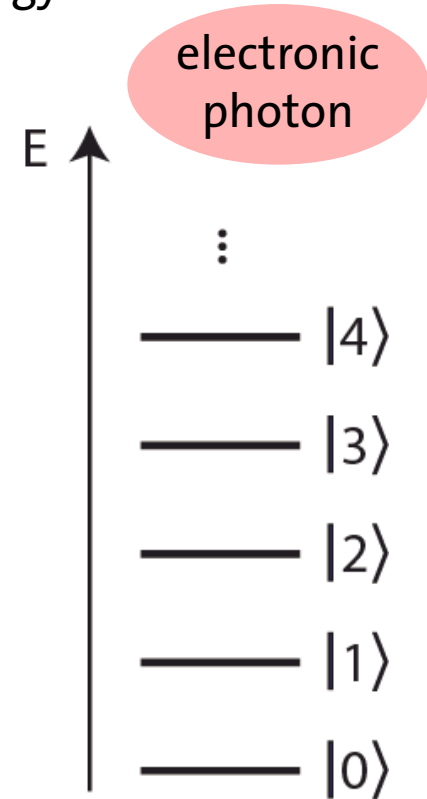
classical physics:

$$H = \frac{\phi^2}{2L} + \frac{q^2}{2C}$$

quantum mechanics:

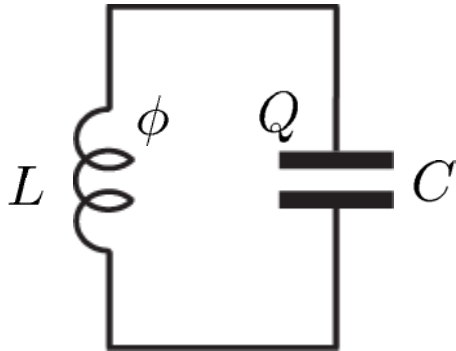
$$\hat{H} = \frac{\hat{\phi}^2}{2L} + \frac{\hat{q}^2}{2C} = \hbar\omega(\hat{a}^\dagger\hat{a} + \frac{1}{2}) \quad [\hat{\phi}, \hat{q}] = i\hbar$$

energy:



Quantization of an Electronic Harmonic Oscillator

Harmonic LC oscillator:



$$Q = CV$$

Charge on capacitor

$$\phi = LI$$

Flux in inductor

$$V = -L\dot{I} = -\dot{\phi}$$

Voltage across inductor

Classical Hamiltonian:

$$H = \frac{CV^2}{2} + \frac{LI^2}{2} = \frac{Q^2}{2C} + \frac{\phi^2}{2L}$$

Conjugate variables:

$$\frac{\partial H}{\partial \phi} = \frac{\phi}{L} = I = \dot{Q}$$

$$\frac{\partial H}{\partial Q} = \frac{Q}{C} = V = -L\dot{I} = -\dot{\phi}$$

Hamilton operator:

$$\hat{H} = \frac{\hat{\phi}^2}{2L} + \frac{\hat{q}^2}{2C}$$

Flux and charge operator:

$$\hat{\phi} = \phi$$

$$\hat{Q} = -i\hbar \frac{\partial}{\partial \phi}$$

Commutation relation:

$$[\hat{\phi}, \hat{q}] = i\hbar$$

Creation and Annihilation Operators for Circuits

Hamilton operator of harmonic oscillator in second quantization:

$$\hat{H} = \frac{\hat{\phi}^2}{2L} + \frac{\hat{Q}^2}{2C} = \hbar\omega(\hat{a}^\dagger\hat{a} + 1/2)$$

$$\hat{a}^\dagger |n\rangle = \sqrt{n+1} |n+1\rangle \quad \text{Creation operator}$$

$$\hat{a} |n\rangle = \sqrt{n} |n-1\rangle \quad \text{Annihilation operator}$$

$$\hat{a}^\dagger\hat{a} |n\rangle = n |n\rangle \quad \text{Number operator}$$

$$\hat{Q} = \sqrt{\frac{\hbar}{2Z_C}}(\hat{a}^\dagger + \hat{a}) \quad \text{Charge/voltage operator}$$

$$\hat{V} = \frac{\hat{Q}}{C}$$

$$\hat{\phi} = i\sqrt{\frac{\hbar Z_C}{2}}(\hat{a}^\dagger - \hat{a}) \quad \text{Flux/current operator}$$

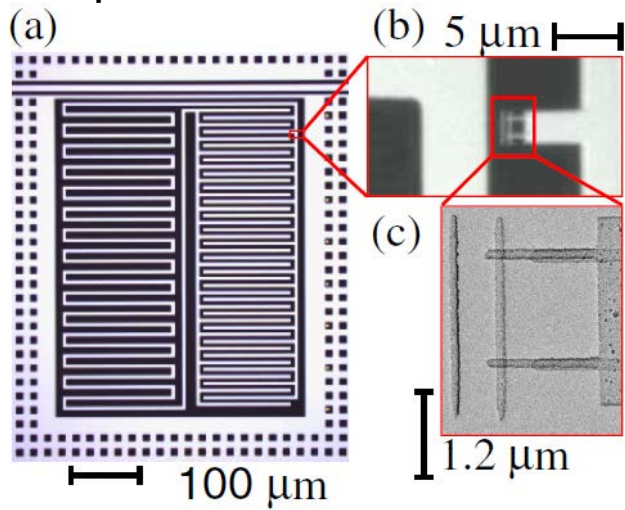
$$\hat{I} = \frac{\hat{\phi}}{L}$$

With characteristic impedance:

$$Z_C = \sqrt{\frac{L}{C}}$$

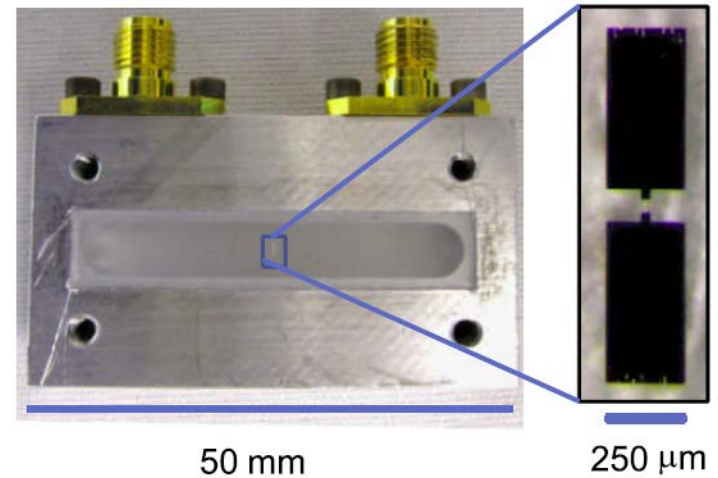
Flavors of Superconducting Harmonic Oscillators

lumped element resonator:



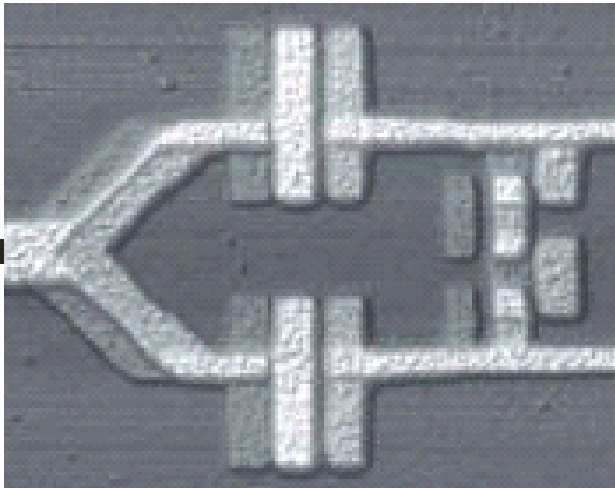
Z. Kim *et al.*, *PRL* 106, 120501 (2011)

3D cavity:



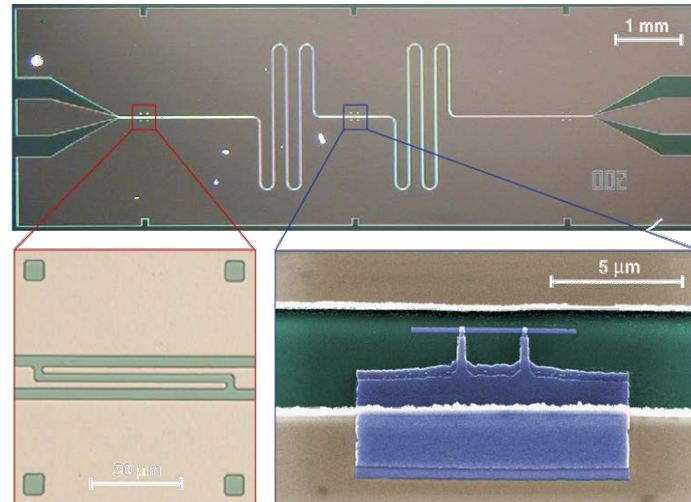
H. Paik *et al.*, *PRL* 107, 240501 (2011)

weakly nonlinear junction:



I. Chiorescu *et al.*, *Nature* 431, 159 (2004)

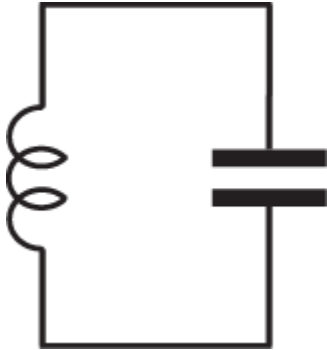
planar transmission line resonator:



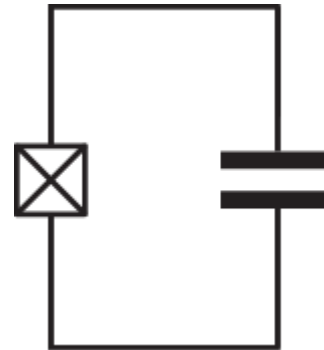
A. Wallraff *et al.*, *Nature* 431, 162 (2004)

Linear vs. Nonlinear Superconducting Oscillators

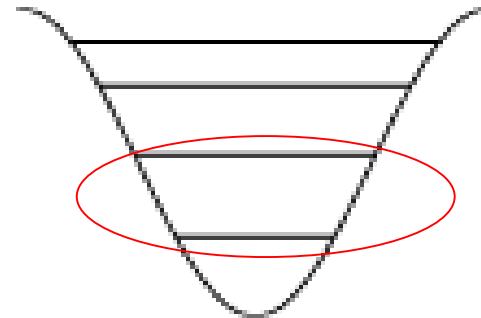
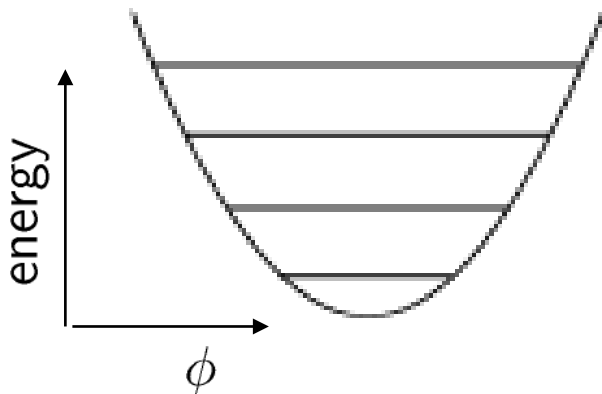
LC resonator:



Josephson junction resonator:
Josephson junction = nonlinear inductor

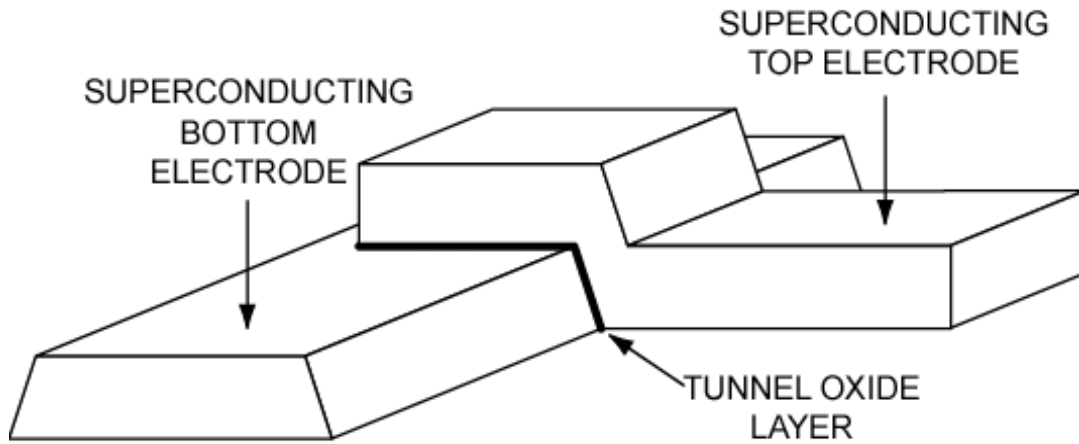


anharmonicity defines effective two-level system



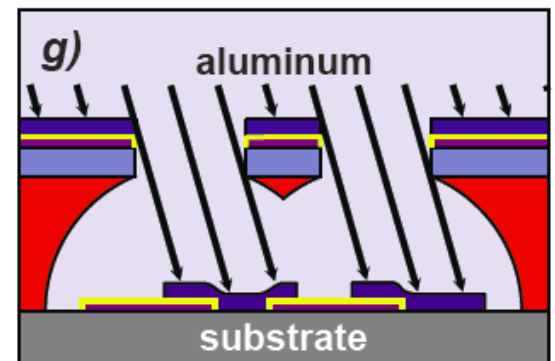
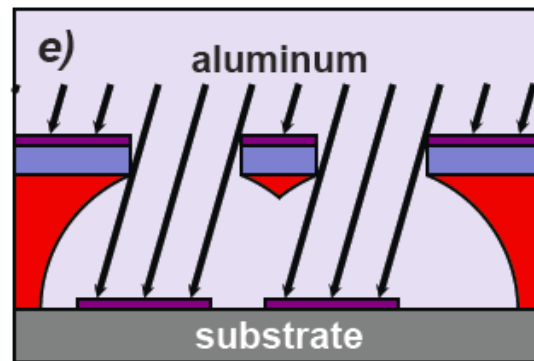
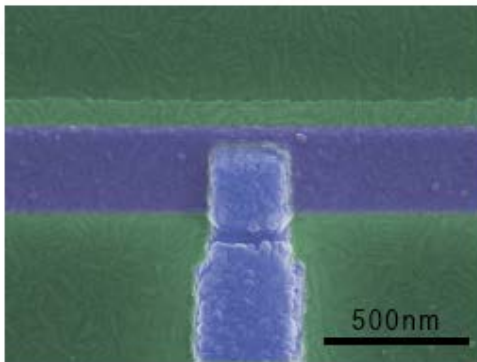
A Low-Loss Nonlinear Element

a (superconducting) Josephson junction:



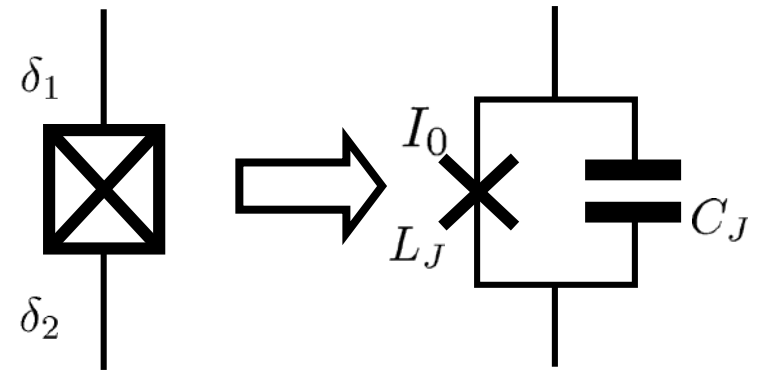
- superconductors: Nb, Al
- tunnel barrier: AlO_x

Josephson junction fabricated by shadow evaporation:



The Josephson Junction as an ideal Non-Linear Inductor

a nonlinear inductor without dissipation



Josephson relations:

$$I = I_0 \sin \delta = I_0 \sin [2\pi\phi(t)/\phi_0]$$

nonlinear
current/phase
relation

$$V = \frac{\phi_0}{2\pi} \dot{\delta} = \dot{\phi}$$

gauge inv. phase difference:

$$\delta = \delta_2 - \delta_1 = 2\pi\phi(t)/\phi_0$$

Josephson inductance:

$$V = -L_J \dot{I} = \frac{\phi_0}{2\pi I_0} \frac{1}{\cos \delta} \dot{I}$$

specific Josephson
inductance L_{J0}

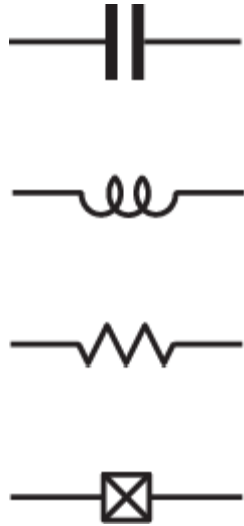
Josephson energy:

$$E_J = \int V I dt = \frac{I_0 \phi_0}{2\pi} \cos \delta$$

specific Josephson
energy E_{J0}

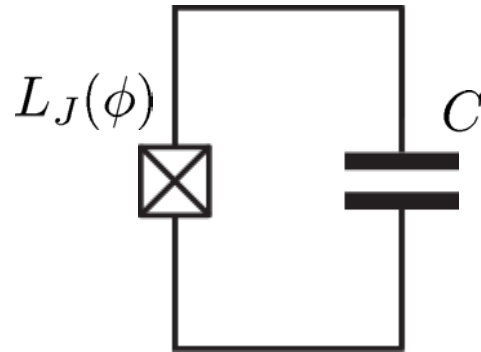
Constructing Non-Linear Quantum Electronic Circuits

circuit elements:



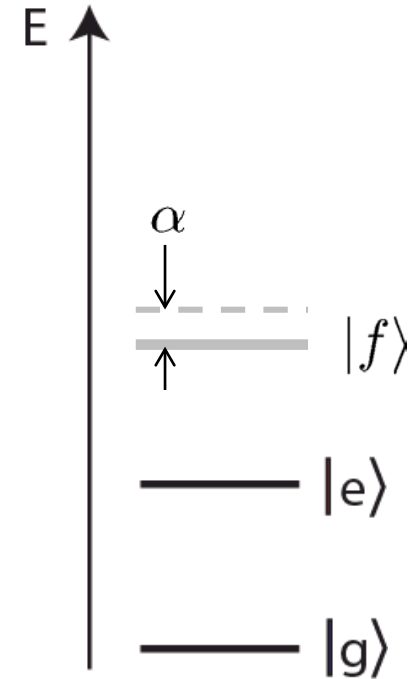
Josephson junction:
a non-dissipative nonlinear
element (inductor)

anharmonic oscillator:



$$H \approx \hbar(\omega_{ge} \hat{b}^\dagger \hat{b} - \frac{\alpha}{2} \hat{b}^{\dagger 2} \hat{b}^2)$$

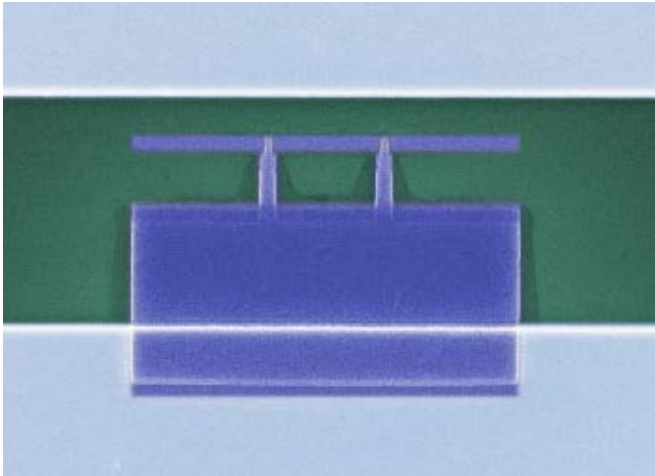
non-linear energy
level spectrum:



electronic
artificial atom

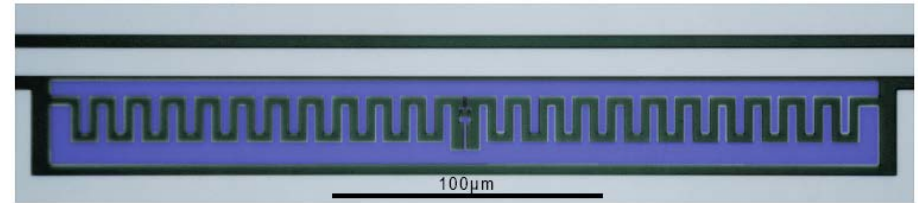
Flavors of Superconducting Artificial Atoms

Cooper pair box:



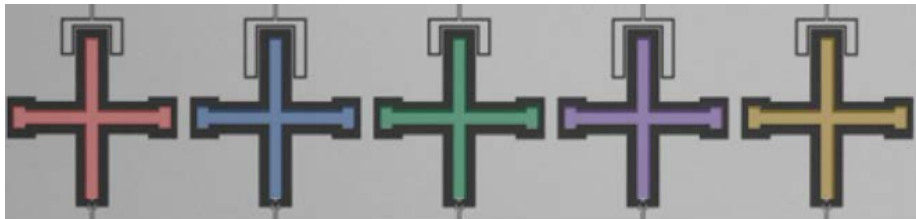
Bouchiat et al., *Physica Scripta* T76, 165 (1998).

Transmon:



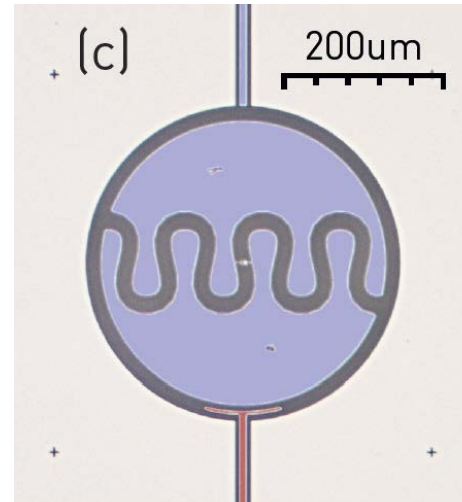
J. Koch et al., *PRA* 76, 042319 (2007)

Xmons:



Barends et al., *Phys. Rev. Lett.* 111, 080502 (2013)

(Jellymon):

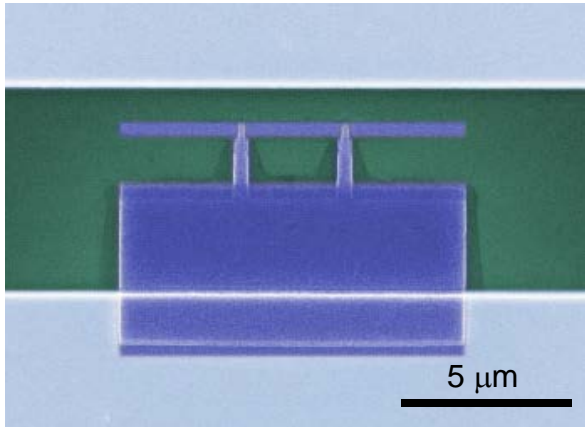


Pechal et al., *arXiv:1606.01031* (2004)

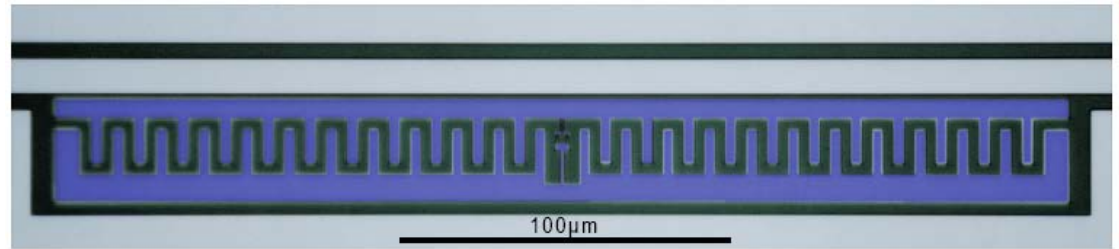
A Variant of the Cooper Pair Box

a Cooper pair box with a small charging energy

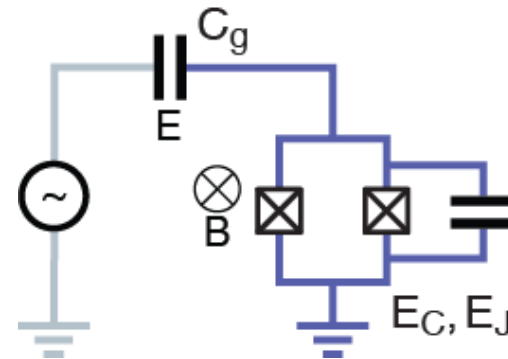
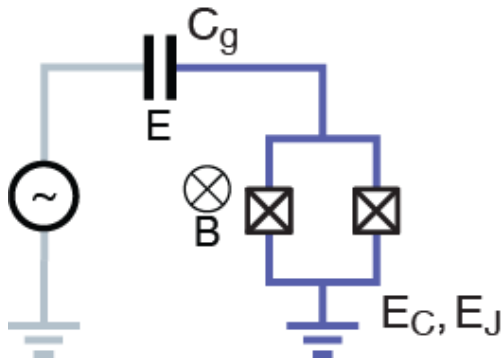
standard CPB:



Transmon qubit:



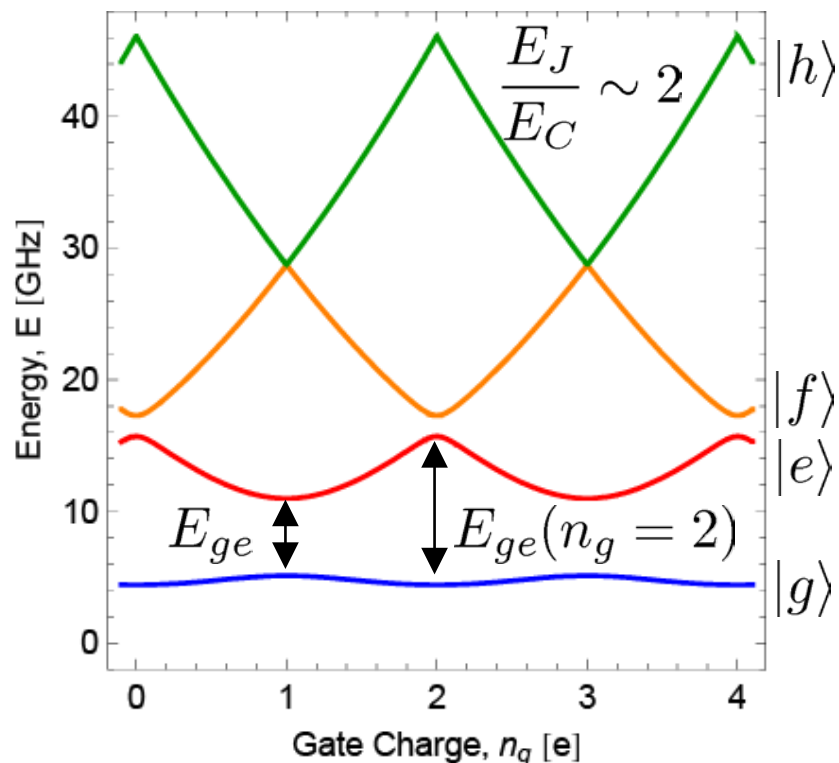
circuit diagram:



J. Koch *et al.*, Phys. Rev. A 76, 042319 (2007)
J. Schreier *et al.*, Phys. Rev. B 77, 180502 (2008)

The Transmon: A Charge Noise Insensitive Qubit

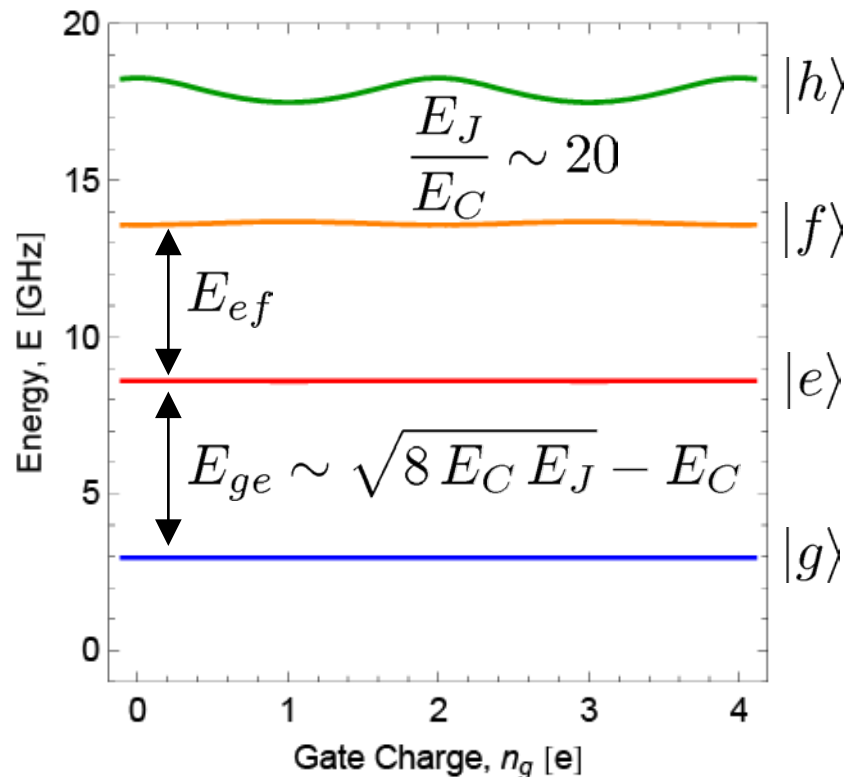
Cooper pair box energy levels:



dispersion:

$$\epsilon = E_{ge}(n_g = 1) - E_{ge}(n_g = 2)$$

Transmon energy levels:

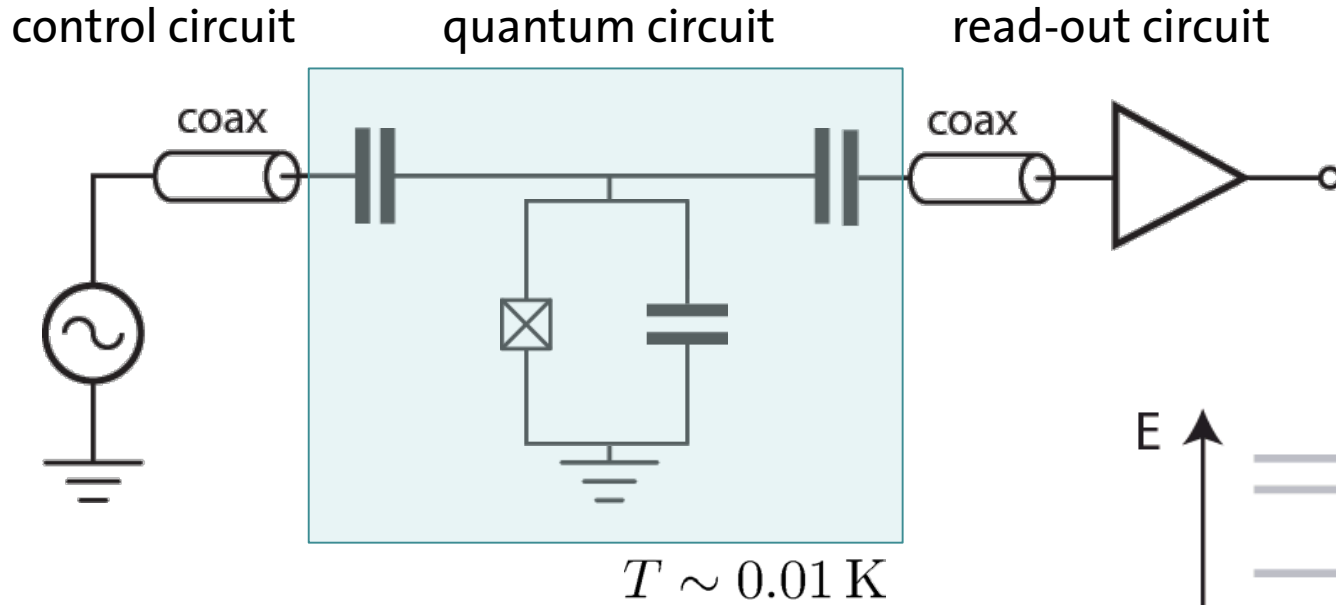


relative anharmonicity:

$$\alpha_r = \frac{E_{ef} - E_{ge}}{E_{ge}}$$

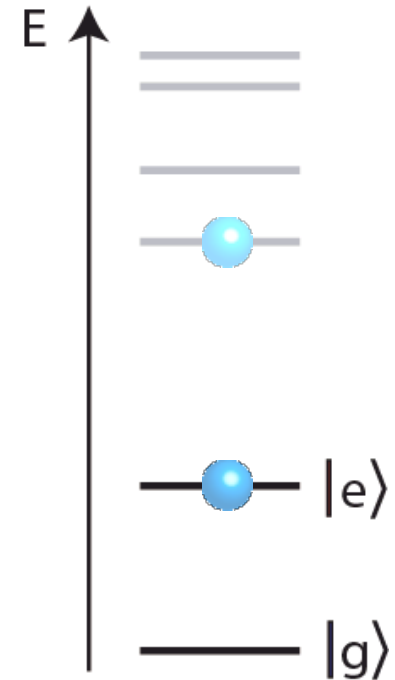
J. Koch *et al.*, Phys. Rev. A 76, 042319 (2007)

How to Operate Circuits in the Quantum Regime?

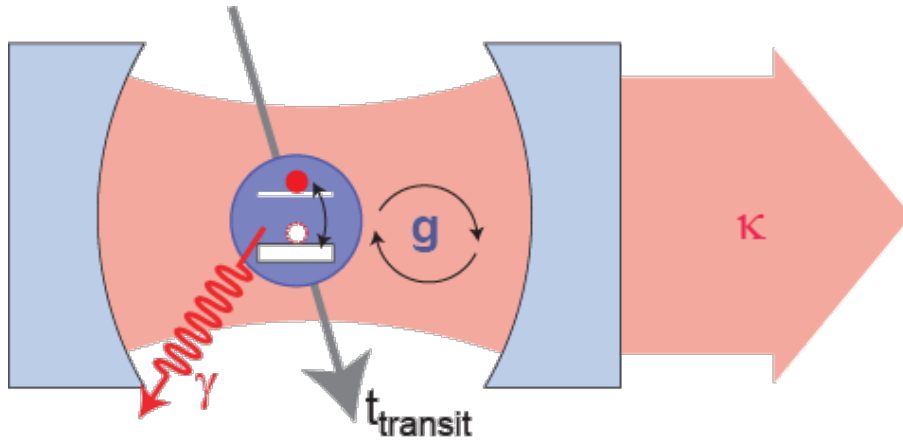


recipe:

- avoid dissipation
- work at low temperatures
- isolate quantum circuit from environment

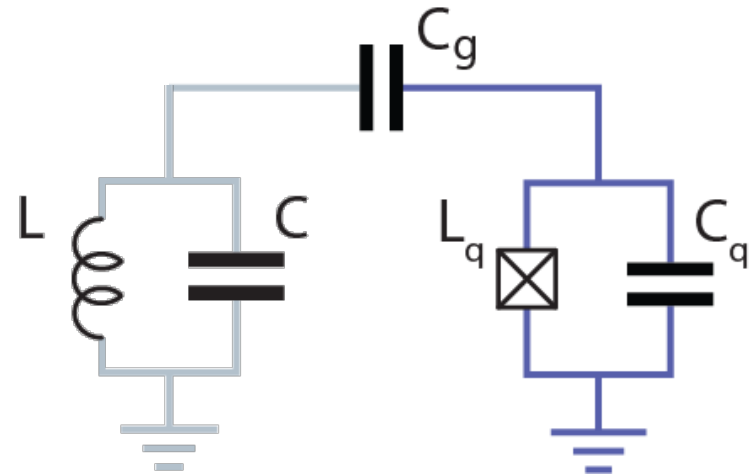


Cavity QED with Superconducting Circuits



coherent quantum mechanics
with individual photons and qubits ...

... basic approach:

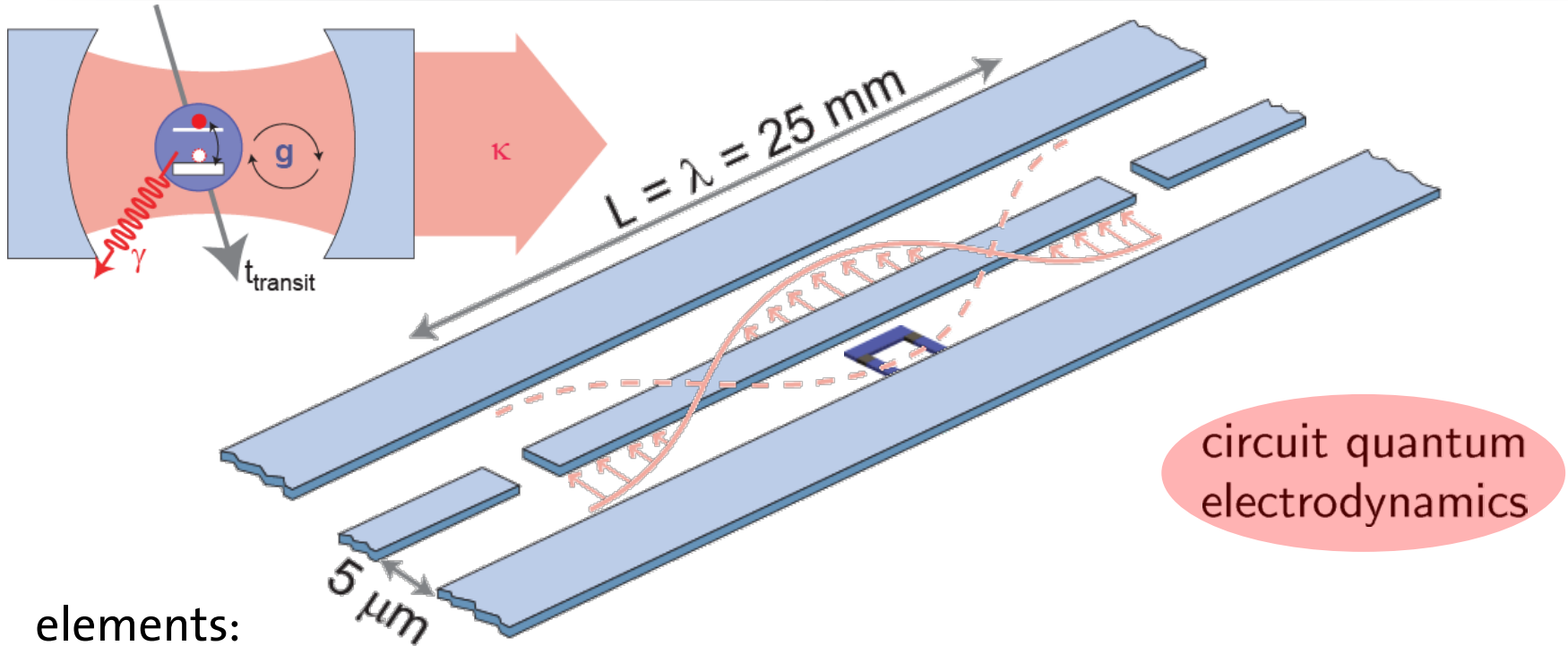


What is this good for?

- Isolating qubits from their environment
- Maintain addressability of qubits
- Reading out the state of qubits
- Coupling qubits to each other
- Converting stationary qubits to flying qubits



Cavity QED with Superconducting Circuits



elements:

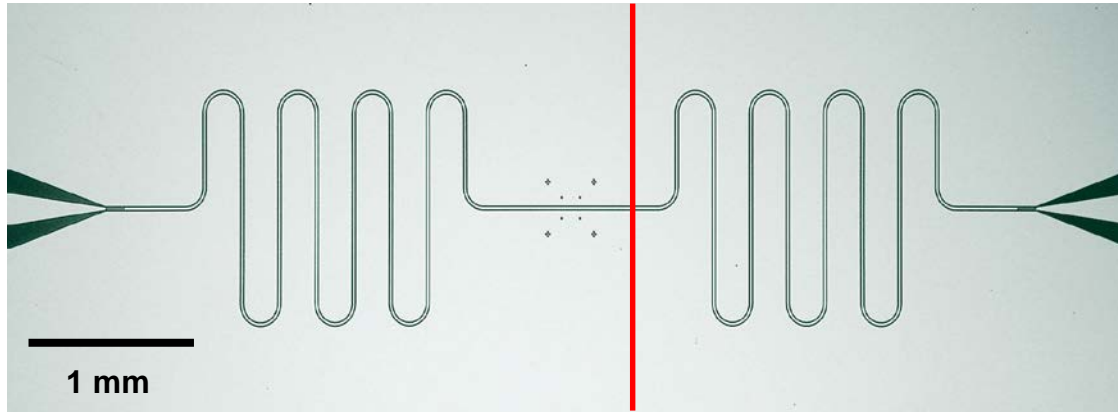
- the cavity: a superconducting 1D transmission line resonator with **large vacuum field** E_0 and **long photon life time** $1/\kappa$
- the atom: a superconducting qubit with **large dipole moment** d and **long coherence time** $1/\gamma$ and **fixed position** ...
- ... or any microscopic/macroscopic quantum element or ensemble thereof with an appreciable dipole moment

A. Blais, *et al.*, *PRA* 69, 062320 (2004)

A. Wallraff *et al.*, *Nature (London)* 431, 162 (2004)

R. J. Schoelkopf, S. M. Girvin, *Nature (London)* 451, 664 (2008)

Large Vacuum Field in 1D Cavity



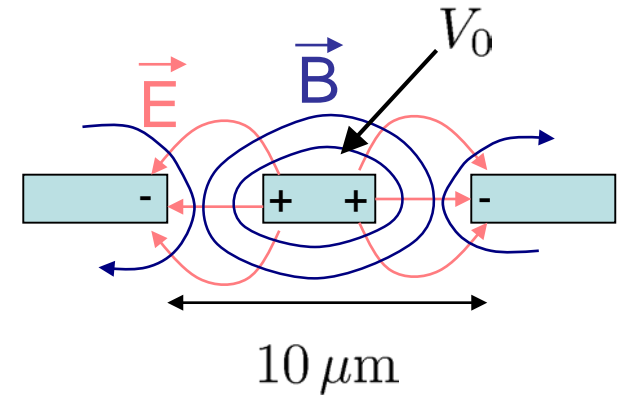
optical microscope image of strip line resonator

electric field across resonator in vacuum state ($n=0$):

$$\int \epsilon_0 E_{0,\text{rms}}^2 dV_{\text{mod}} = \frac{\hbar\omega_r}{2}$$

$$E_{0,\text{rms}} \approx 0.2 \text{ V/m}$$

for $\omega_r/2\pi \approx 6 \text{ GHz}$

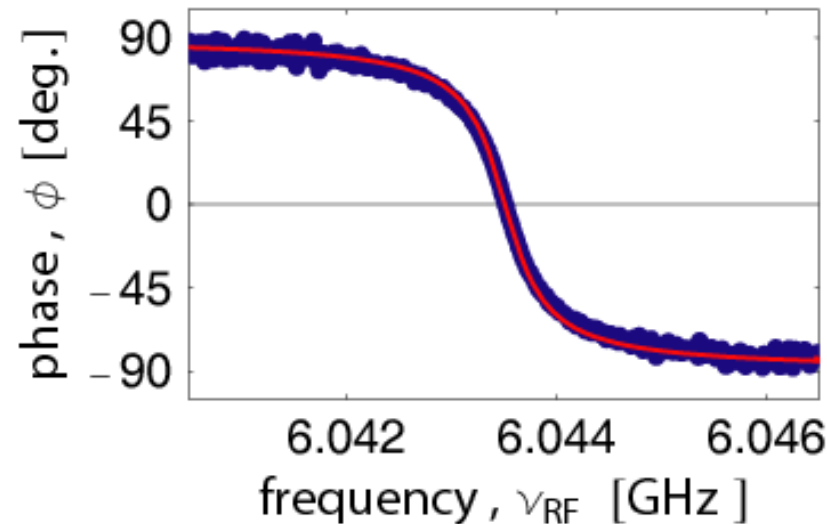
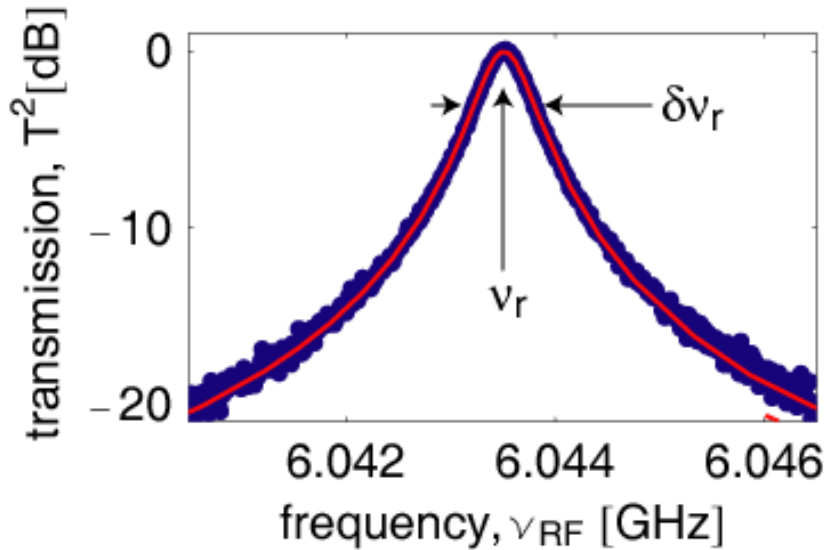


harmonic oscillator

$$H_r = \hbar\omega_r \left(a^\dagger a + \frac{1}{2} \right)$$

$\times 10^6$ larger than E_0
in 3D microwave cavity

Resonator Quality Factor and Photon Lifetime



resonance frequency:

$$\nu_r = 6.04 \text{ GHz}$$

quality factor:

$$Q = \frac{\nu_r}{\delta\nu_r} \approx 10^4$$

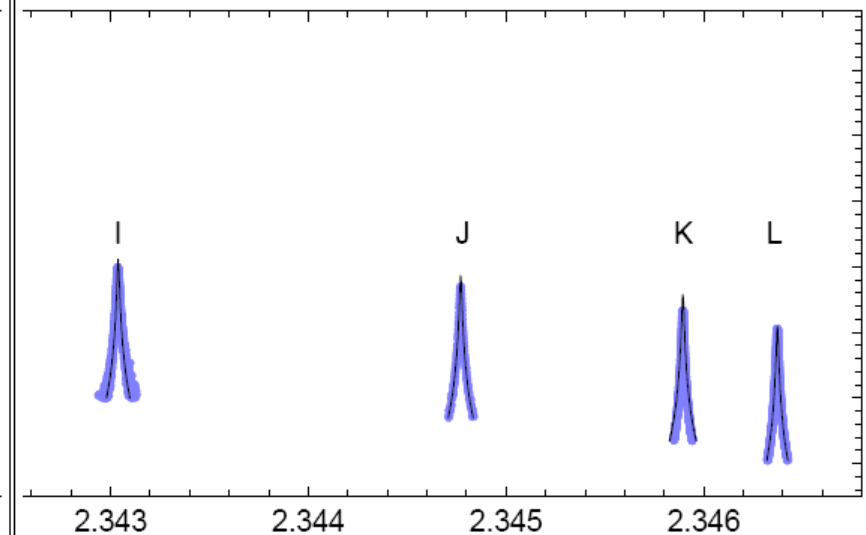
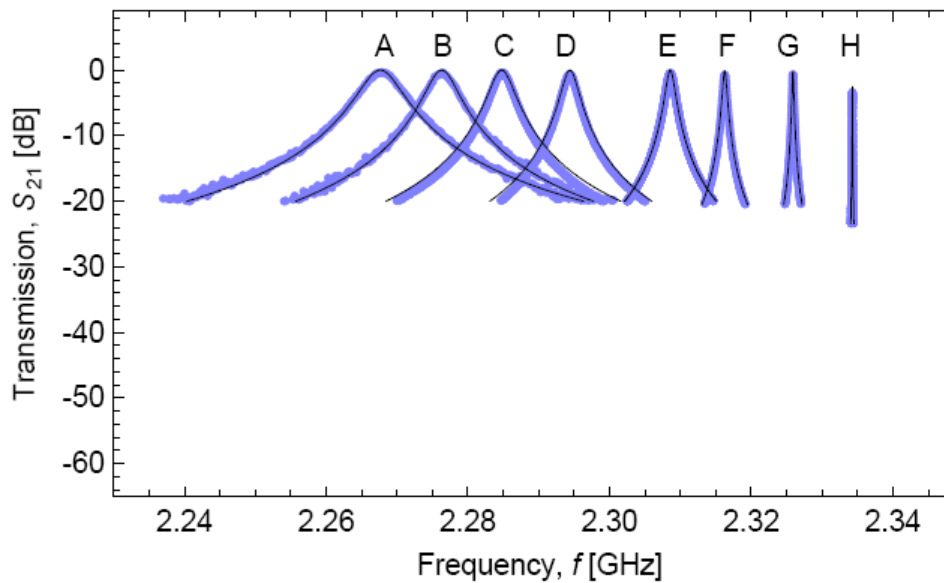
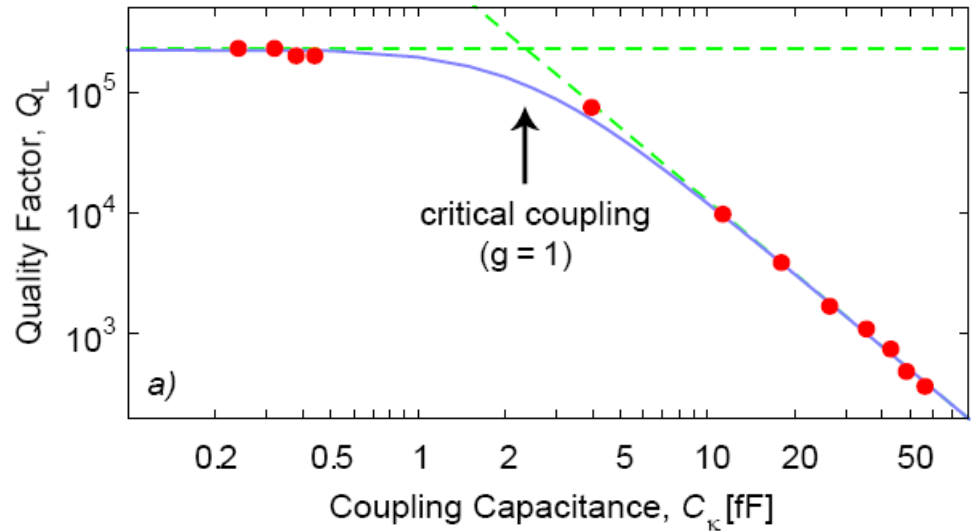
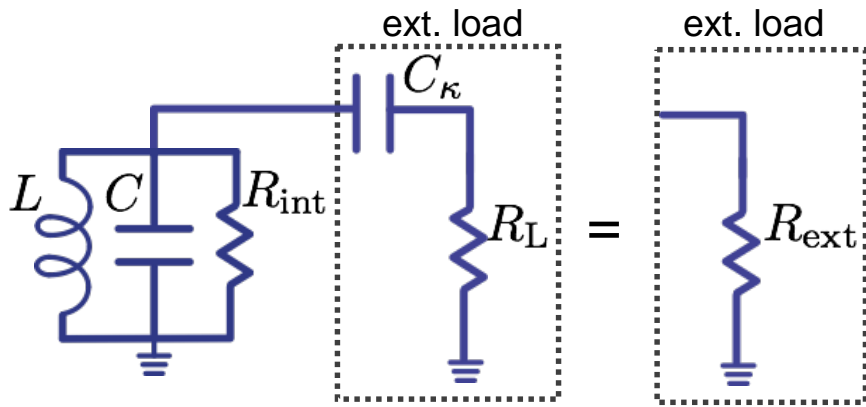
photon decay rate:

$$\frac{\kappa}{2\pi} = \frac{\nu_r}{Q} \approx 0.8 \text{ MHz}$$

photon lifetime:

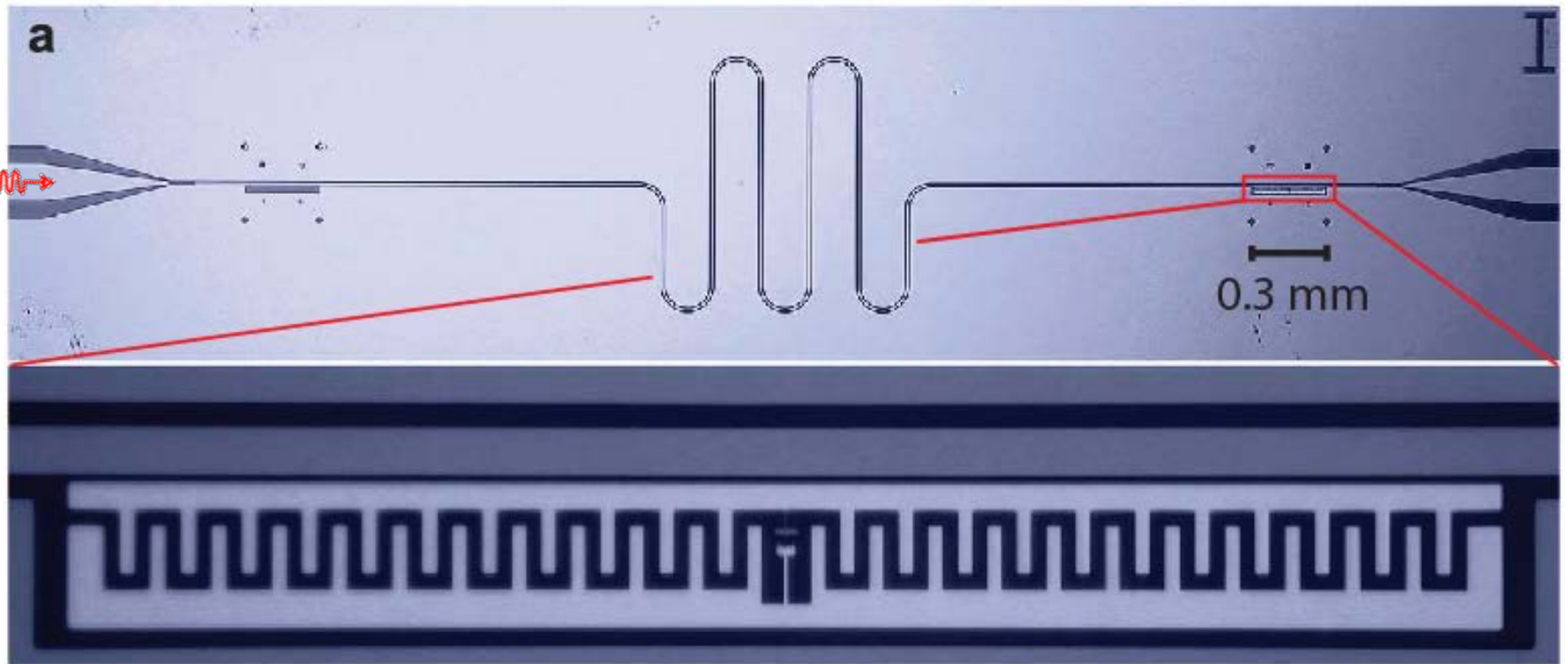
$$T_\kappa = 1/\kappa \approx 200 \text{ ns}$$

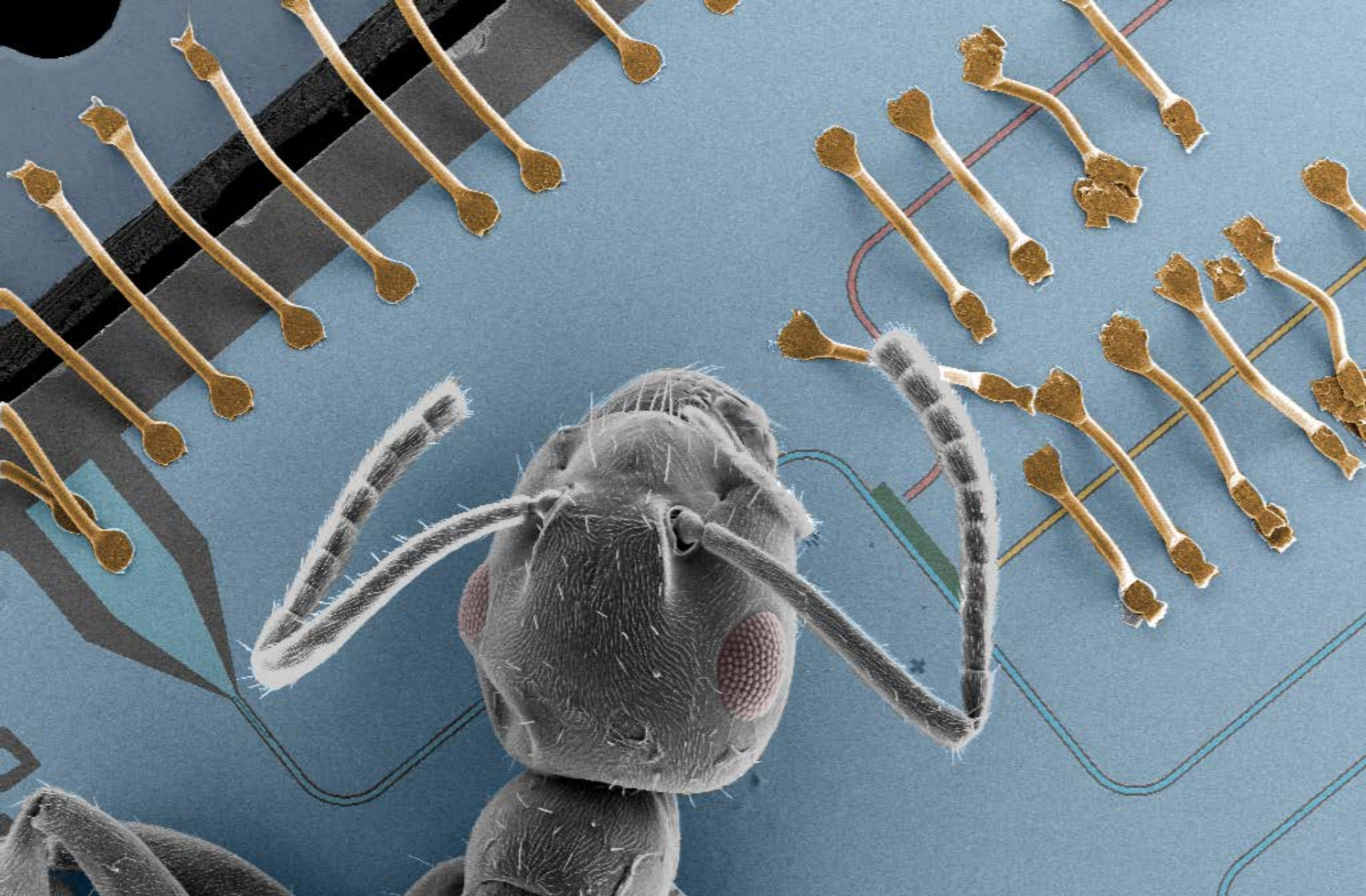
The Quality Factor



M. Goeppel, et al. *J. Appl. Phys.* 104, 113904 (2008)

Realization

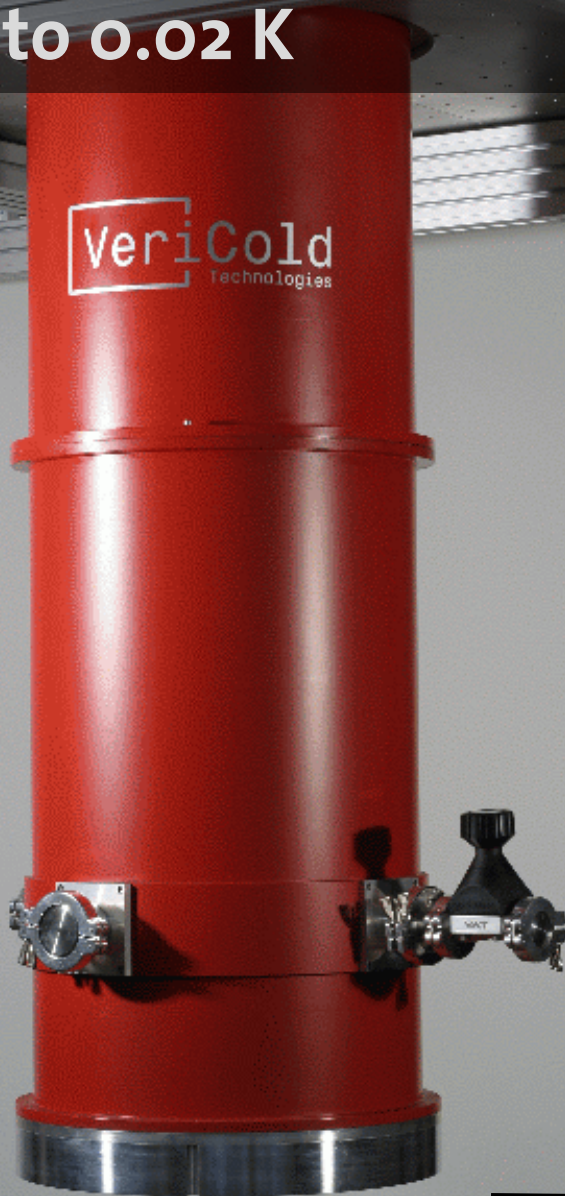




Sample Mount

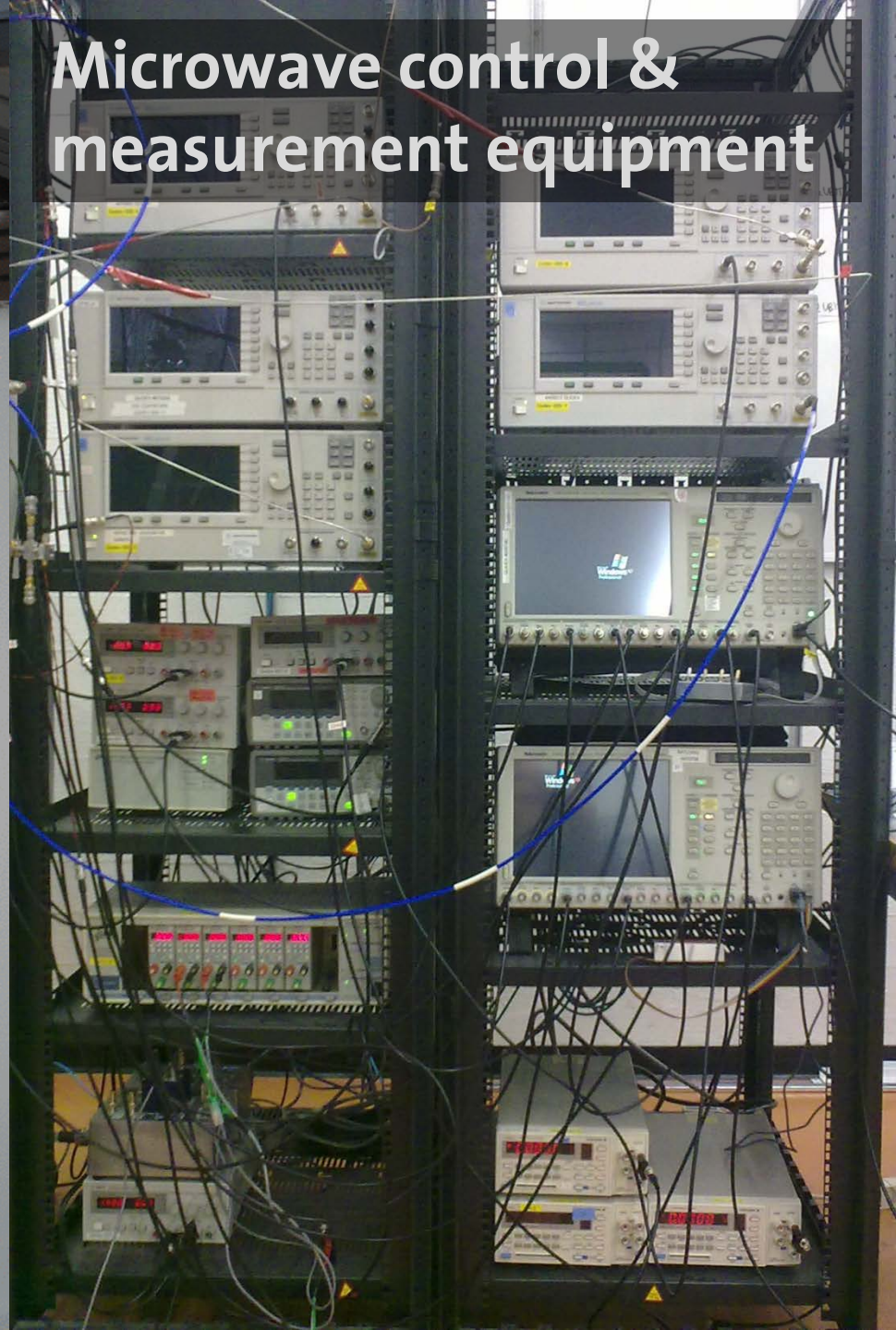


Cryostat for temperatures down to 0.02 K

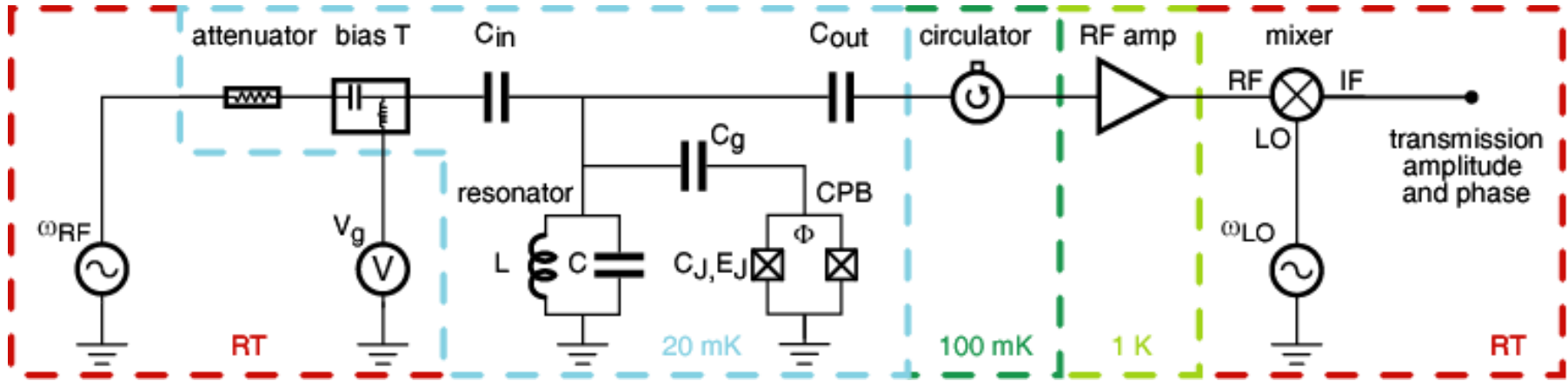


~ 20 cm

Microwave control & measurement equipment



How to do the Measurement



- prevent leakage of thermal photons (cold attenuators and circulators)
- average power to be detected ($\omega_r/2\pi = 6$ GHz, $\kappa/2\pi = 1$ MHz)

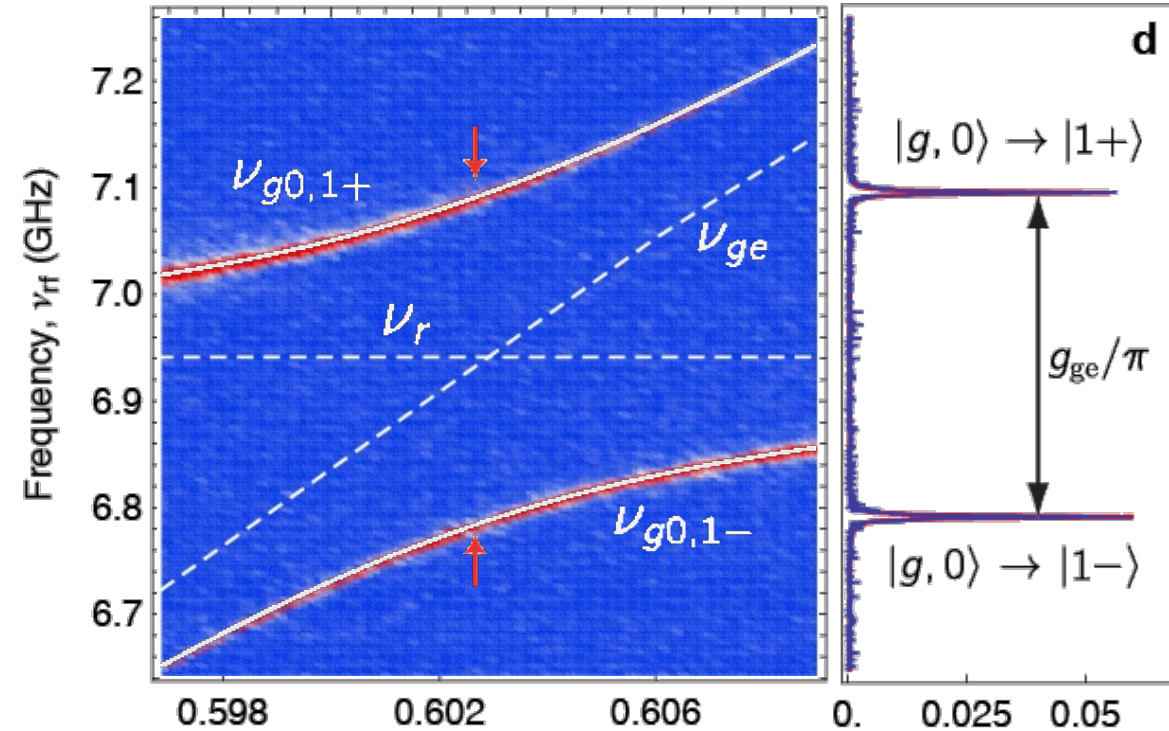
$$\langle n = 1 \rangle \hbar \omega_r \kappa / 2 \approx P_{RF} = -140 \text{ dBm} = 10^{-17} \text{ W}$$

- efficient with cryogenic low noise HEMT amplifier $T_N = 6$ K

Resonant Vacuum Rabi Mode Splitting ...

... with one photon ($n=1$):

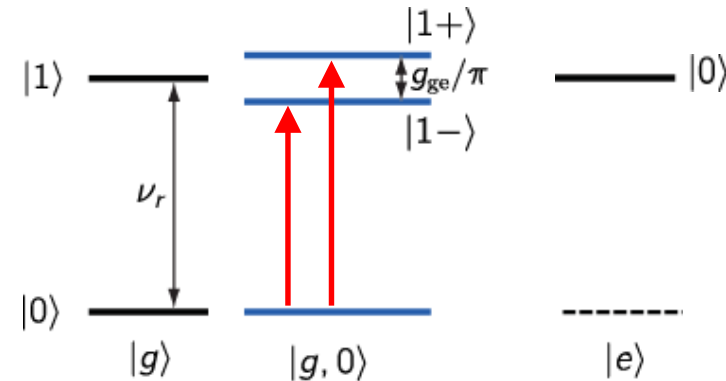
very strong coupling:



$$g_{ge}/\pi = 308 \text{ MHz}$$

$$\kappa, \gamma < 1 \text{ MHz}$$

$$g_{ge} \gg \kappa, \gamma$$



forming a 'molecule' of a qubit and a photon

first demonstration in a solid: A. Wallraff *et al.*, *Nature (London)* 431, 162 (2004)

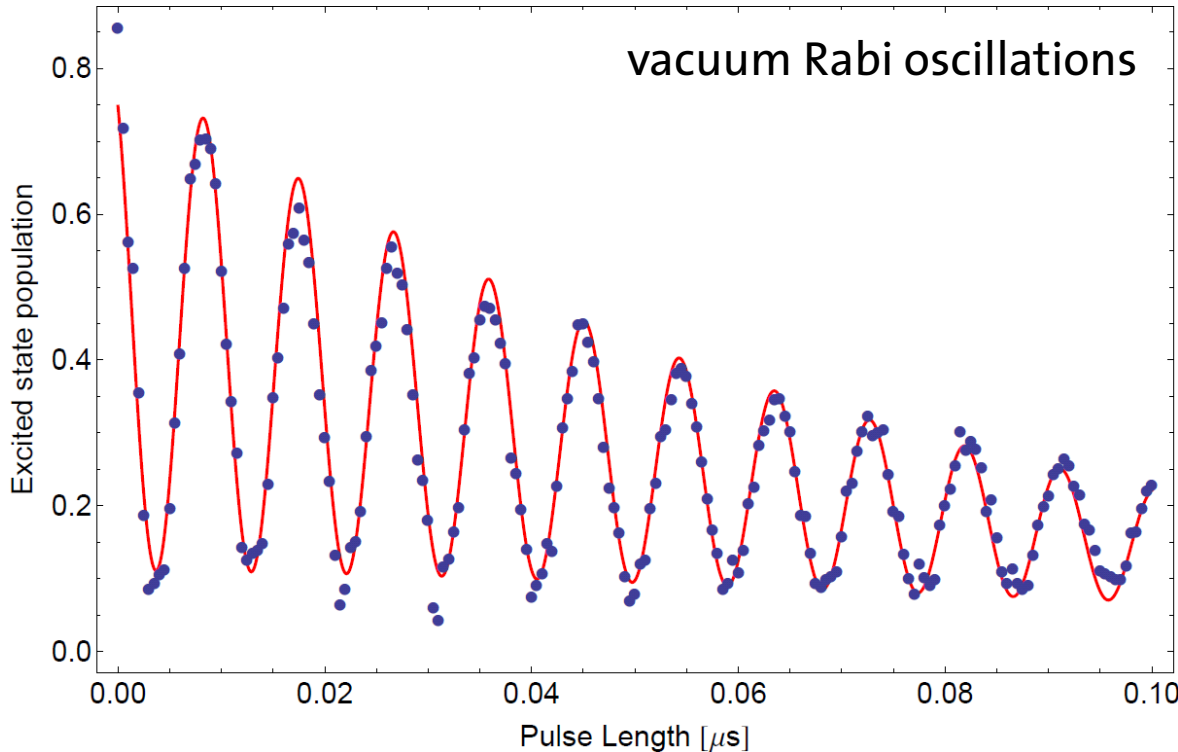
this data: J. Fink *et al.*, *Nature (London)* 454, 315 (2008)

R. J. Schoelkopf, S. M. Girvin, *Nature (London)* 451, 664 (2008)

Resonant Vacuum Rabi Mode Splitting ...

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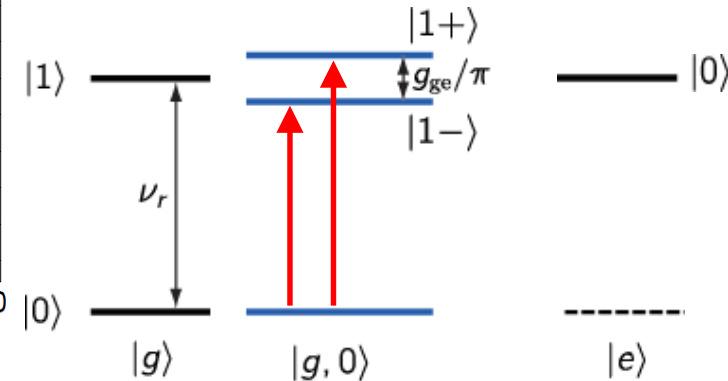
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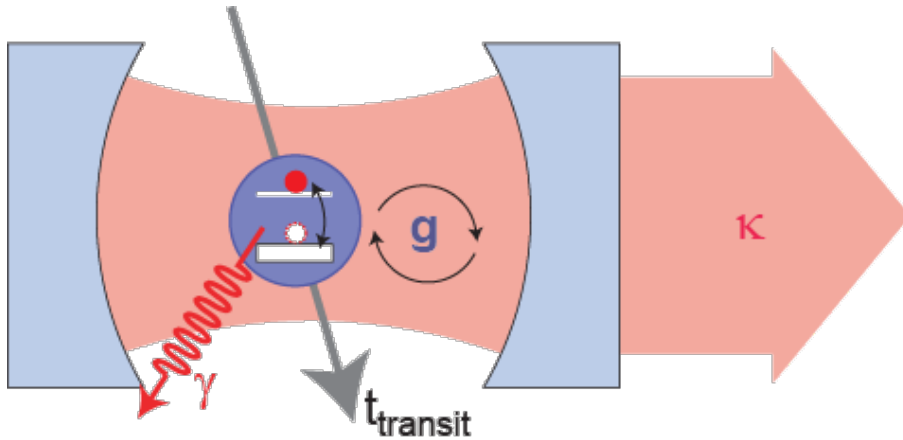
forming a 'molecule' of a qubit and a photon

first demonstration in a solid: A. Wallraff *et al.*, *Nature (London)* 431, 162 (2004)

this data: J. Fink *et al.*, *Nature (London)* 454, 315 (2008)

R. J. Schoelkopf, S. M. Girvin, *Nature (London)* 451, 664 (2008)

Cavity QED with Superconducting Circuits



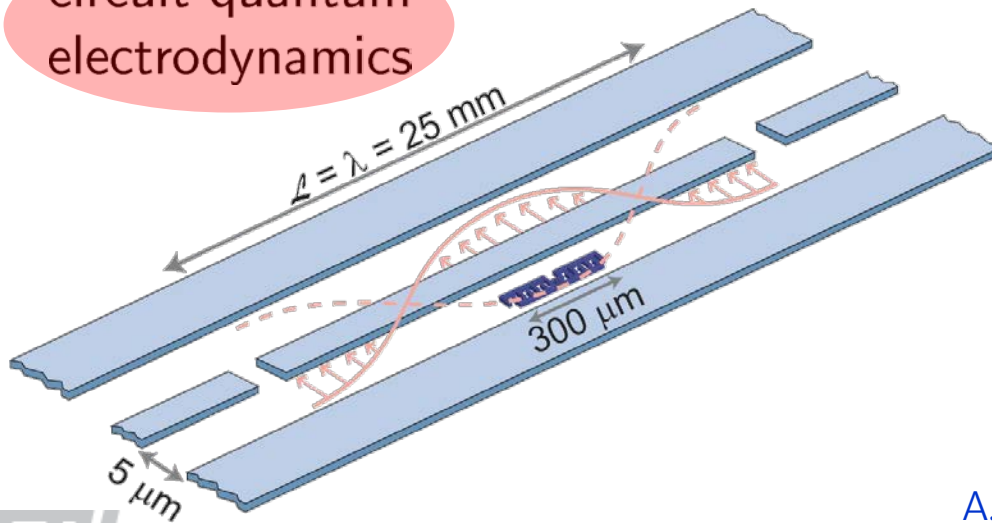
coherent interaction of photons with quantum two-level systems ...

J. M. Raimond *et al.*, *Rev. Mod. Phys.* 73, 565 (2001)

S. Haroche & J. Raimond, *oup Oxford* (2006)

J. Ye., H. J. Kimble, H. Katori, *Science* 320, 1734 (2008)

circuit quantum electrodynamics



Properties:

- strong coupling in solid state sys.
- 'easy' to fabricate and integrate

Research directions:

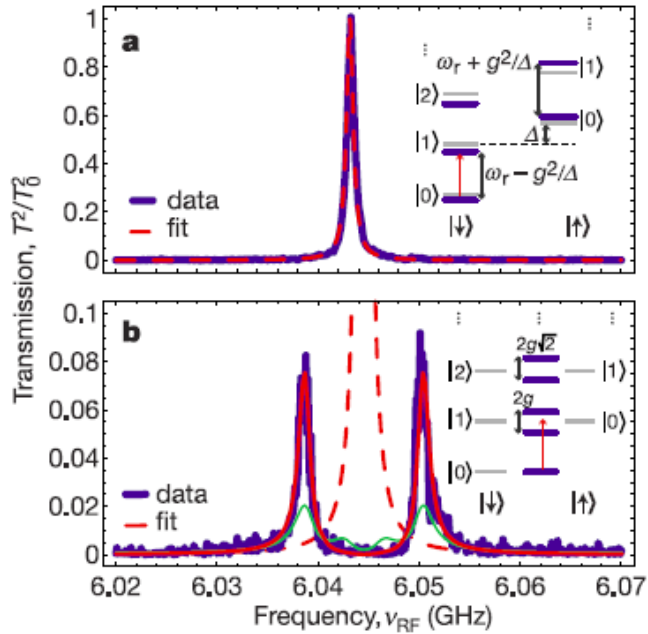
- quantum optics
- quantum information
- hybrid quantum systems

A. Blais, *et al.*, *PRA* 69, 062320 (2004)

A. Wallraff *et al.*, *Nature (London)* 431, 162 (2004)

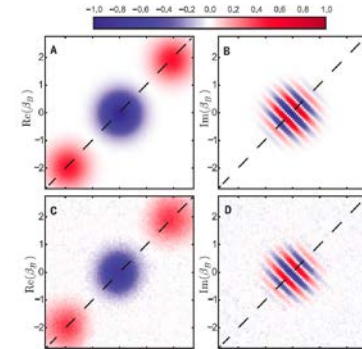
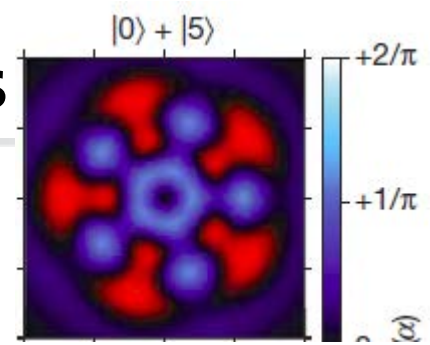
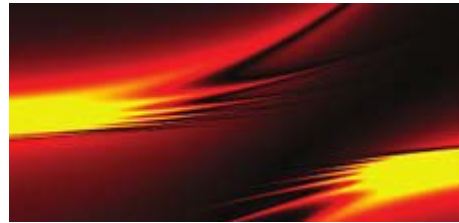
R. J. Schoelkopf, S. M. Girvin, *Nature (London)* 451, 664 (2008)

Quantum Optics with Supercond. Circuits

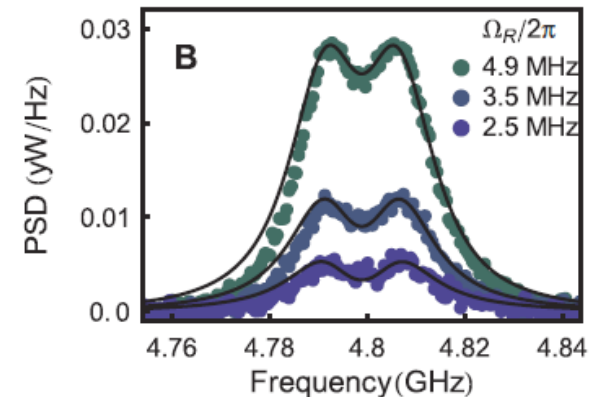


Strong Coherent Coupling
 Chiorescu *et al.*, *Nature* 431, 159 (2004)
 Wallraff *et al.*, *Nature* 431, 162 (2004)
 Schuster *et al.*, *Nature* 445, 515 (2007)

Root n Nonlinearities
 Fink *et al.*, *Nature* 454, 315 (2008)
 Deppe *et al.*, *Nat. Phys.* 4, 686 (2008)
 Bishop *et al.*, *Nat. Phys.* 5, 105 (2009)

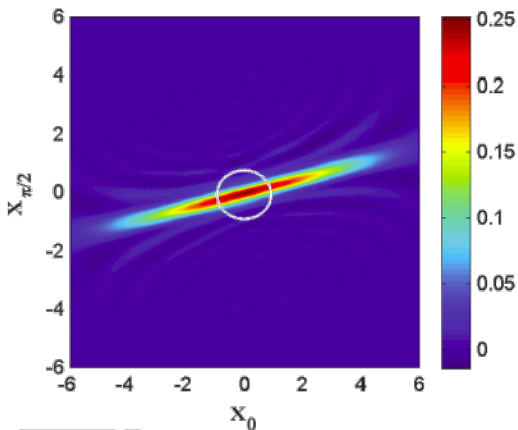


Microwave Fock and Cat States
 Hofheinz *et al.*, *Nature* 454, 310 (2008)
 Hofheinz *et al.*, *Nature* 459, 546 (2009)
 Kirchmair *et al.*, *Nature* 495, 205 (2013)
 Vlastakis *et al.*, *Science* 342, 607 (2013)
 Wang *et al.*, *Science* 352, 1087 (2016)



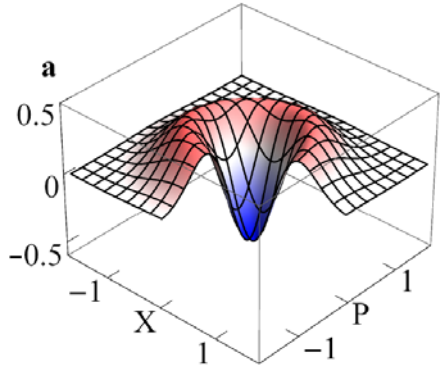
Parametric Amplification & Squeezing
 Castellanos-Beltran *et al.*, *Nat. Phys.* 4, 928 (2008)
 Abdo *et al.*, *PRX* 3, 031001 (2013)

Waveguide QED – Qubit Interactions in Free Space
 Astafiev *et al.*, *Science* 327, 840 (2010)
 van Loo *et al.*, *Science* 342, 1494 (2013)



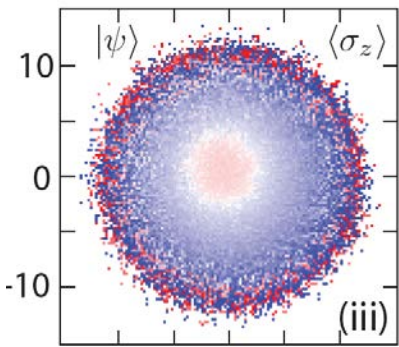
Experiments with Propagating Microwaves

Full state tomography and Wigner functions of propagating photons



Eichler *et al.*, *PRL* 106, 220503 (2011)

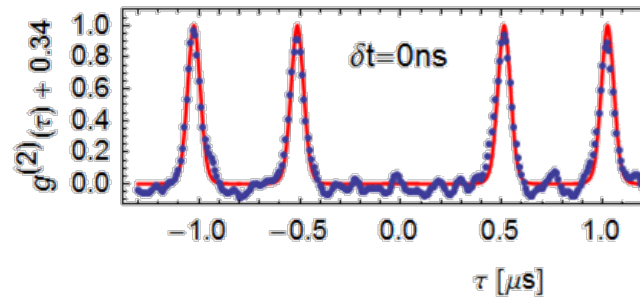
Preparation and characterization of qubit-propagating photon entanglement



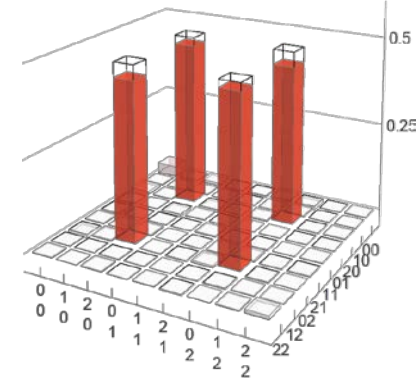
Eichler *et al.*, *PRL* 109, 240501 (2012)

Eichler *et al.*, *PRA* 86, 032106 (2012)

Hong-Ou-Mandel: Two-photon interference incl. msmrnt of coherences at microwave freq.

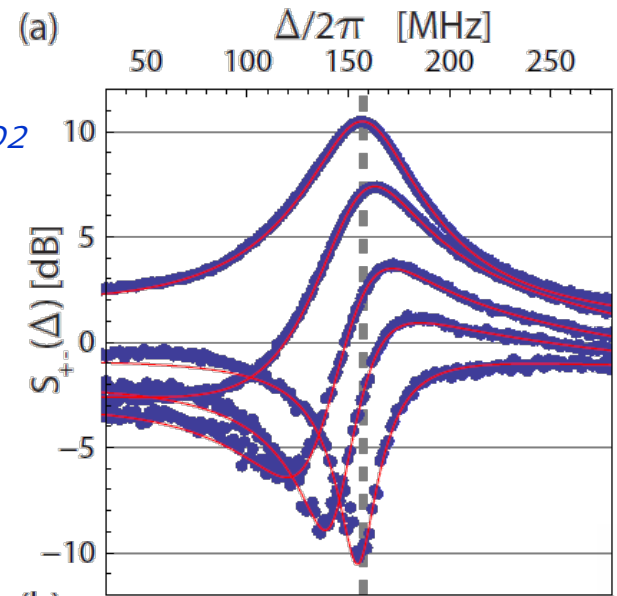


Lang *et al.*, *Nat. Phys.* 9, 345 (2013)



Squeezing in a Josephson parametric dimer

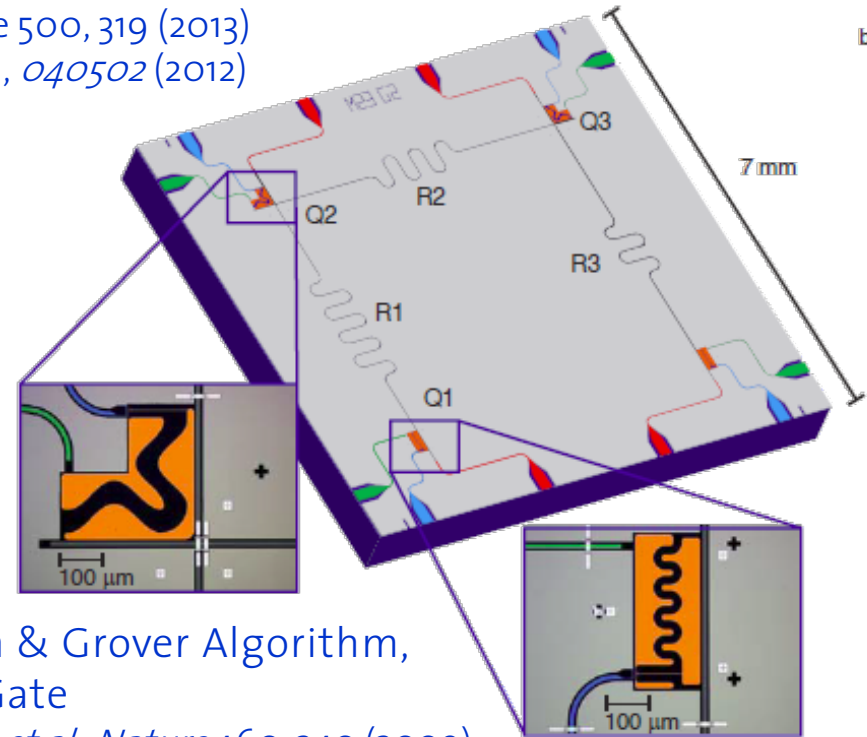
Eichler *et al.*,
PRL 113, 110502
(2014)



Quantum Computing with Superconducting Circuits

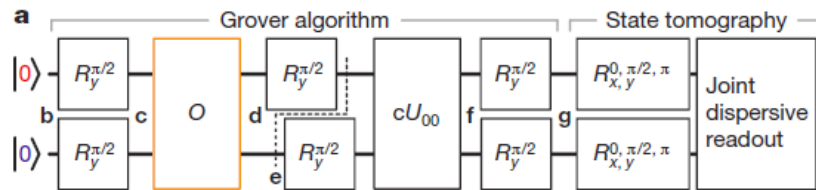
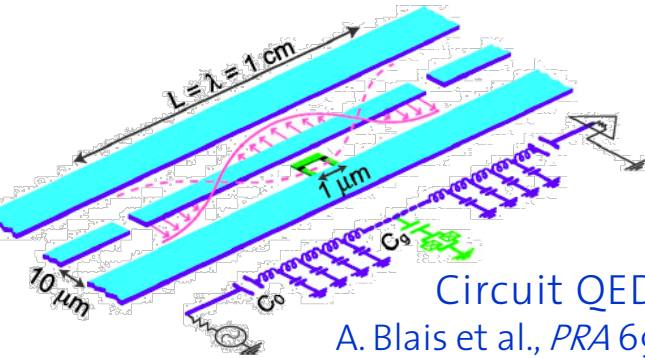
Teleportation

L. Steffen *et al.*, *Nature* 500, 319 (2013)
 M.. Baur *et al.*, *PRL* 108, 040502 (2012)



Circuit QED Architecture

A. Blais *et al.*, *PRA* 69, 062320 (2004)
 A. Wallraff *et al.*, *Nature* 431, 162 (2004)
 M. Sillanpaa *et al.*, *Nature* 449, 438 (2007)
 H. Majer *et al.*, *Nature* 449, 443 (2007)
 M. Mariani *et al.*, *Science* 334, 61 (2011)
 R. Barends *et al.*, *Nature* 508, 500 (2014)

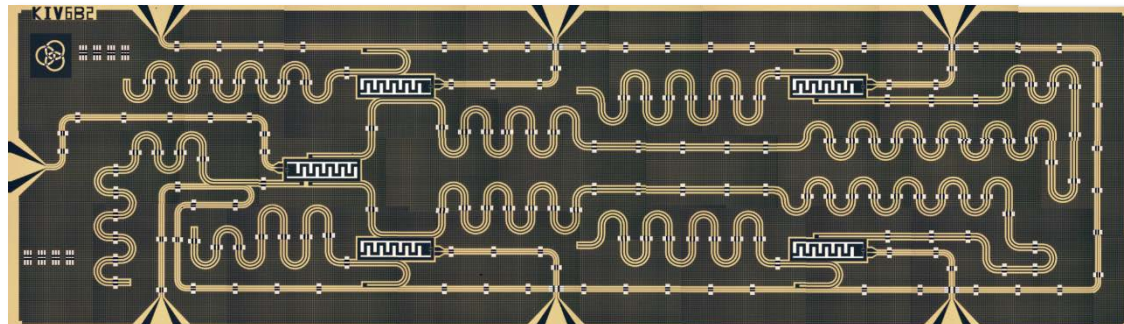


Deutsch & Grover Algorithm, Toffoli Gate

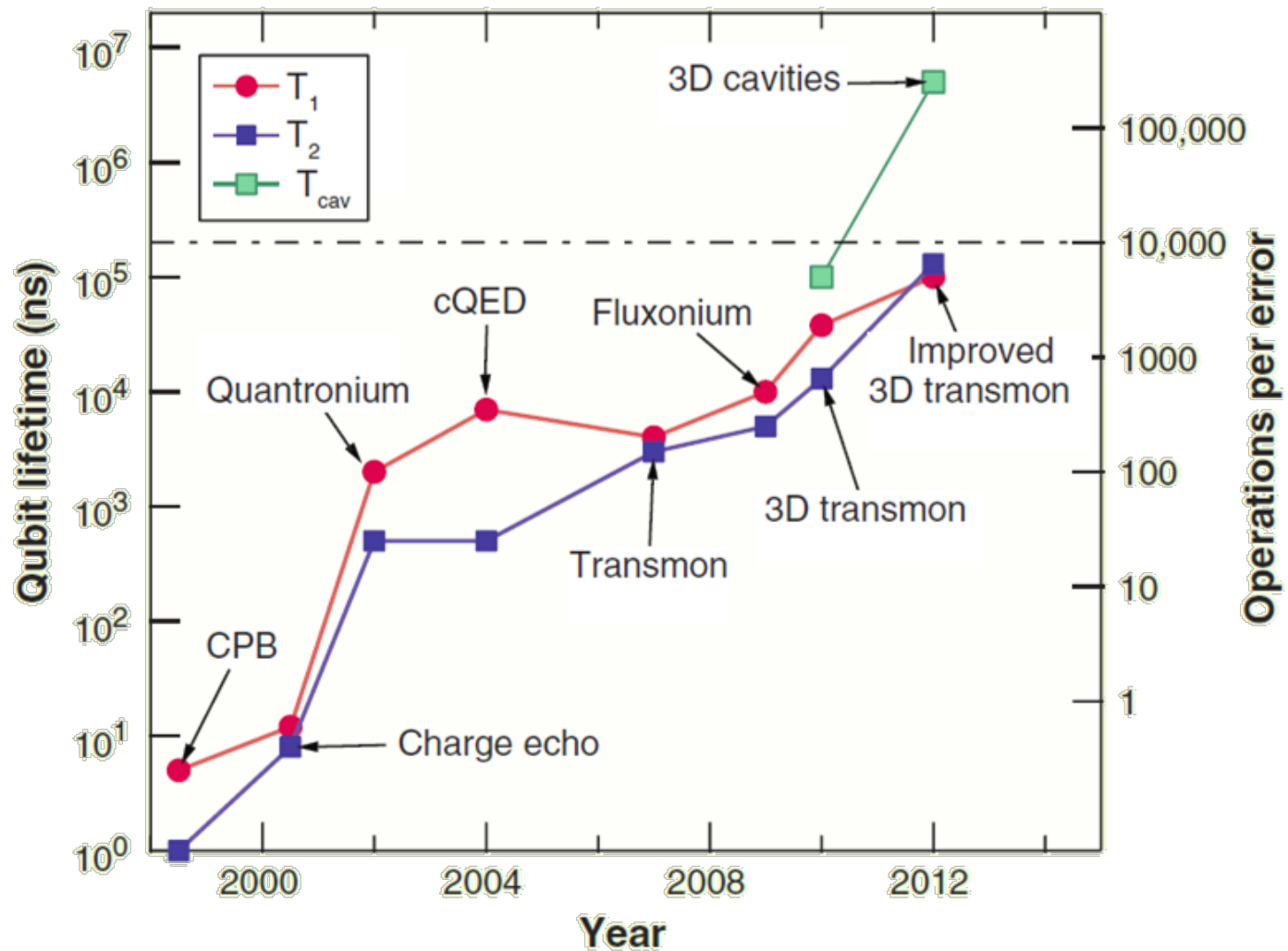
L. DiCarlo *et al.*, *Nature* 460, 240 (2009)
 L. DiCarlo *et al.*, *Nature* 467, 574 (2010)
 A. Fedorov *et al.*, *Nature* 481, 170 (2012)

Error Correction

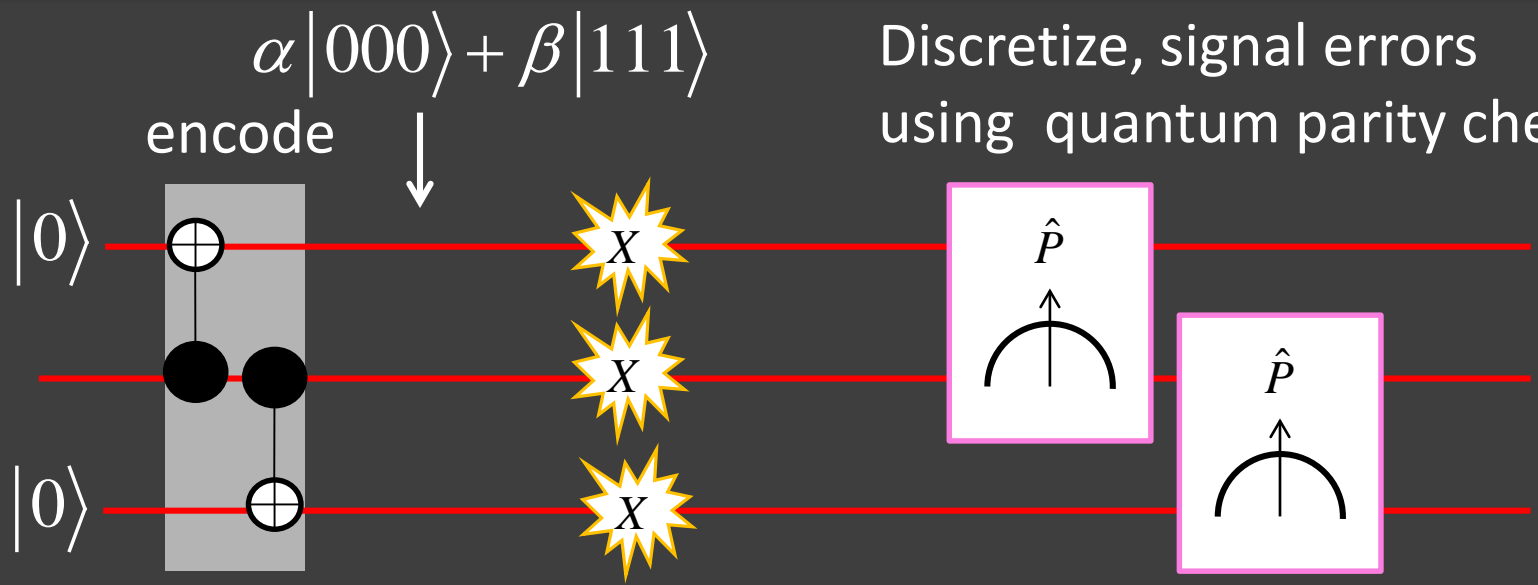
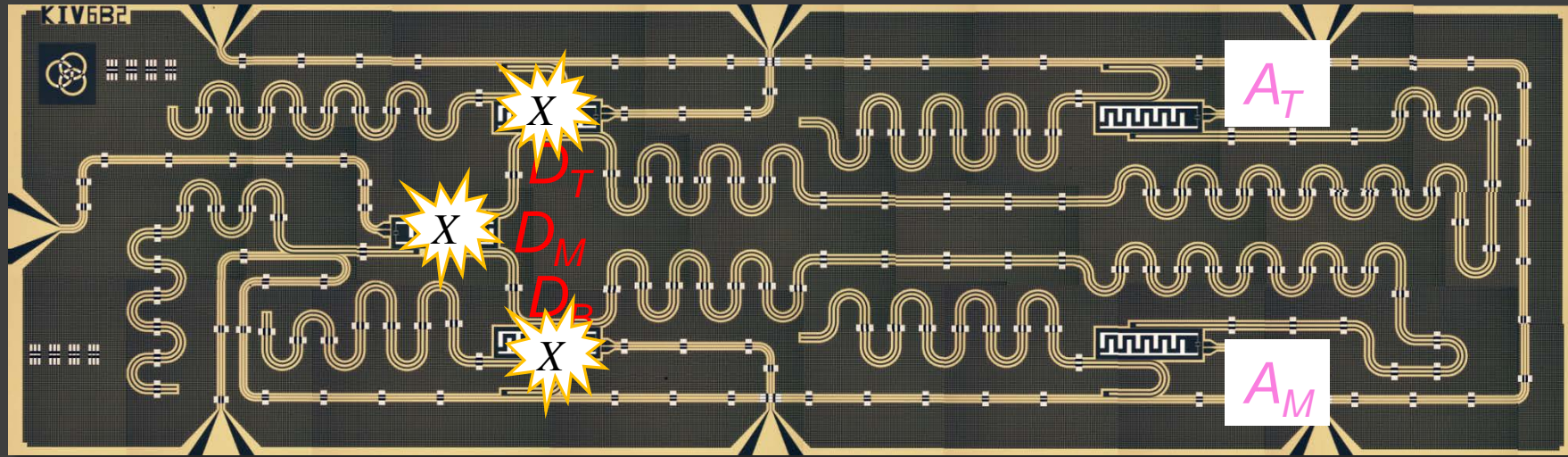
M. Reed *et al.*, *Nature* 481, 382 (2012)
 Corcoles *et al.*, *Nat. Com.* 6, 6979 (2015)
 Ristè *et al.*, *Nat. Com.* 6, 6983 (2015)
 Kelly *et al.*, *Nature* 519, 66-69 (2015)



10^5 Improvement in Coherence Time in 13 Years



Recent Progress in Quantum Error Correction



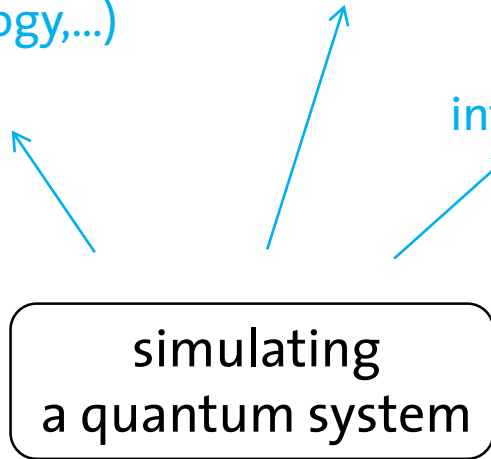
IBM: Corcoles *et al.*, *Nat. Com.* **6**, 6979 (2015), ArXiv:1410.6419
 QuTech: Ristè, Poletto, Huang *et al.*, *Nat. Com.* **6**, 6983 (2015), ArXiv:1411.5542
 UCSB/Google: Kelly *et al.*, *Nature* **519**, 66-69 (2015), ArXiv:1411.7403

Quantum Simulation

describes physical system of interest (physics, chemistry, biology,...)

encodes hard classical problem

interesting toy model



... difficult on classical computer!

use

universal quantum computer

... still to be realized!

... OR ...

map!

quantum simulator

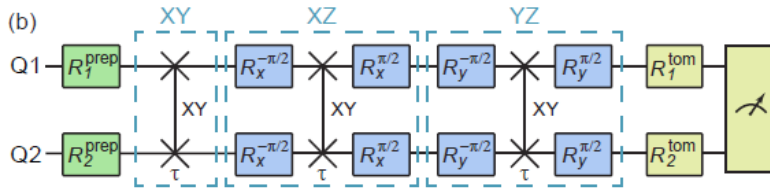
... sufficient controllability, flexibility!

Feynman, *Int. Journal of Th. Phys.* 21, 467 (1982)

Lloyd, *Science* 273, 5278 (1996)

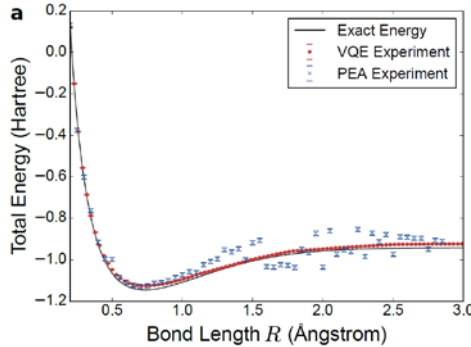
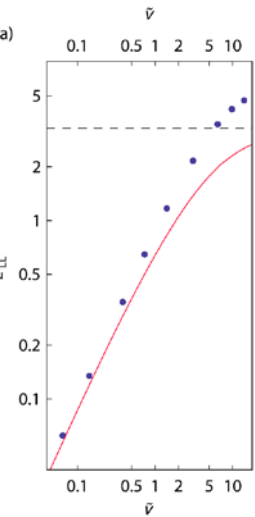
Quantum Simulation with Superconducting Circuits

Digital simulation of exchange, Heisenberg, Ising spin models



Salathe *et al.*, *PRX* 5, 021027 (2015)

Quantum simulation of correlated systems with variational Ansatz



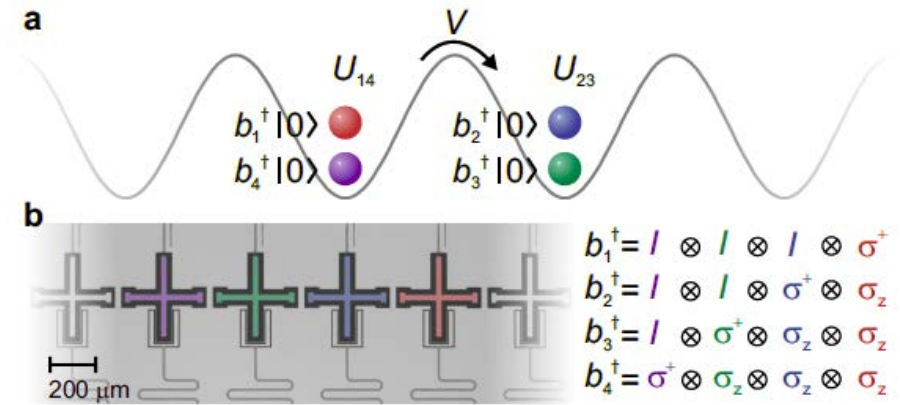
Eichler *et al.*, *Phys. Rev. X* 5, 041044 (2015)

O'Malley *et al.*, *arXiv:1512.06860* (2015)

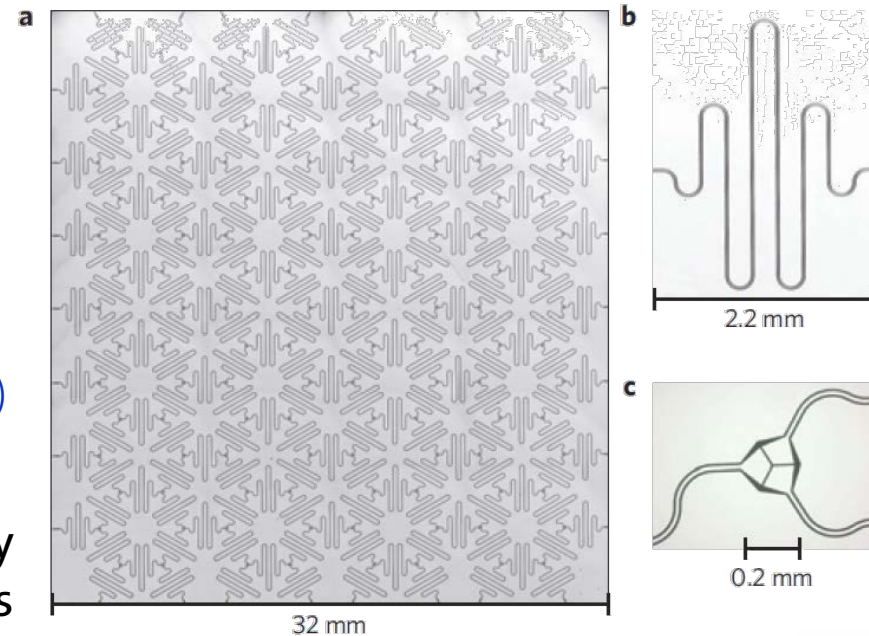
Analog simulations with cavity and/or qubit arrays

Houck *et al.*, *Nat Phys.* 8, 292 (2012)

... two-mode fermionic Hubbard models



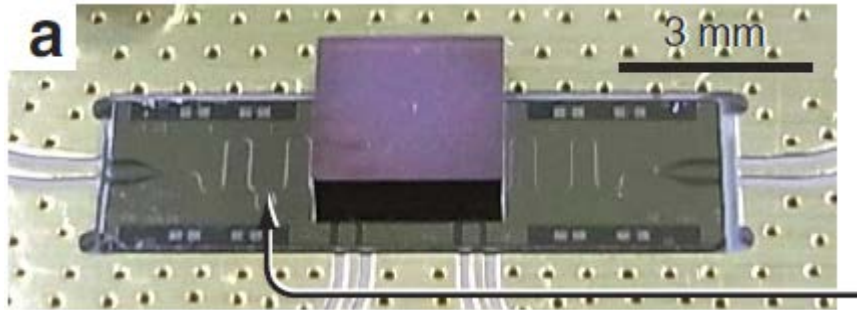
Barends *et al.*, *Nat. Com.* 6, 7654 (2015)



Raferly *et al.*, *Phys. Rev. X* 4, 031043 (2014)

Hybrid Systems with Superconducting Circuits

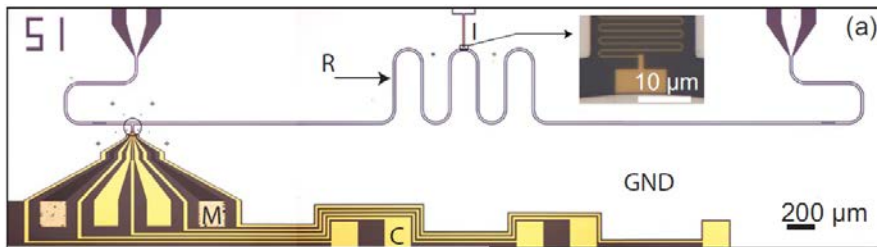
Spin Ensembles: e.g. NV centers
 D. Schuster *et al.*, *PRL* 105, 140501 (2010)
 Y. Kubo *et al.*, *PRL* 105, 140502 (2010)



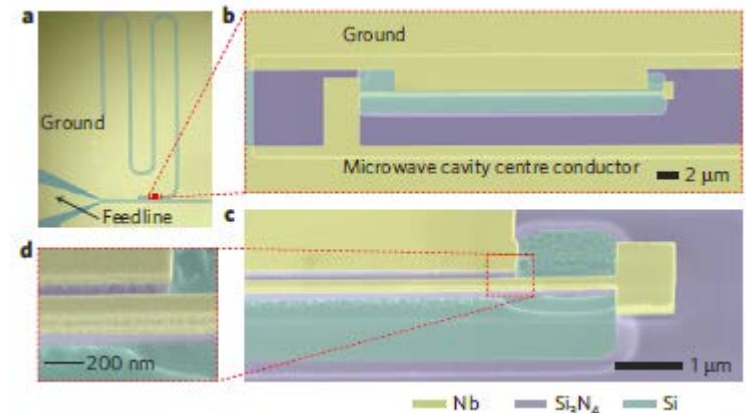
Polar Molecules, Rydberg, BEC
 P. Rabl *et al.*, *PRL* 97, 033003 (2006)
 A. Andre *et al.*, *Nat. Phys.* 2, 636 (2006)
 D. Petrosyan *et al.*, *PRL* 100, 170501 (2008)
 J. Verdu *et al.*, *PRL* 103, 043603 (2009)



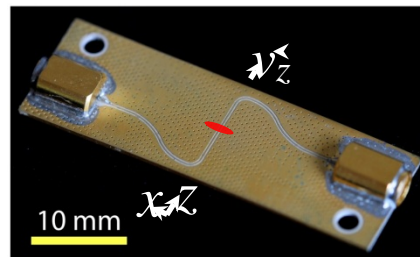
CNT, Gate Defined 2DEG, or nanowire Quantum Dots
 M. Delbecq *et al.*, *PRL* 107, 256804 (2011)
 T. Frey *et al.*, *PRL* 108, 046807 (2012)
 K. Petersson *et al.*, *Nature* 490, 380 (2013)



Nano-Mechanics
 J. Teufel *et al.*, *Nature* 475, 359 (2011)
 X. Zhou *et al.*, *Nat. Phys.* 9, 179 (2013)



Rydberg Atoms
 S. Hogan *et al.*, *PRL* 108, 063004 (2012)

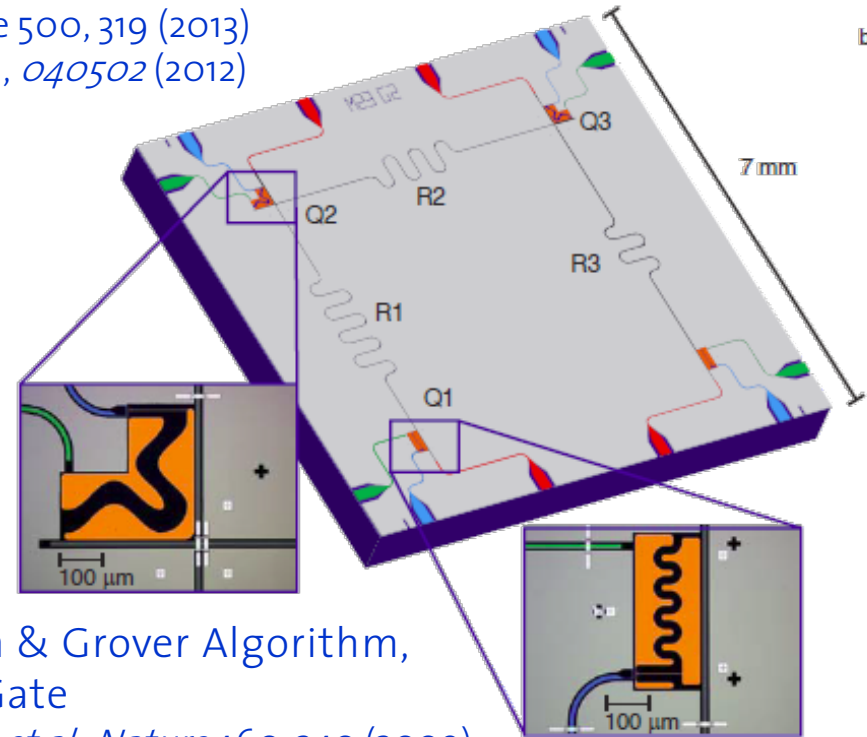


... and many more

Quantum Computing with Superconducting Circuits

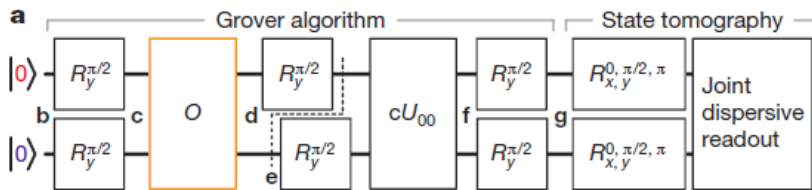
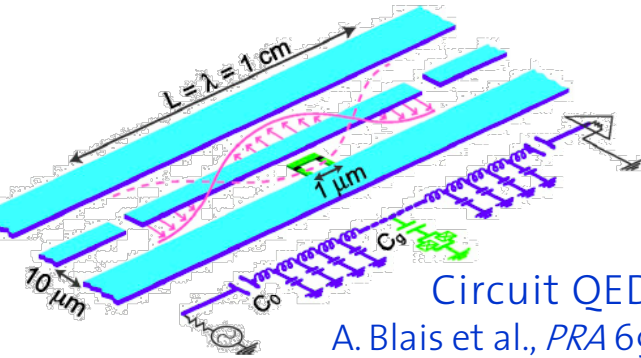
Teleportation

L. Steffen *et al.*, *Nature* 500, 319 (2013)
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 A. Wallraff *et al.*, *Nature* 431, 162 (2004)
 M. Sillanpaa *et al.*, *Nature* 449, 438 (2007)
 H. Majer *et al.*, *Nature* 449, 443 (2007)
 M. Mariani *et al.*, *Science* 334, 61 (2011)
 R. Barends *et al.*, *Nature* 508, 500 (2014)

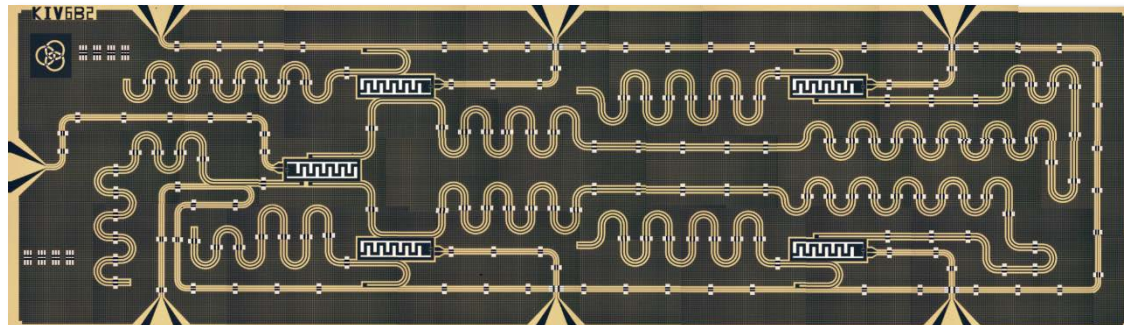


Deutsch & Grover Algorithm, Toffoli Gate

L. DiCarlo *et al.*, *Nature* 460, 240 (2009)
 L. DiCarlo *et al.*, *Nature* 467, 574 (2010)
 A. Fedorov *et al.*, *Nature* 481, 170 (2012)

Error Correction

M. Reed *et al.*, *Nature* 481, 382 (2012)
 Corcoles *et al.*, *Nat. Com.* 6, 6979 (2015)
 Ristè *et al.*, *Nat. Com.* 6, 6983 (2015)
 Kelly *et al.*, *Nature* 519, 66-69 (2015)



Teleportation ... what one may wish for !?



Teleportation in the Quantum World

Objective:

- transfer information stored in a quantum bit from a sender to receiver

Resources:

- a pair of entangled qubits shared between the sender and receiver
- a small quantum computer at the sender and at the receiver
- a classical communication channel

Alice



classical communication

Bob



Features:

- exploits non-local quantum correlations
- uses all essential ingredients required for realizing a universal quantum computer
- full protocol demonstrates use of real-time feed-forward

Applications:

- universal quantum computation
- simplification of quantum circuits
- repeaters for quantum comm.

Has been demonstrated for photons, ions and now also in solid state systems.

Teleportation Protocol

Task:

- transfer unknown quantum state from Alice to Bob

Resources:

- a pair of entangled qubits ($Q_1 + Q_2$)

Alice



Bell measurement

Qubit: Q_1 Q_2



arbitrary unknown qubit state

entangled qubit state

Bob



Q_3



proposal: Bennett *et al.*, *Phys. Rev. Lett.* 70, 1895 (1993)

Teleportation Protocol

Task:

- transfer unknown quantum state from Alice to Bob

Resources:

- a pair of entangled qubits ($Q_1 + Q_2$)
- classical communication

Alice



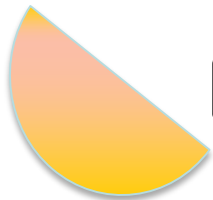
Q_1

classical communication

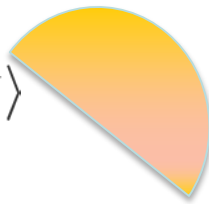
Bob



Qubit: Q_1



Q_2



$|\Psi^+\rangle$

Q_3



$|\psi\rangle$

Teleportation in other Systems

Single photons

- D. Bouwmeester, *et al.*, *Nature*, 390, 575–579 (1997)
- I. Marcikic, *et al.*, *Nature*, 421, 509–513 (2003)
- J. Yin, *et al.*, *Nature*, 488, 185–188 (2012)
- X.-S. Ma, *et al.*, *Nature*, 489, 269–273 (2012)

Ion traps

- M. Riebe *et al.*, *Nature*, 429, 734–737 (2004)
- M. Barrett, *et al.*, *Nature*, 429, 737–739 (2004)
- S. Olmschenk, *et al.*, *Science*, 323, 486–489 (2009)

Atomic ensembles

- X.-H. Bao, *et al.*, *PNAS*, 109, 20347 (2012)

Single atoms

- C. Nölleke, *et al.*, *Phys. Rev. Lett.*, 110, 140403 (2013)

NMR

- M. A. Nielsen, *et al.*, *Nature*, 396, 52–55 (1998)

Continuous variables

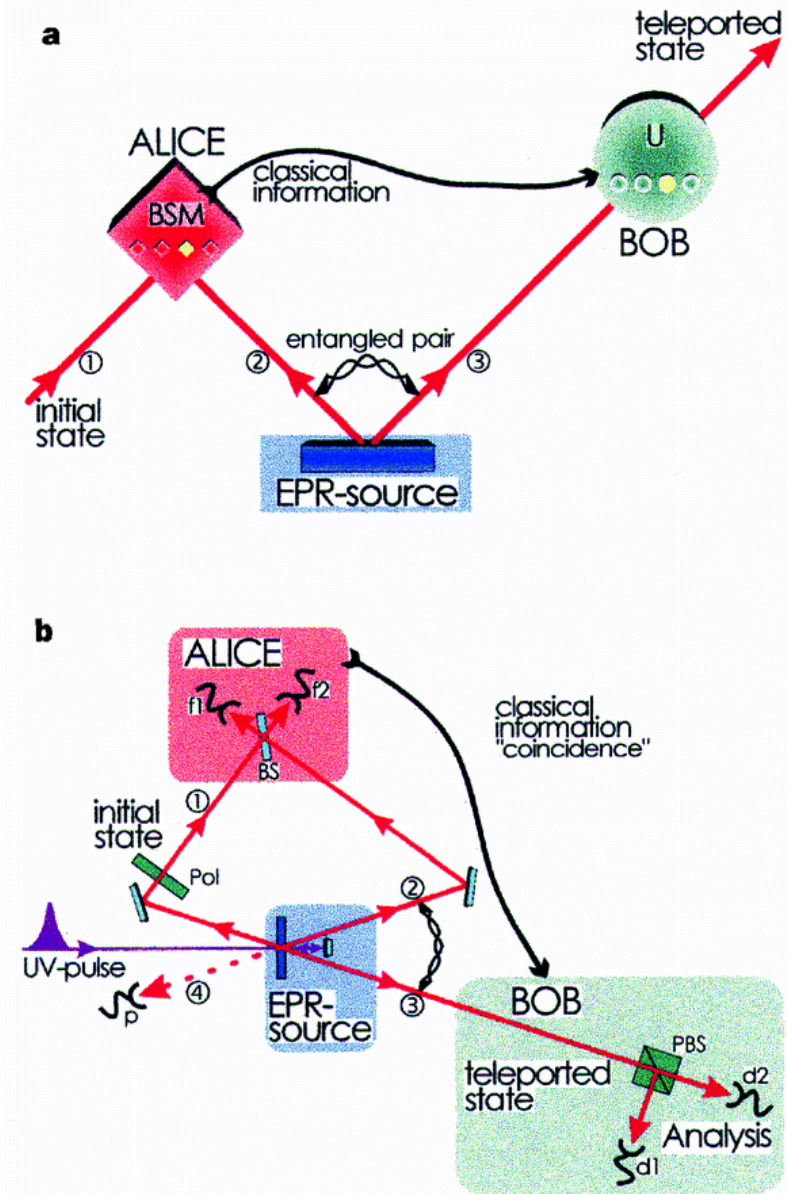
- A. Furusawa, *et al.*, *Science*, 282, 706–709 (1998)
- N. Lee, *et al.*, *Science*, 332, 330–333 (2011)
- S. Takeda, *et al.*, *Nature*, 500, 315–318 (2013)

Semiconductor Quantum Dots

- W.B. Gao, *et al.*, *Nat. Comm* 4, 2744 (2013)

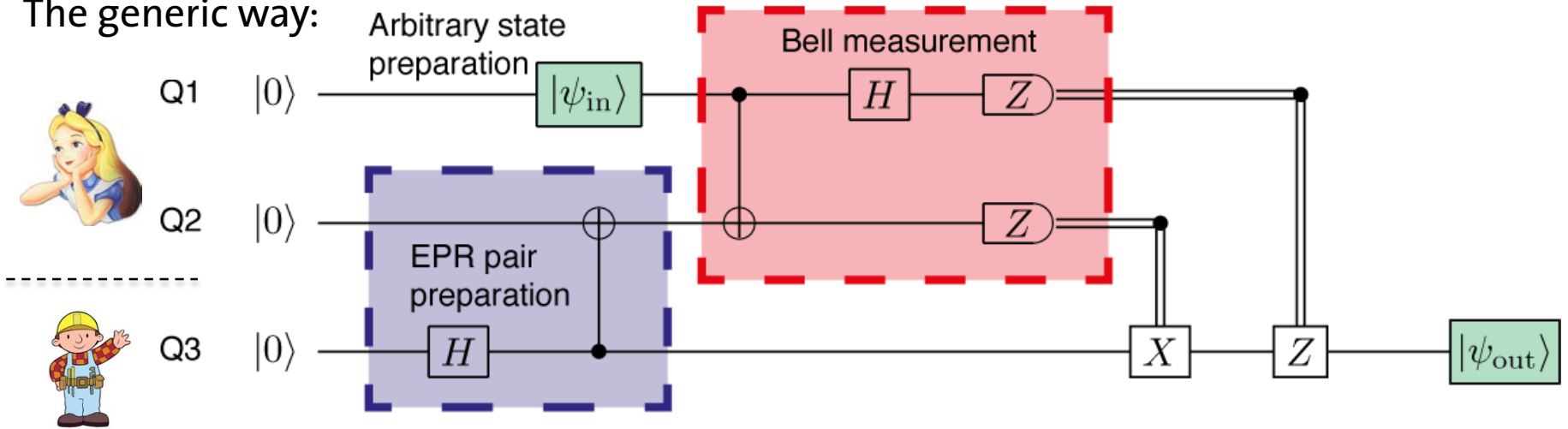
NV Centers

- W. Pfaff, *et al.*, *Science* 345, 532 (2014)



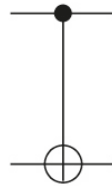
Implementation of the Teleportation Protocol

The generic way:



Hadamard

Rotation around Y-axis



Controlled NOT

Controlled phase gate



Measurement along Z-axis

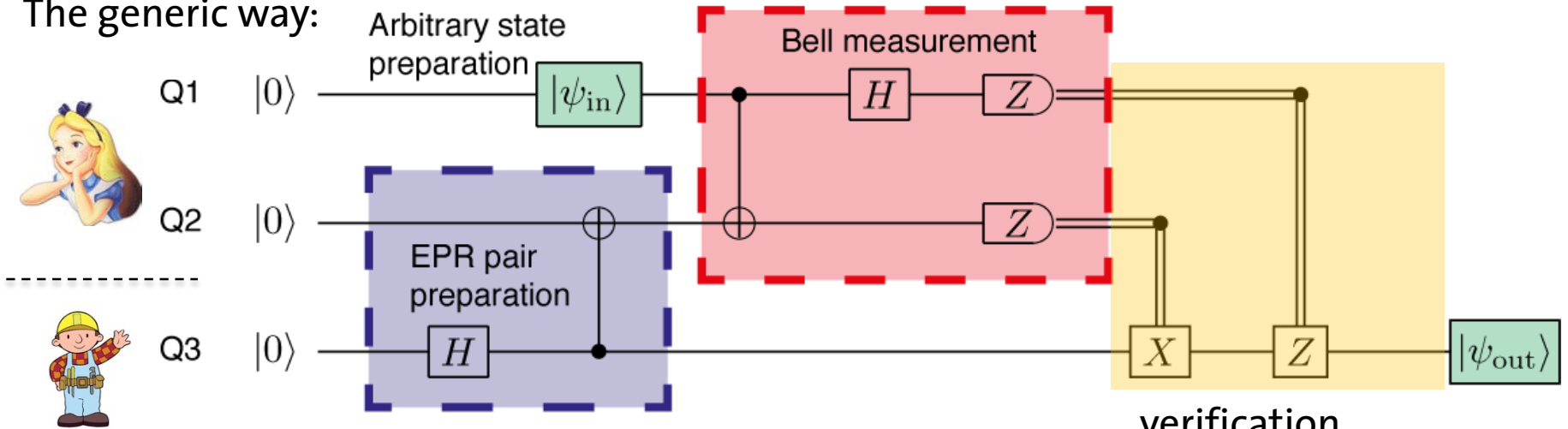
$$\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$

proposal: F. W. Strauch, *Phys. Rev. Lett.* 91, 167005 (2003).

implementation: L. DiCarlo, *Nature* 460, 240 (2010).

Implementation of the Teleportation Protocol

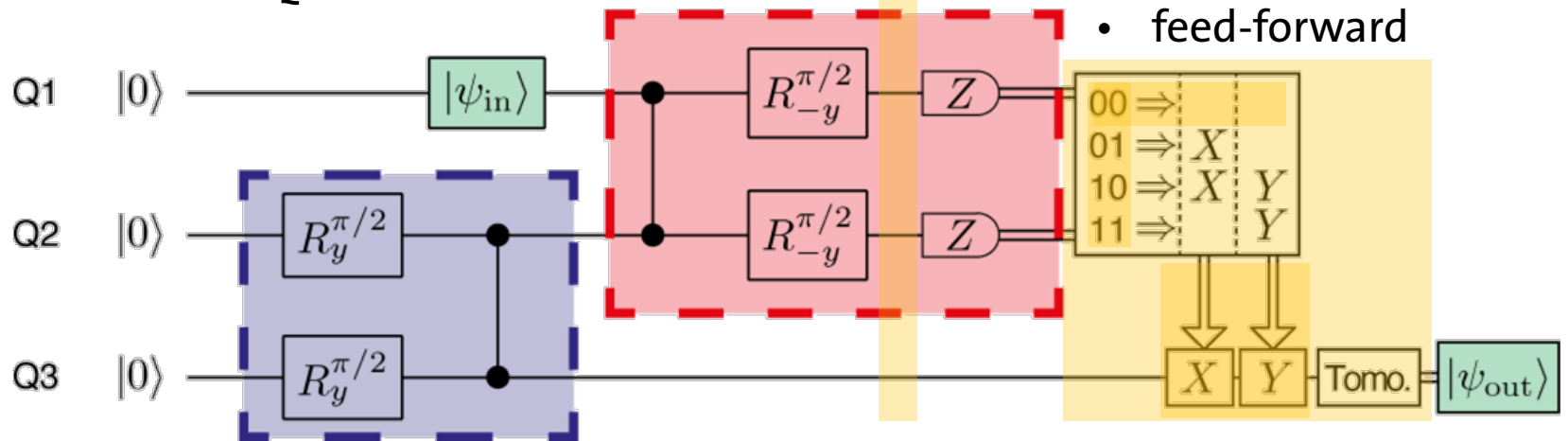
The generic way:



verification

- 3-qubit tomography
- post selection (1 state)
- deterministic (4 states)
- feed-forward

Realization in circuit QED:



Steffen *et al.*, *Nature* 500, 319 (2013)

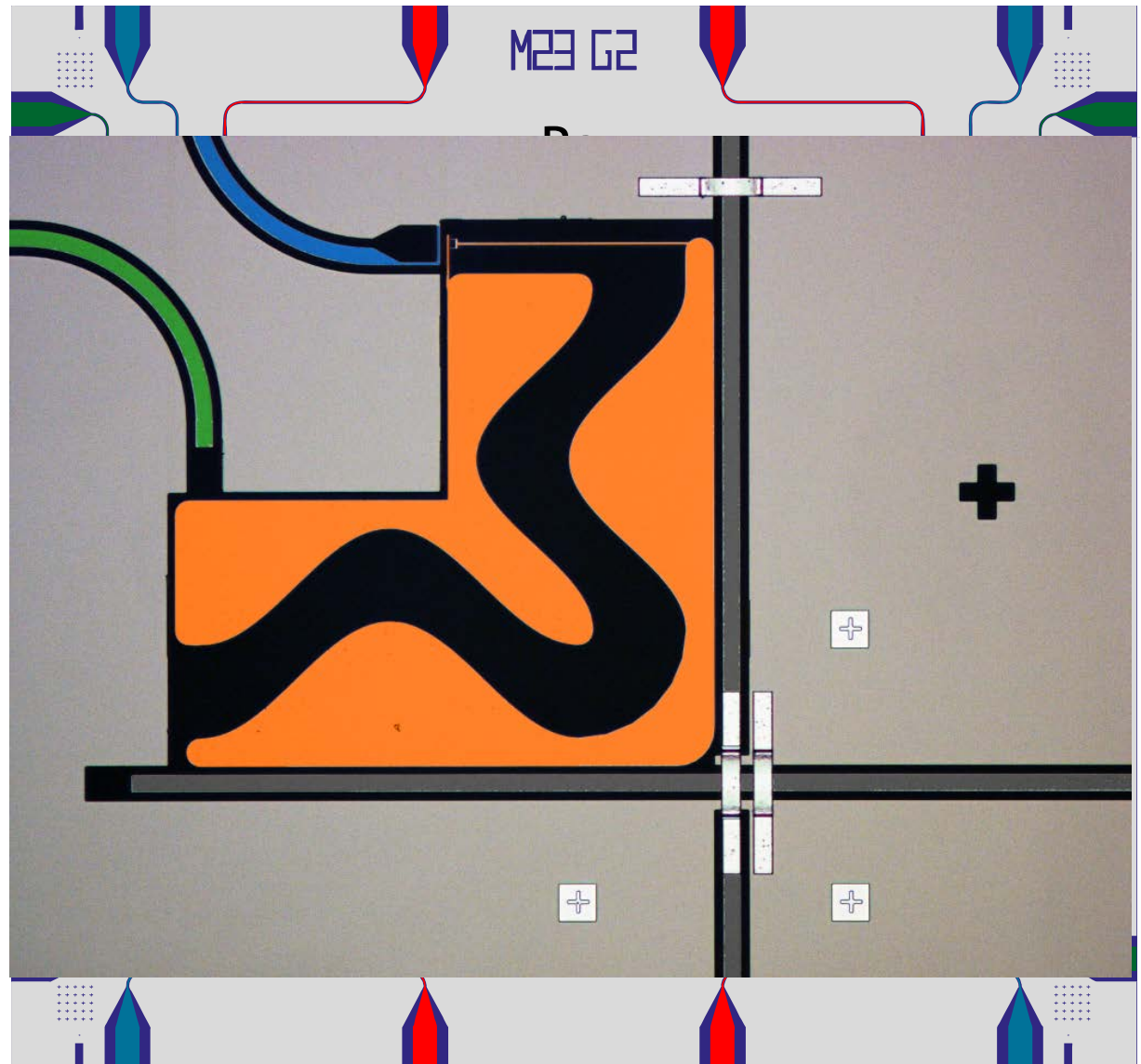
The Sample

- 3 Resonators
- 3 Qubits
- single-qubit gates
- two-qubit gates (qubits in the same resonator)
- joint single-shot readout of qubits 1 & 2
- single-shot readout of qubit 3
- with two parametric amplifiers

Yurke and Buks, *J. Lightwave Tech.* 24, 5054 (2006).

Castellanos-Beltran et al., *Nat. Phys.* 4, 929 (2008).

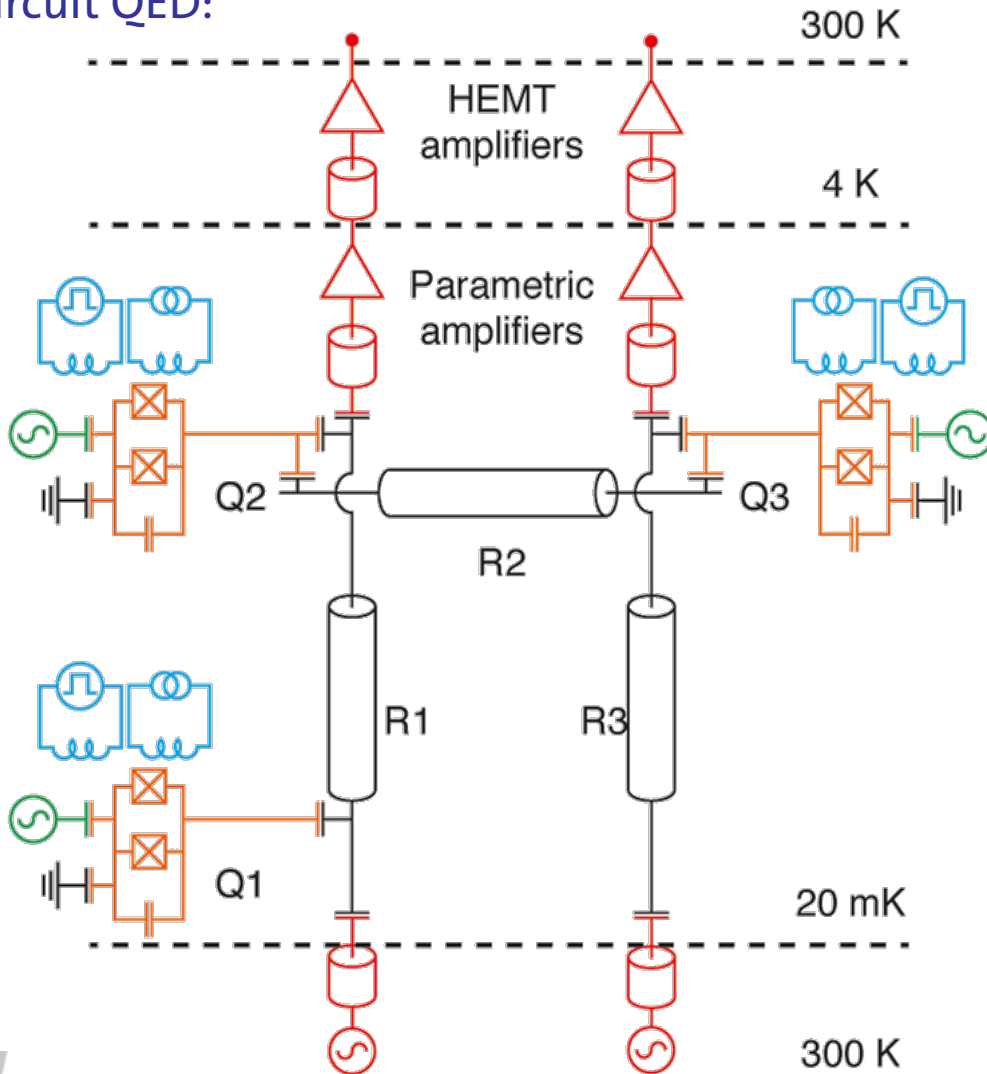
Eichler et al., *PRL* 107, 113601 (2011).



Steffen et al., *Nature* 500, 319 (2013)

The Circuit

12-port quantum device based
on circuit QED:

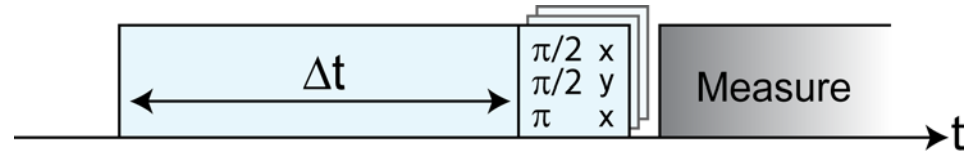


Device highlights:

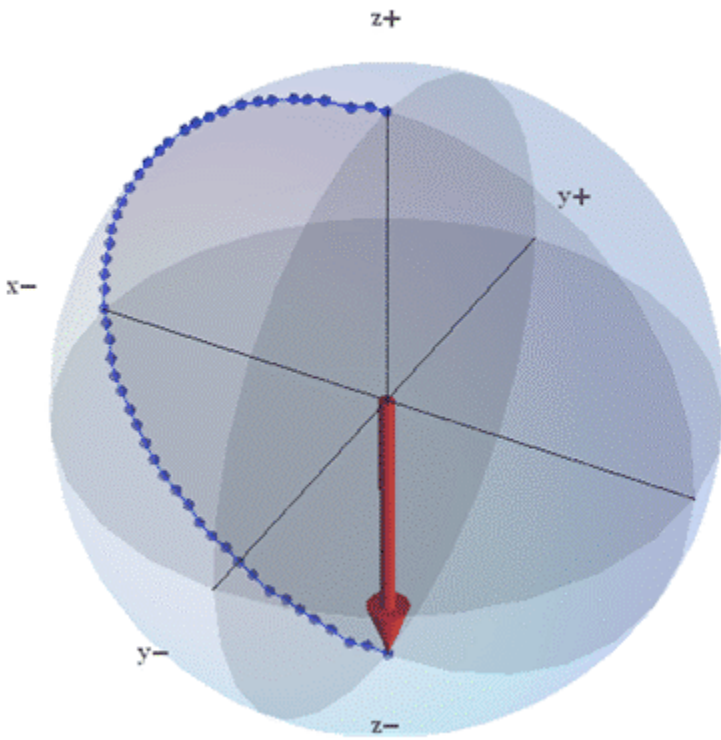
- 3 high-Q resonators
- 4 transmon **qubits**
- individual control of all qubits
- nearest neighbor interaction via quantum bus
- individual read-out for pairs of **qubits 1-2** and **3-4** through resonators
- single-shot read-out using parametric amplifiers
- qubit separation ~ 10 mm
- cross-overs for resonators

Single Qubit Gates

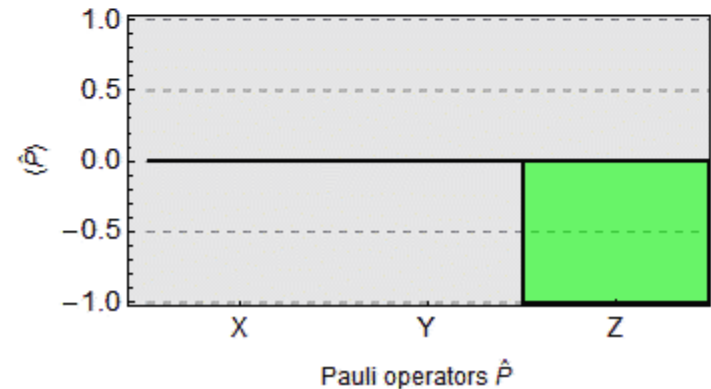
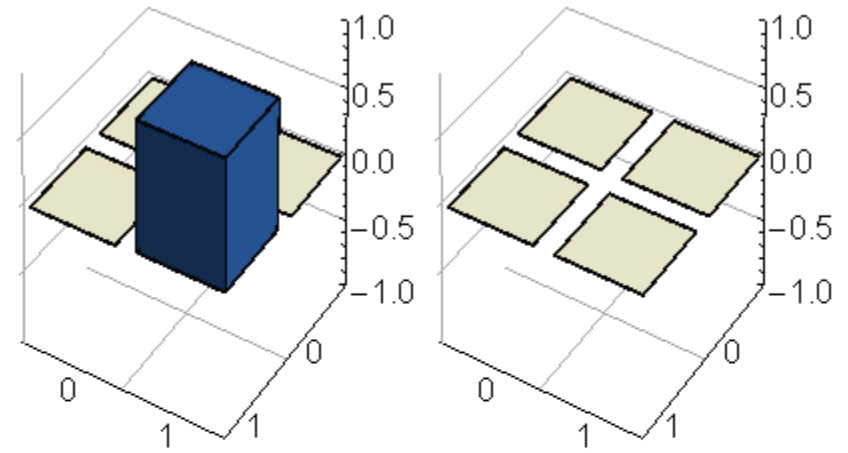
Pulse sequence for qubit rotation and readout:



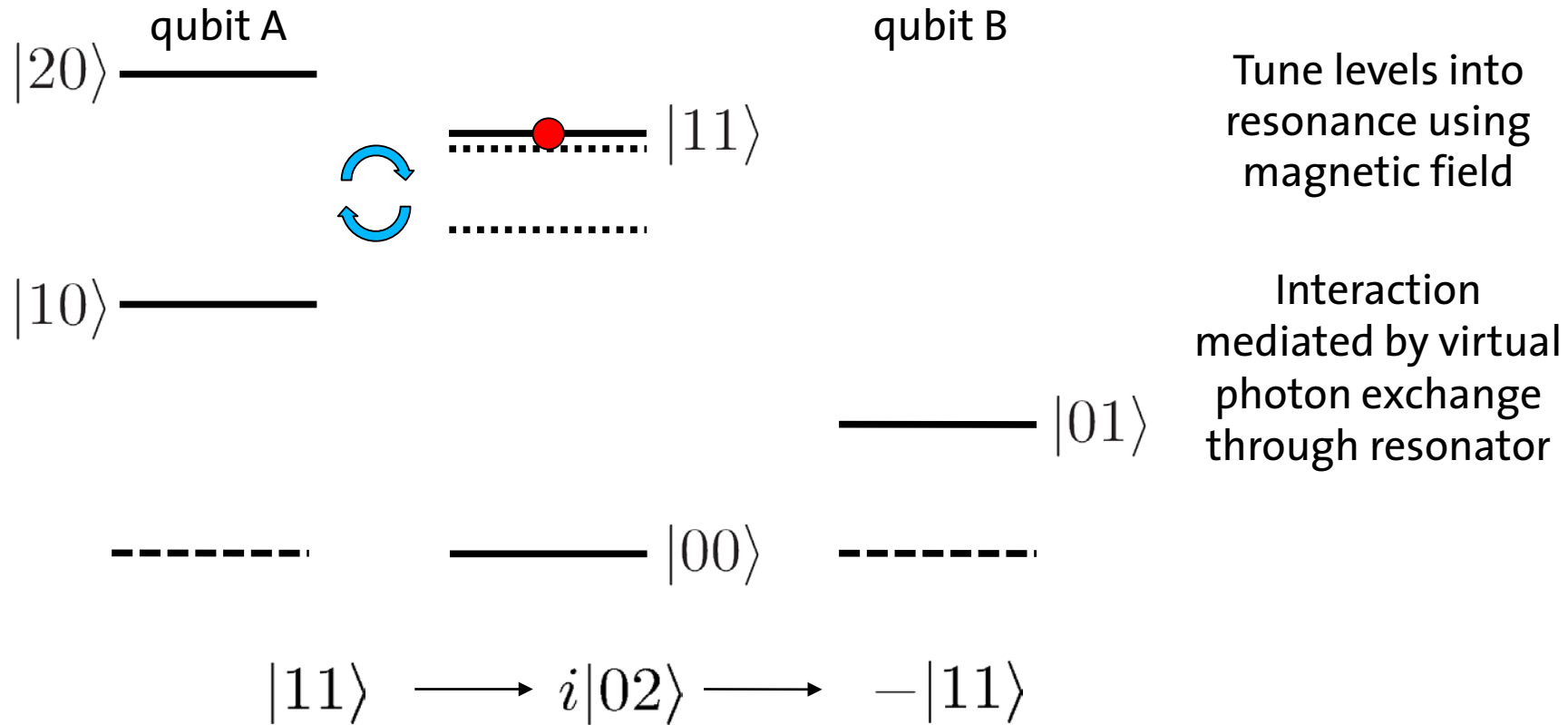
experimental Bloch vector:



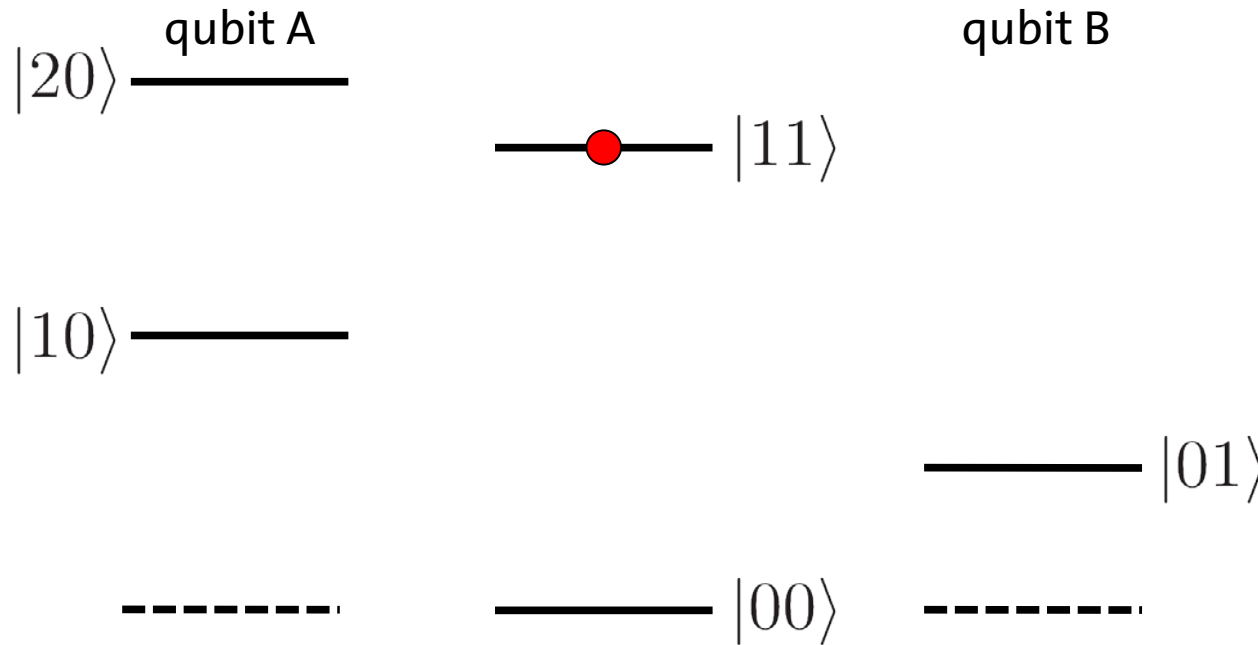
experimental density matrix and Pauli set:



Controlled Phase Gate



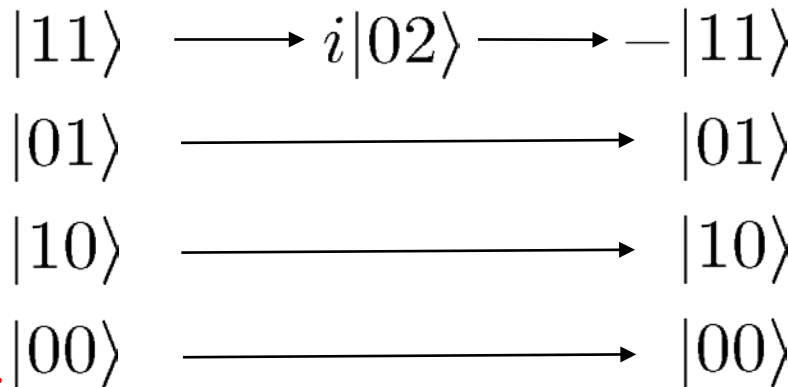
Controlled Phase Gate



Tune levels into resonance using magnetic field

Interaction mediated by virtual photon exchange through resonator

This experiment:
 Universal two-qubit gates: $F \sim 90\%$.
 Single-qubit gates: $F \sim 98\%$.
 .. to realize needed quantum operations.



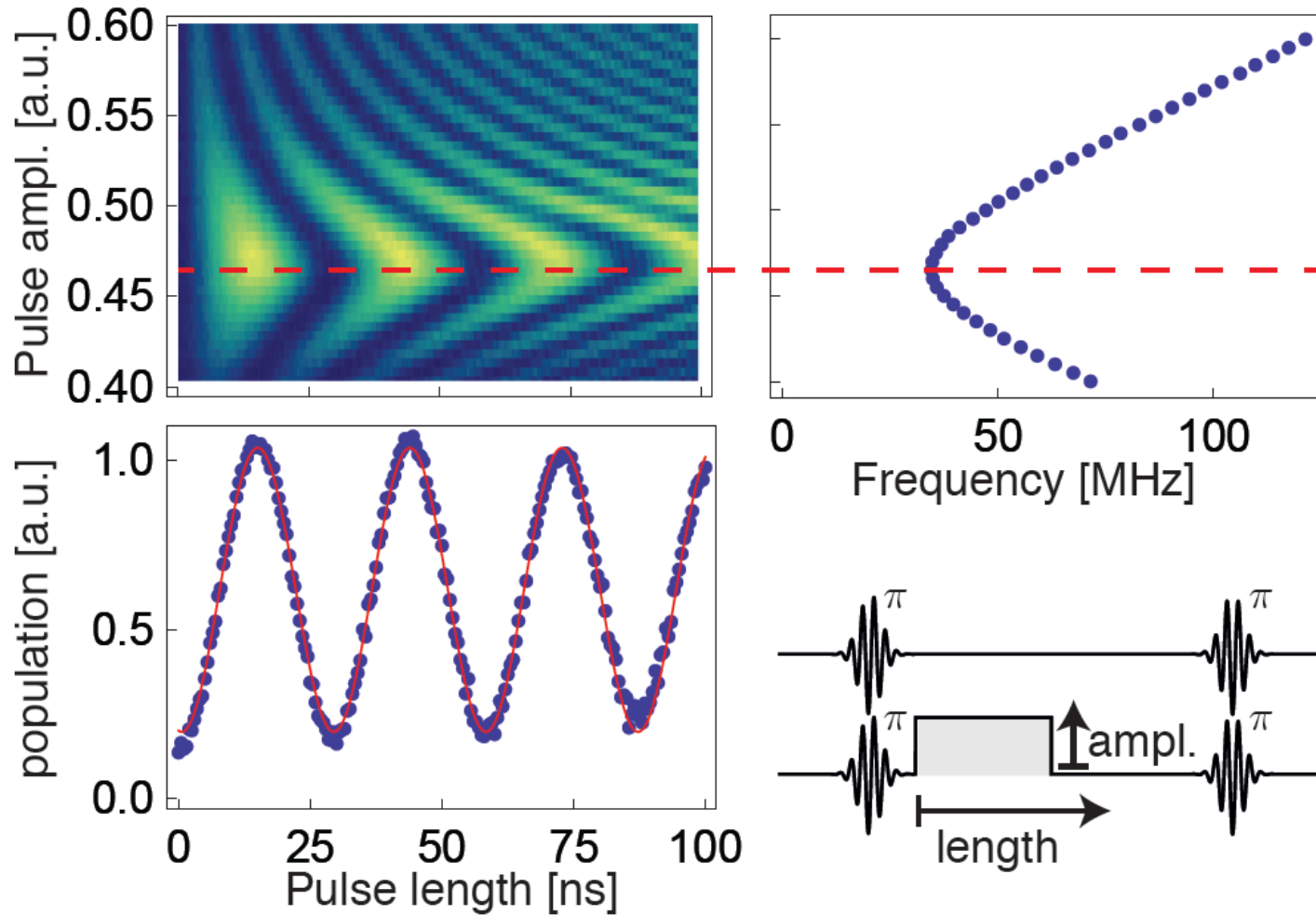
C-Phase gate:

$$\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$

proposal: F. W. Strauch *et al*, *PRL* 91, 167005 (2003)

first implementation: L. DiCarlo *et al.*, *Nature* 467, 467 (2010)

Experimental C-PHASE Gate



Process Tomography: C-Phase Gate

arbitrary quantum process

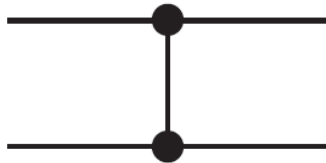
$$\rho' = \mathcal{E}(\rho)$$

decomposed into

$$\mathcal{E}(\rho) = \sum_{mn} \tilde{E}_m \rho \tilde{E}_n^\dagger \chi_{mn}$$

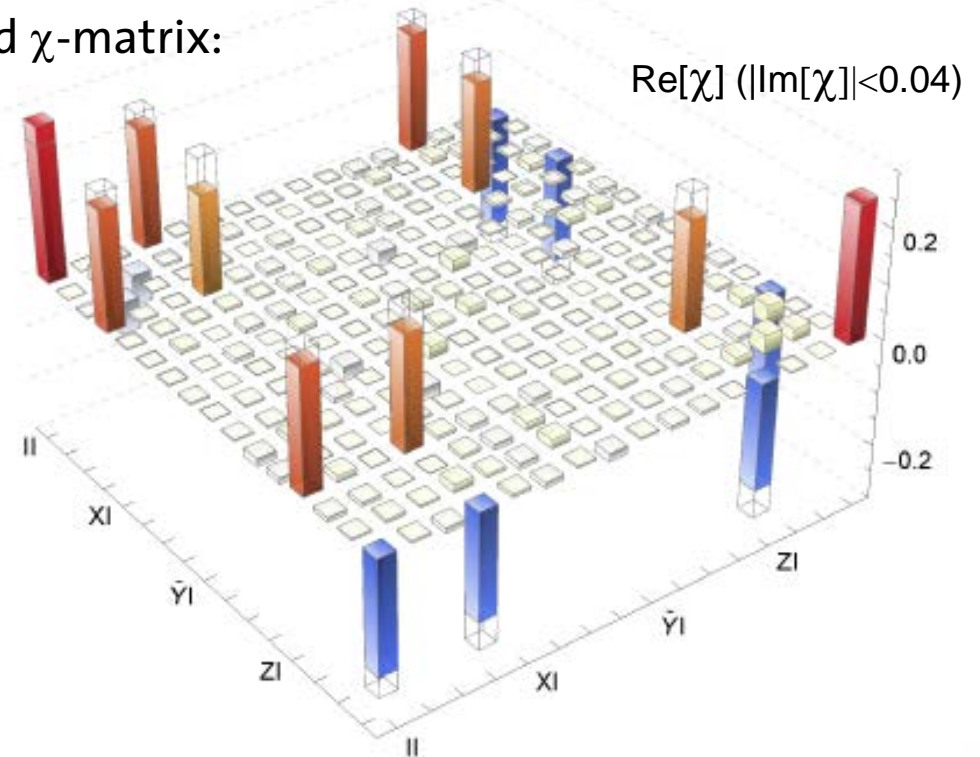
$\{\tilde{E}_k\}$ is an operator basis
 χ is a positive semi definite Hermitian matrix characteristic for the process

Controlled phase gate



$$cZ_{00} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$

Measured χ -matrix:



$$F = \text{Tr}[\chi_{\text{meas}} \chi_{\text{ideal}}] > 0.90$$

Process Tomography: C-NOT Gate

arbitrary quantum process

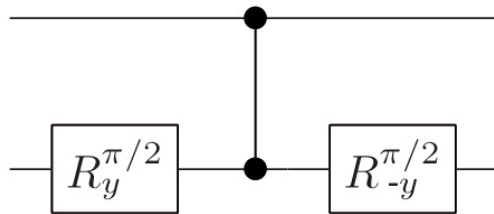
$$\rho' = \mathcal{E}(\rho)$$

decomposed into

$$\mathcal{E}(\rho) = \sum_{mn} \tilde{E}_m \rho \tilde{E}_n^\dagger \chi_{mn}$$

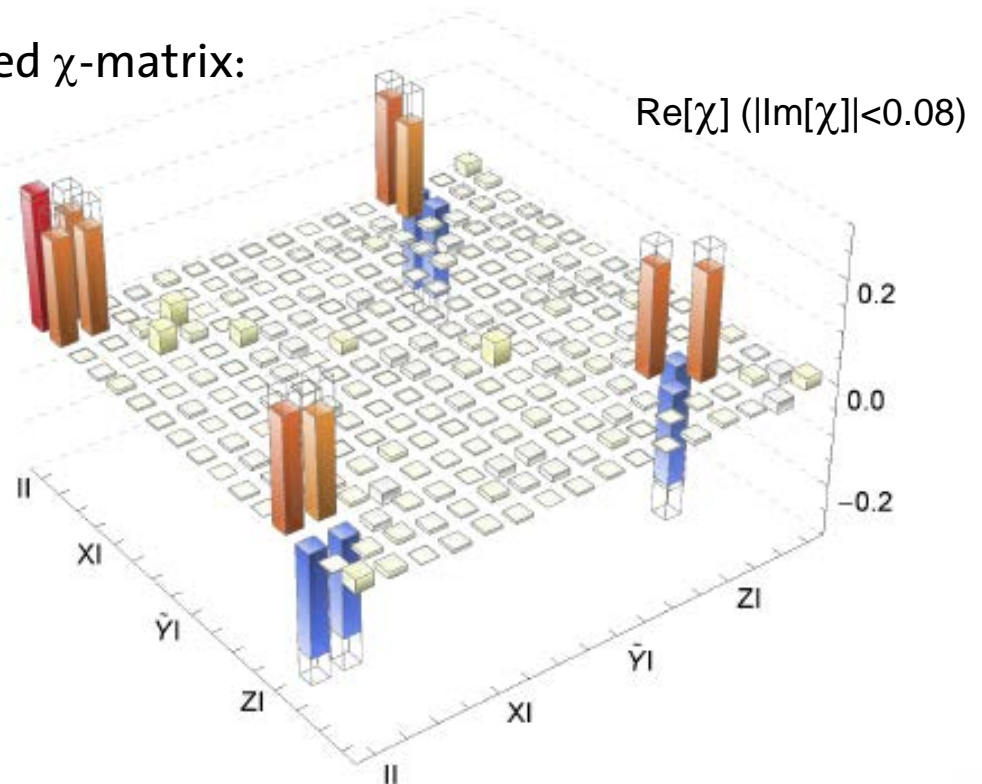
$\{\tilde{E}_k\}$ is an operator basis
 χ is a positive semi definite Hermitian matrix characteristic for the process

Controlled-NOT gate



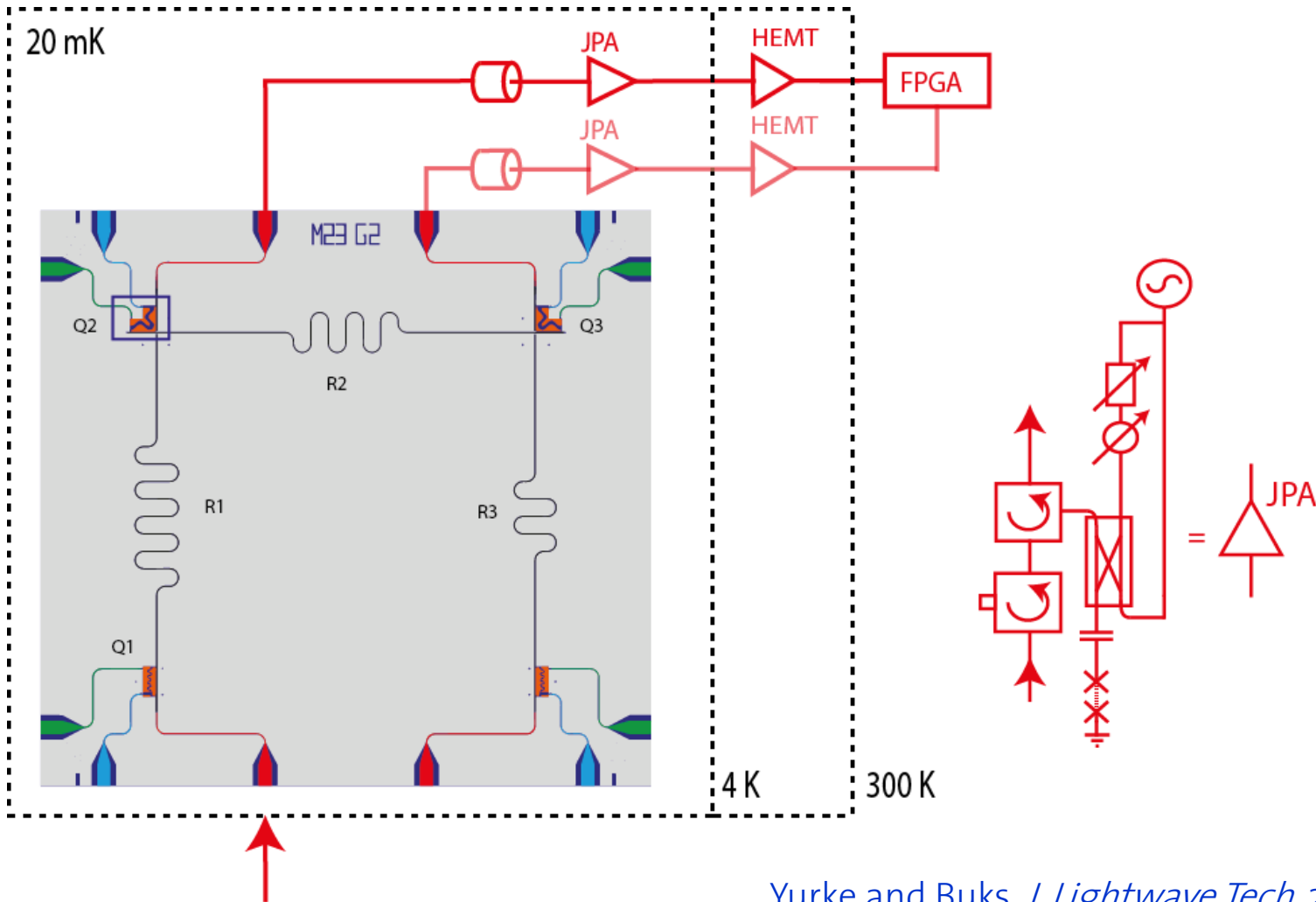
$$\text{C - NOT} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$$

Measured χ -matrix:



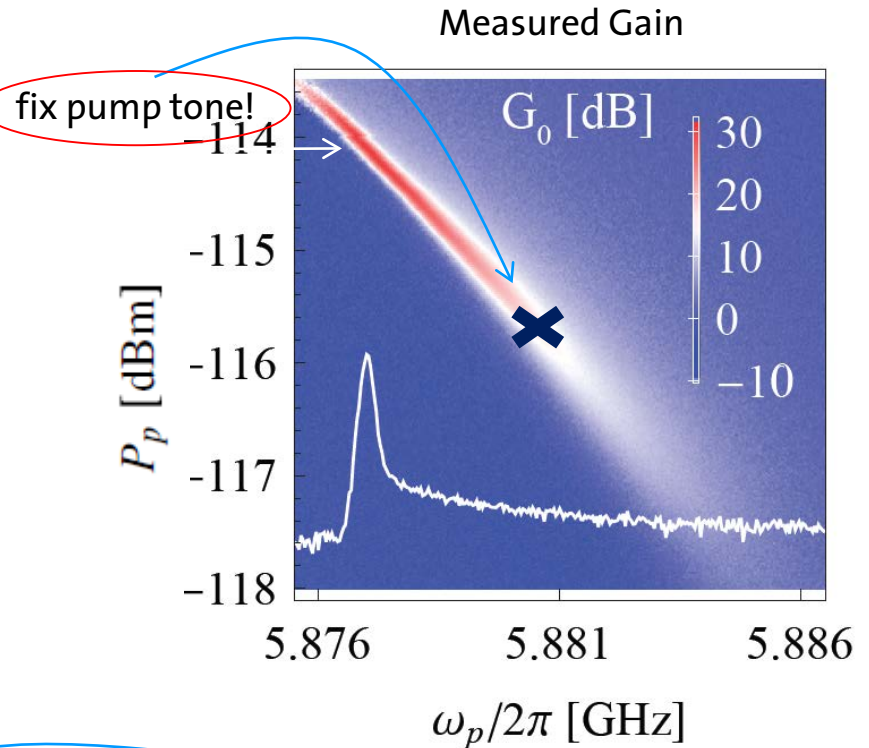
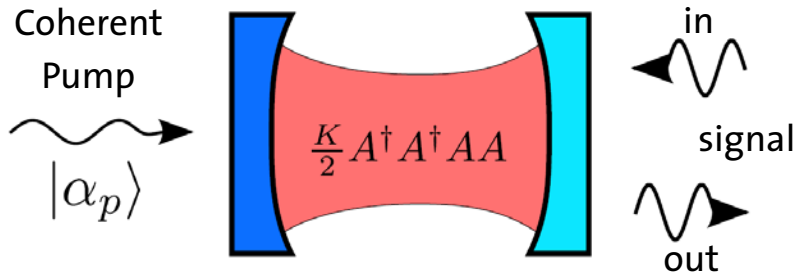
$$F = \text{Tr}[\chi_{\text{meas}} \chi_{\text{ideal}}] = 0.81$$

Dispersive Qubit Readout with Parametric Amplifiers

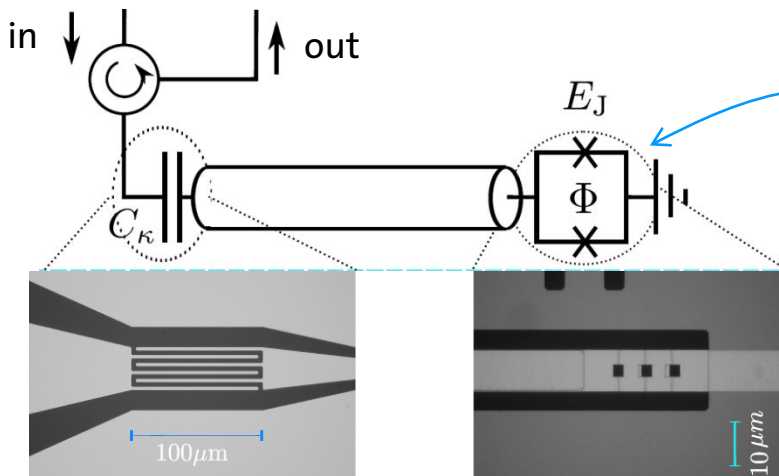


Yurke and Buks, *J. Lightwave Tech.* 24, 5054 (2006)
Castellanos-Beltran et al., *Nat. Phys.* 4, 929 (2008)
Eichler et al., *PRL* 107, 113601 (2011)
R. Vijay et al., *PRL* 106, 110502 (2011)

Near Quantum-Limited Parametric Amplifier



Circuit QED implementation:



SQUID(-array) provides required nonlinearity

Eichler *et al.*, EPJ Quantum Technology 1, 2 (2014)

Eichler *et al.*, *Phys. Rev. Lett.* 107, 113601 (2011)

Eichler *et al.*, *Phys. Rev. Lett.* 113, 110502 (2014)

Caves, *Phys. Rev. D* 26, 1817 (1982)

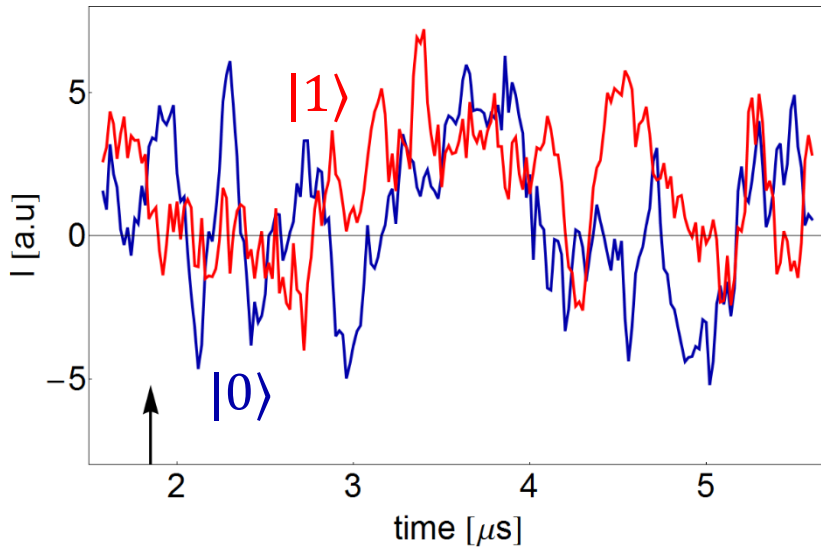
Yurke and Buks, *J. Lightwave Tech.* 24, 5054 (2006)

Castellanos-Beltran *et al.*, *Nat. Phys.* 4, 929 (2008)

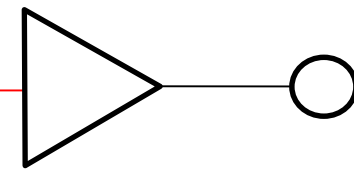
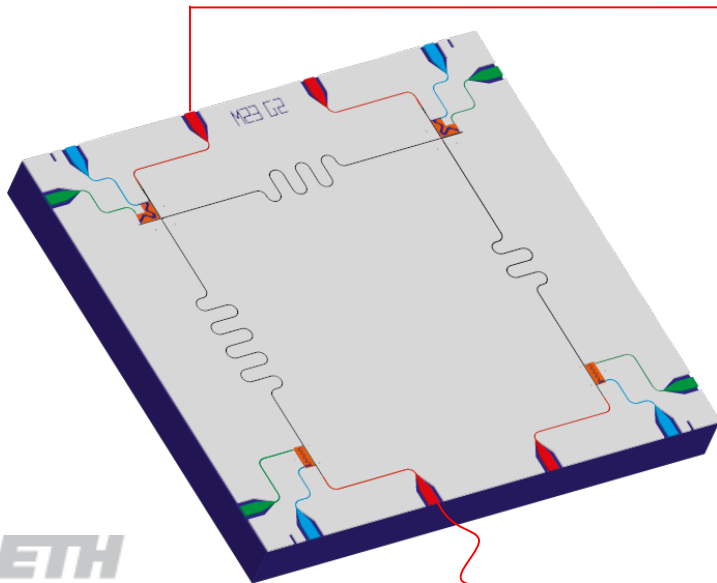
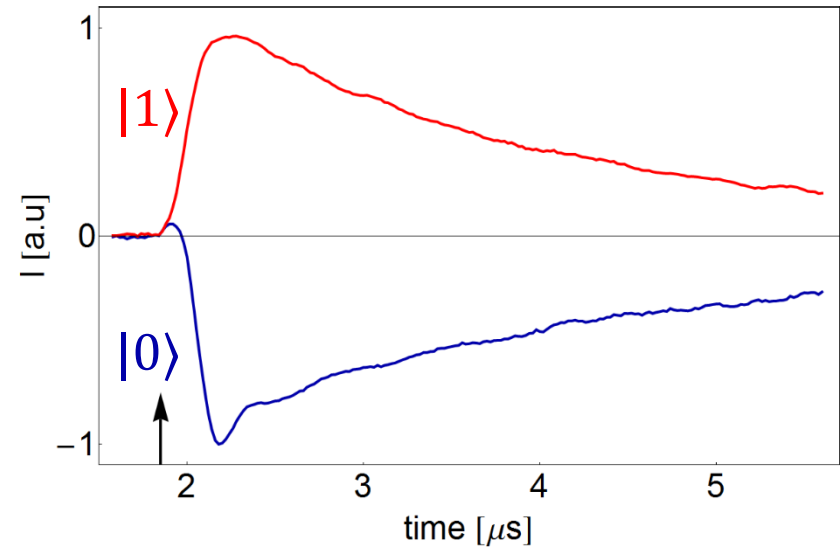
Single-Shot Single-Qubit Readout

single-shot measurements:

Conventional HEMT



averaged measurements ($8 \cdot 10^4$):



HEMT amplifier 4 K

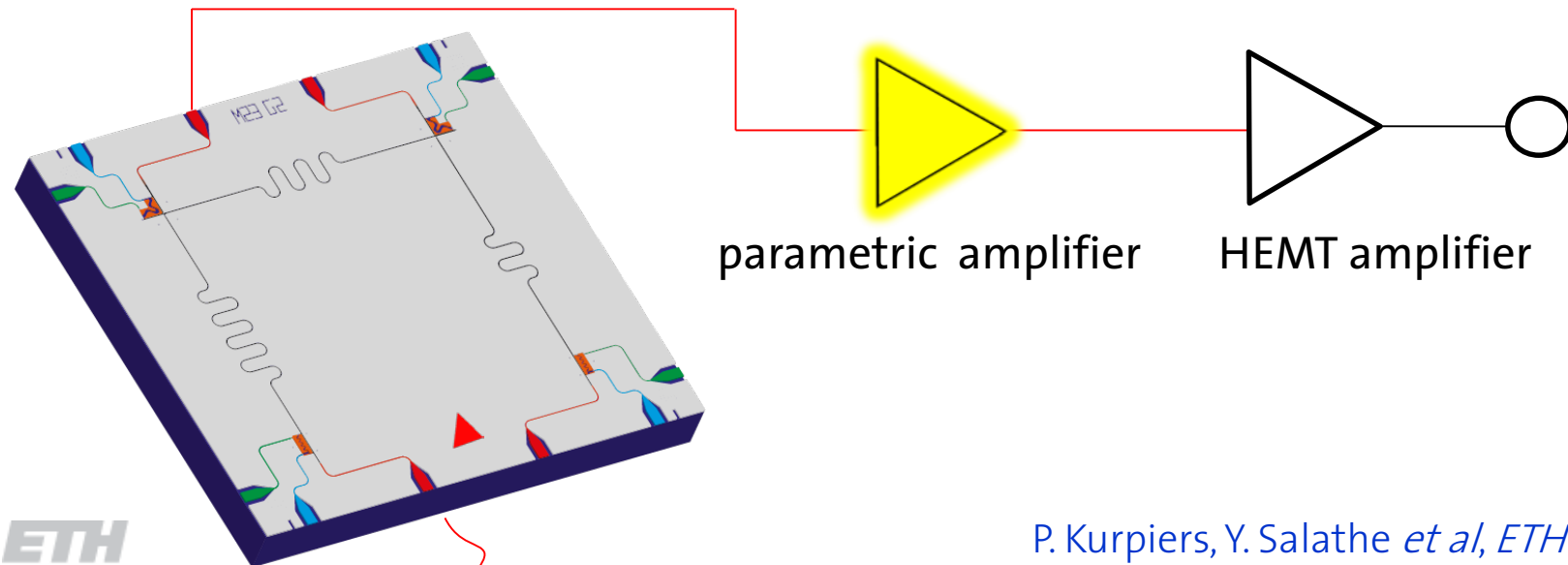
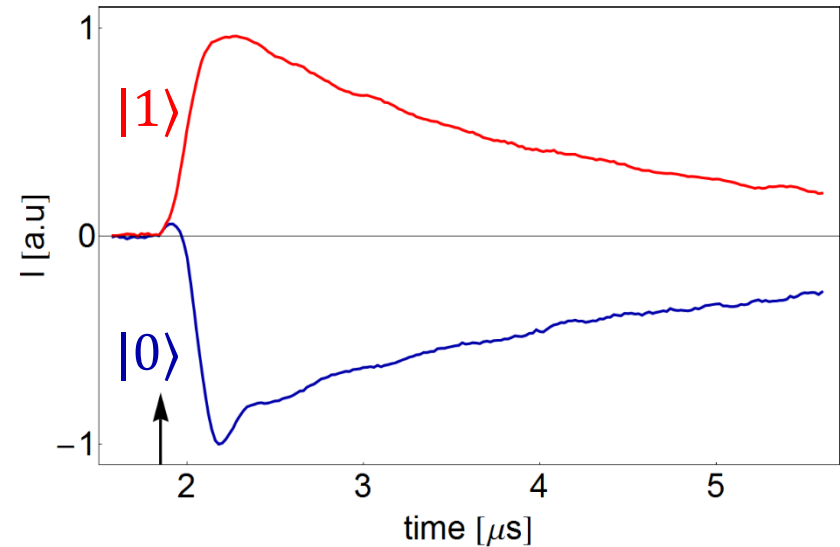
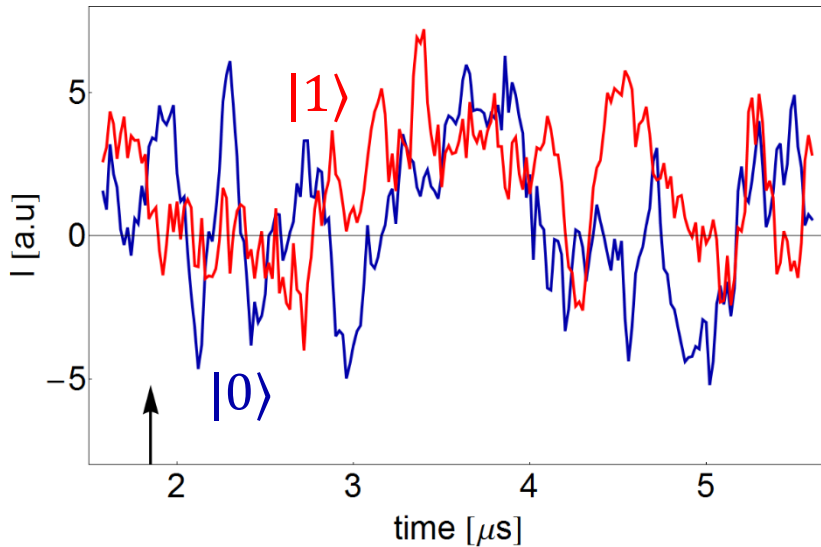
P. Kurpiers, Y. Salathe *et al*, *ETH Zurich* (2013)
R. Vijay *et al.*, *PRL* 106, 110502 (2011)

Single-Shot Single-Qubit Readout

single-shot measurements:

averaged measurements ($8 \cdot 10^4$):

Conventional HEMT



parametric amplifier

HEMT amplifier

P. Kurpiers, Y. Salathe *et al*, *ETH Zurich* (2013)

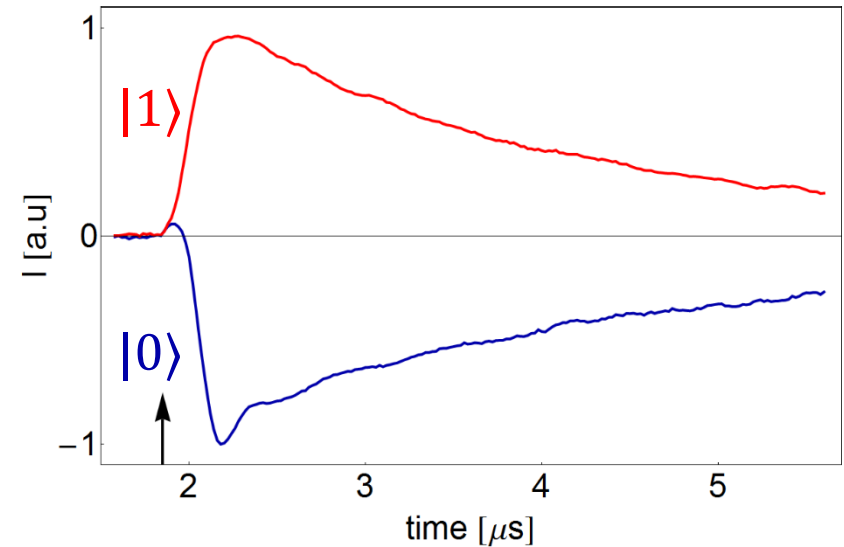
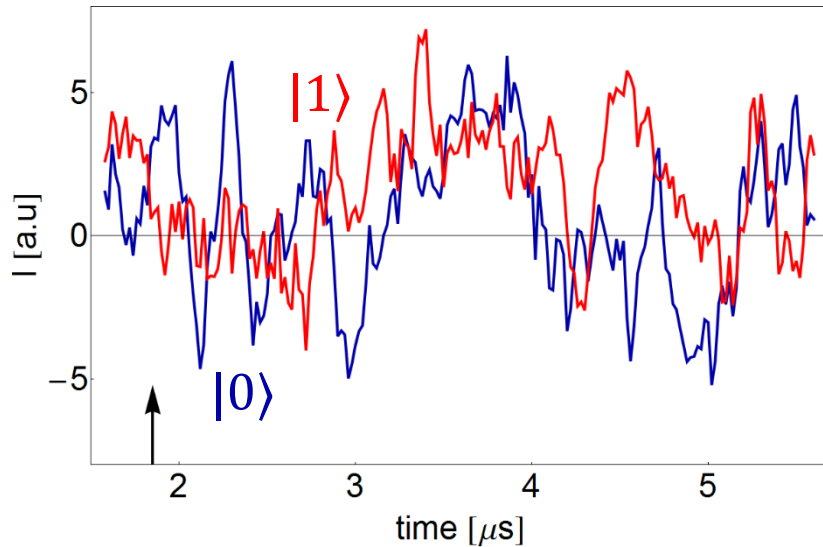
R. Vijay *et al*, *PRL* 106, 110502 (2011)

Single-Shot Single-Qubit Readout

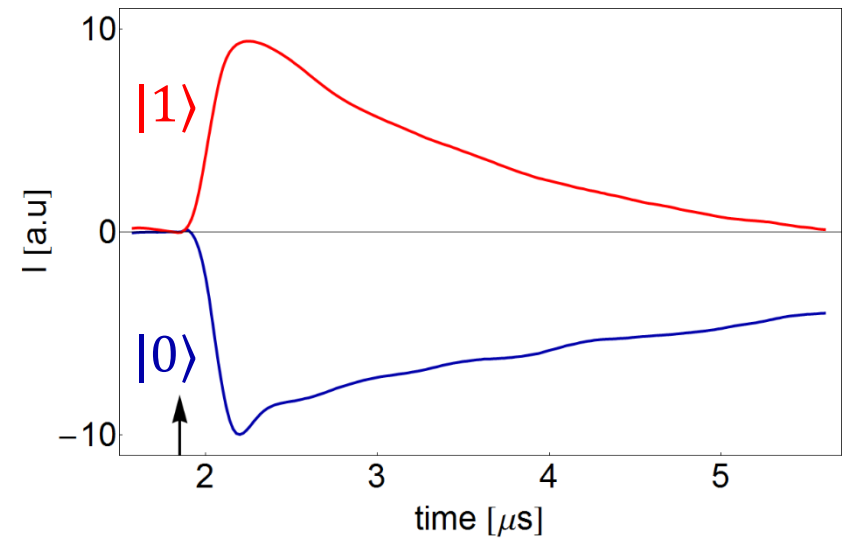
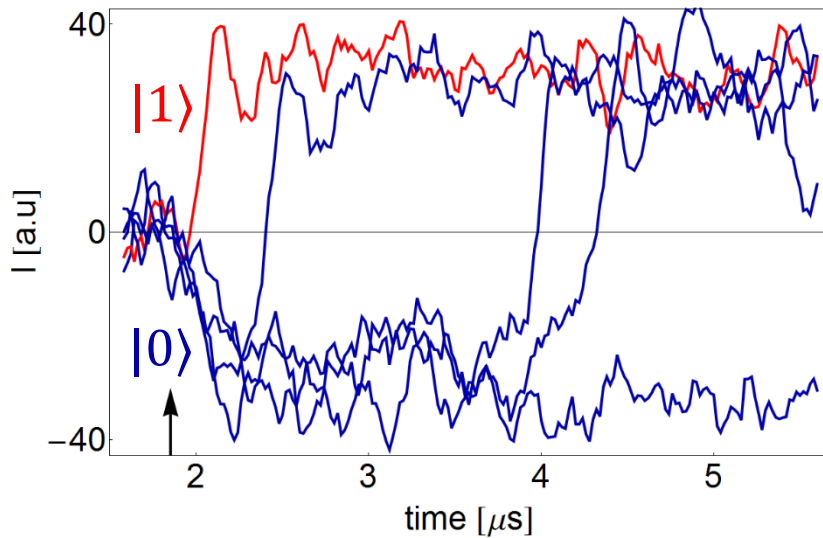
single-shot measurements:

averaged measurements ($8 \cdot 10^4$):

Conventional HEMT

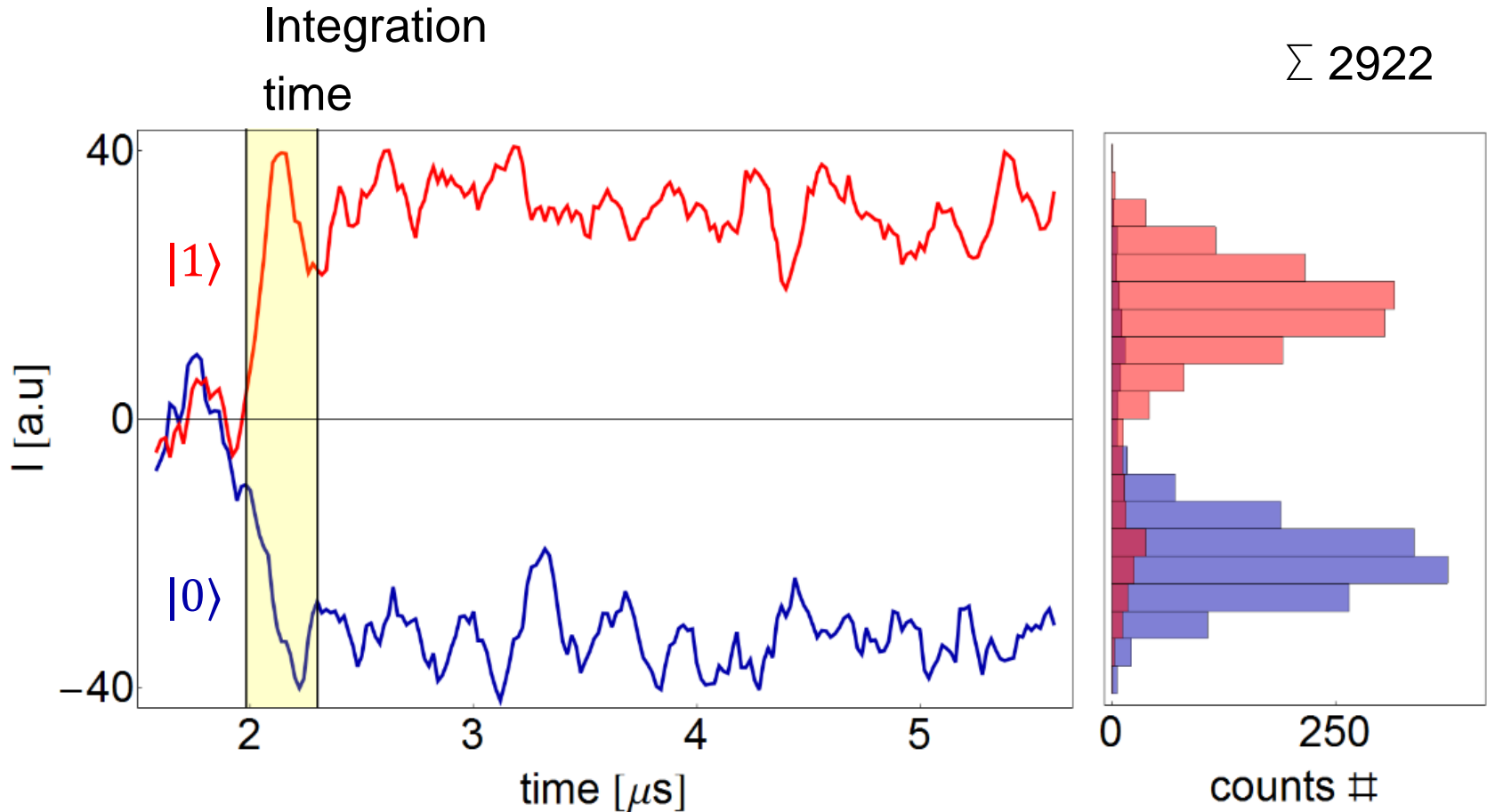


Parametric Amplifier

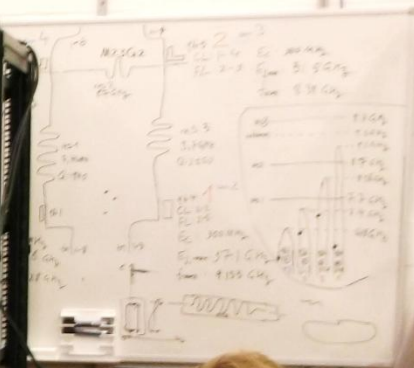


P. Kurpiers, Y. Salathe *et al*, *ETH Zurich* (2013)
R. Vijay *et al.*, *PRL* 106, 110502 (2011)

Statistics of Integrated Single-Shot Readout



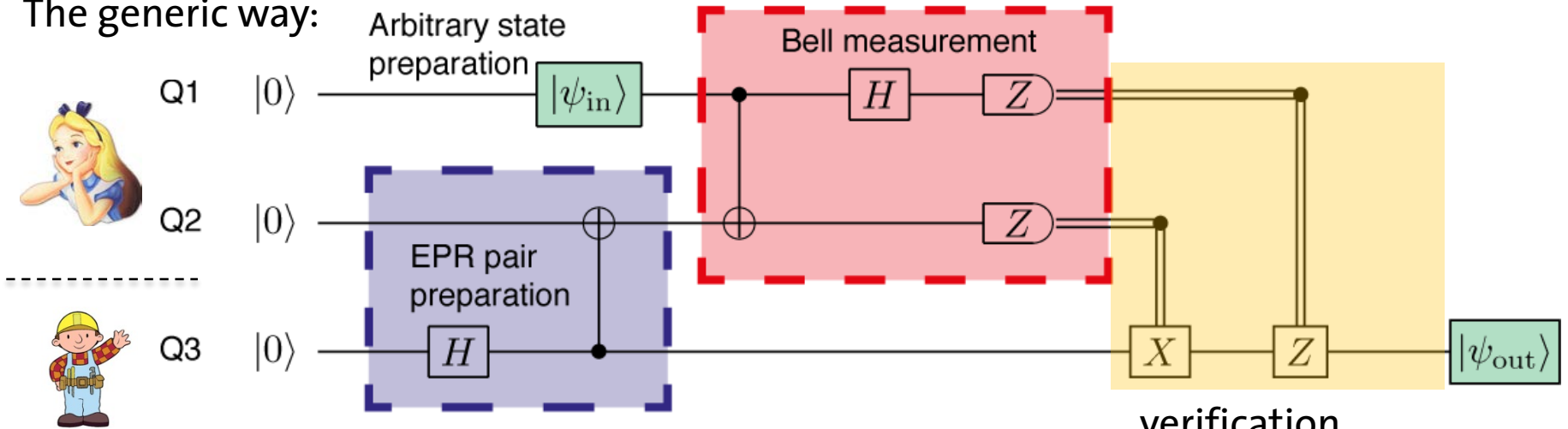
Quantum Teleportation via Superconducting Circuits



1 μm V_{cos} 2.4 GHz
10 μm V_{cos} 2.4 GHz
100 μm V_{cos} 2.4 GHz
1 mV \sin 2.4 GHz
10 mV \cos 2.4 GHz

Implementation of the Teleportation Protocol

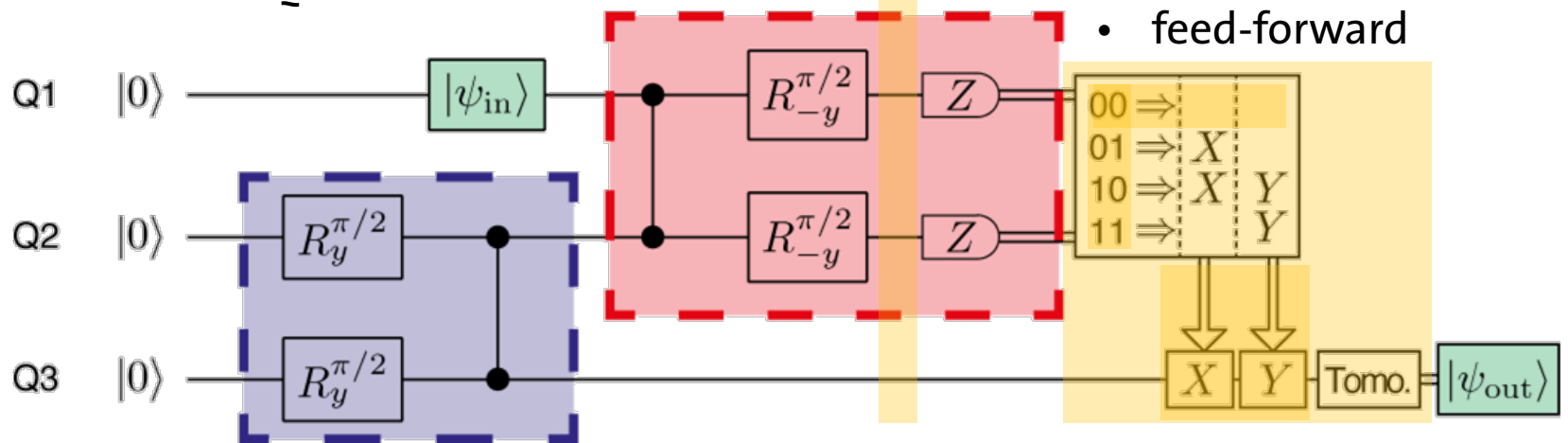
The generic way:



verification

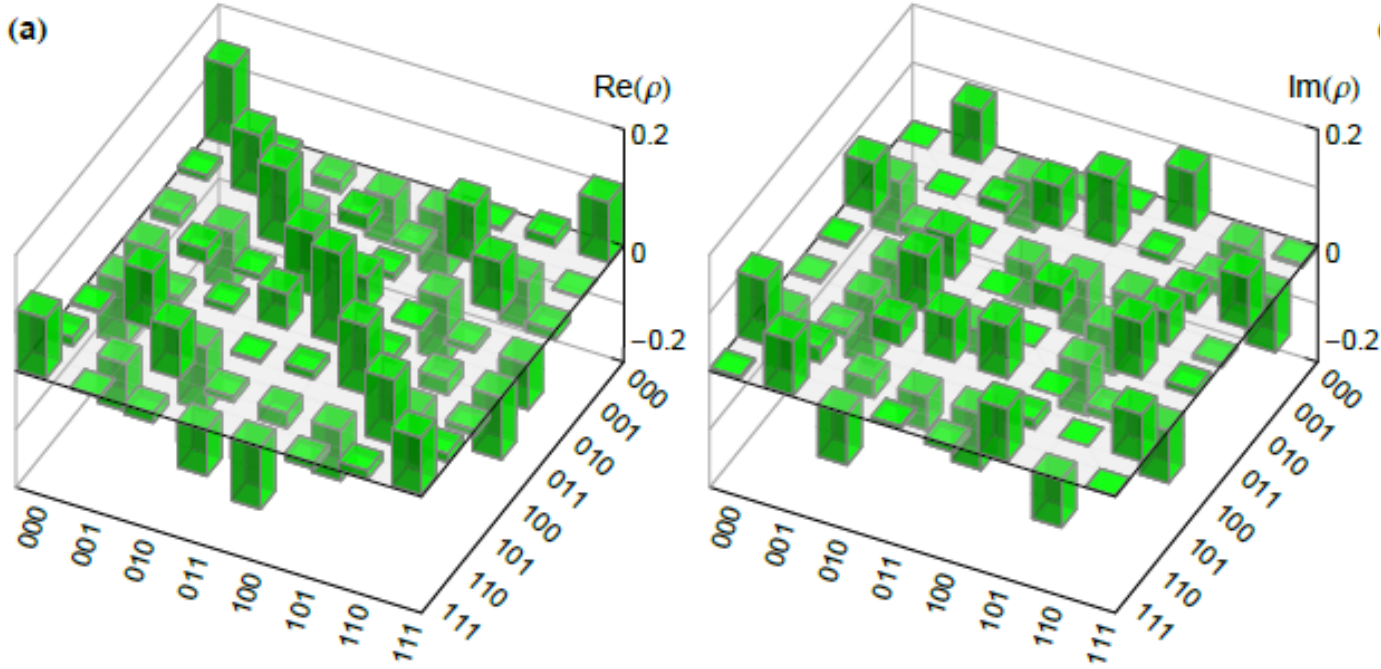
- 3-qubit tomography
- post selection (1 state)
- deterministic (4 states)
- feed-forward

Realization in circuit QED:



Steffen *et al.*, *Nature* 500, 319 (2013)

Characterizing the Full System State



measured 3-qubit entangled state after execution of circuit:

$$\text{input state: } |\psi\rangle = (|0\rangle - i|1\rangle)/\sqrt{2}$$

density matrix of
teleported state
when projecting
qubits A & B onto
 $|00\rangle$

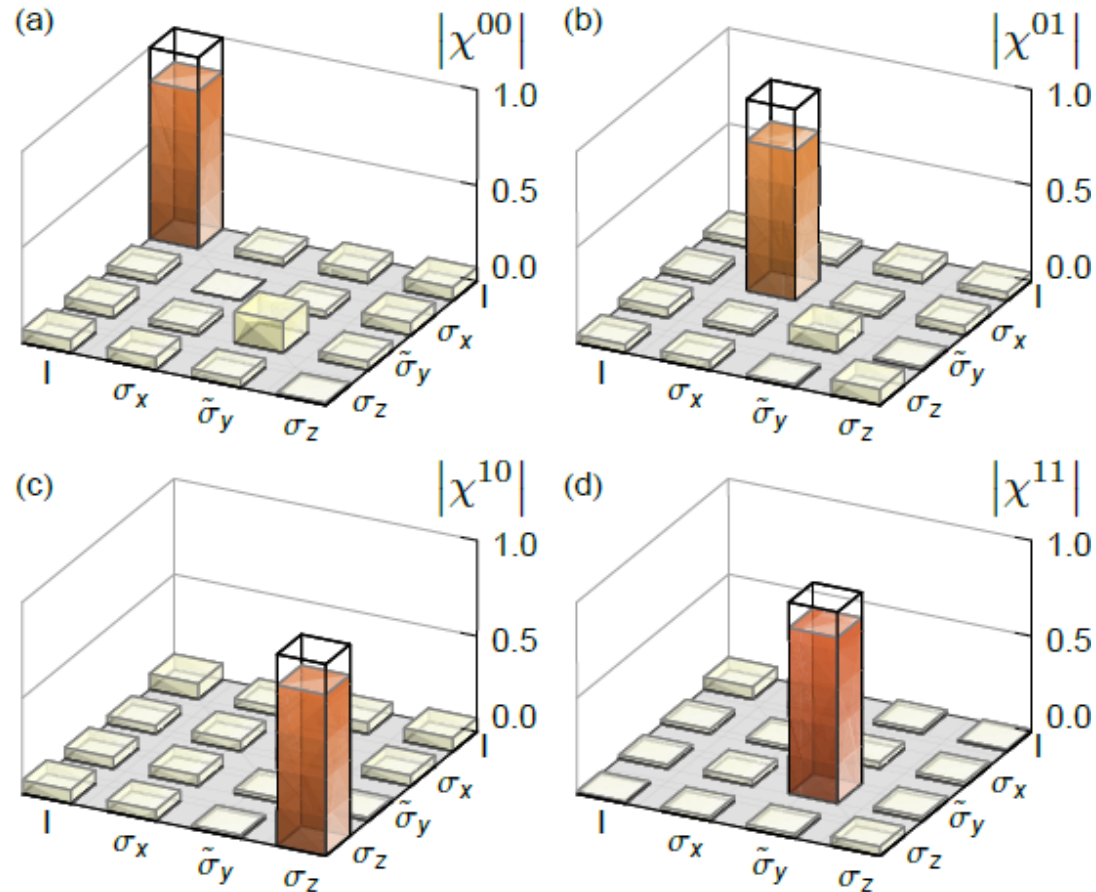
Benchmarking the Teleportation Process

procedure:

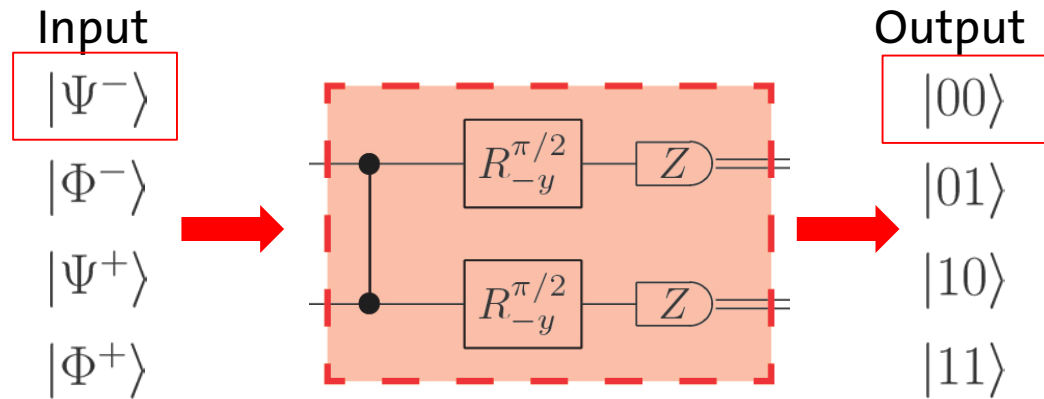
- prepare set of input states
- run circuit
- perform tomography on full system state
- project on basis states for qubits A and B
- analyze process acting on qubit C (see a-d)
- average fidelity of teleportation 83%

summary:

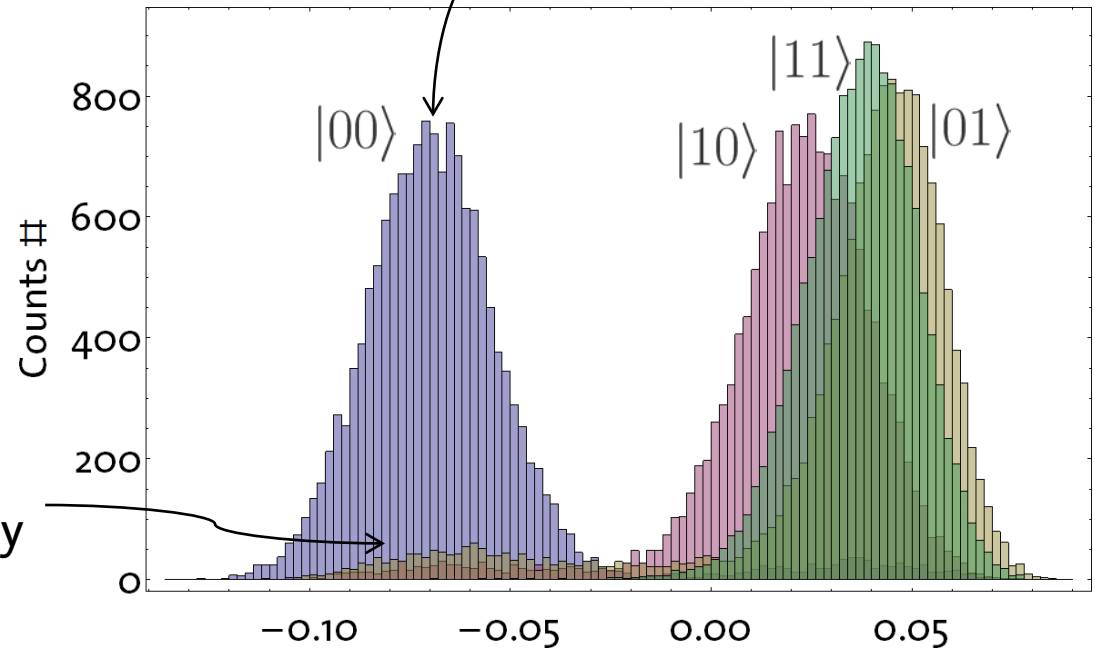
- verified operation of teleportation circuit & demonstrated 3-qubit quantum processor
- without single-shot projective read-out or feed-back



Post-Selected Teleportation: Bell Measurement



detection fidelity of 94 %



Limited by decay

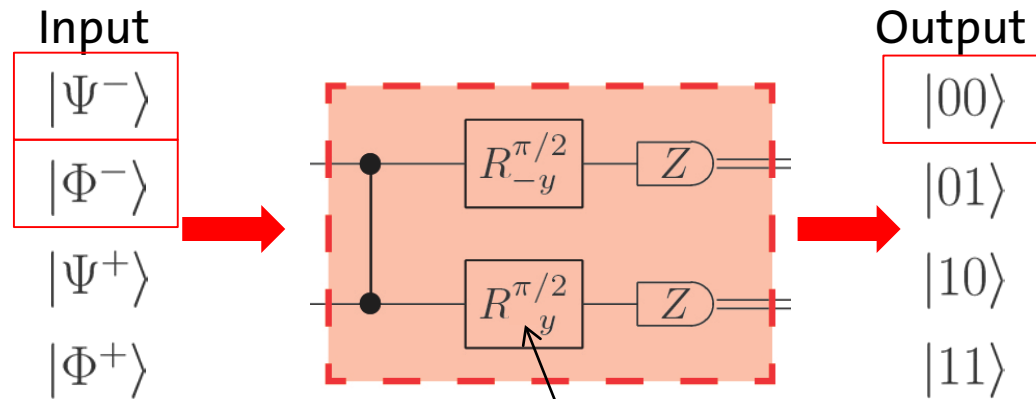
Operate parametric amplifier in phase sensitive mode

Maximize contrast of $|00\rangle$ to other states

Steffen *et al.*, *Nature* 500, 319 (2013)

ETH

Post-Selecting Every State Individually



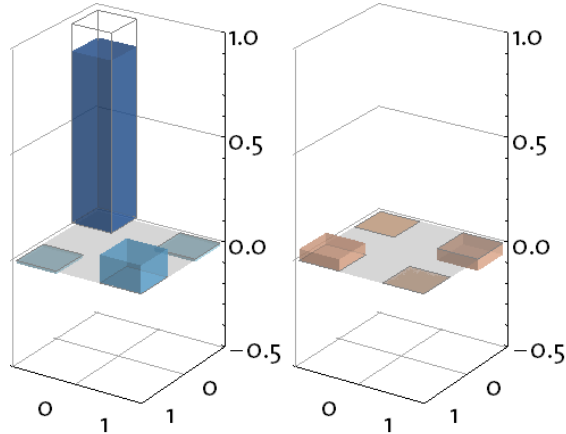
How to post-select on the other Bell states:

- Change the phases of the $\pi/2$ pulses
- Another Bell state is transformed to the $|00\rangle$ state
- Possibility to post-select on all four Bell states

Tomography of Teleported States with Post-Selection

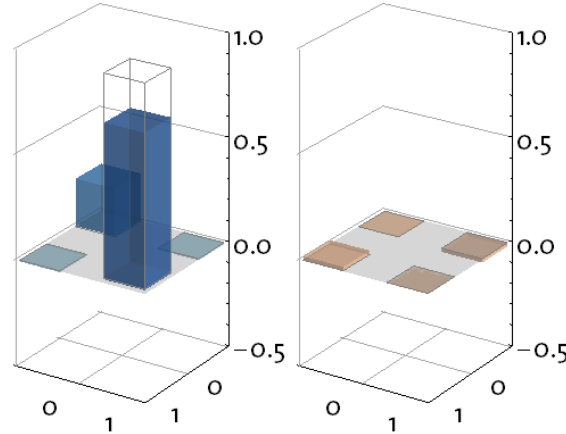
$$\psi_{in} = |0\rangle$$

82.2 %



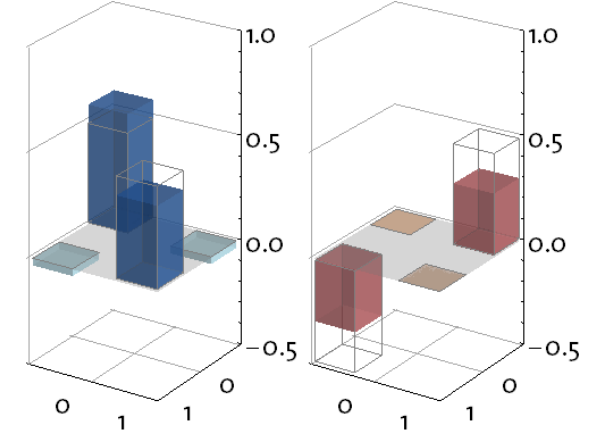
$$\psi_{in} = |1\rangle$$

80.5 %



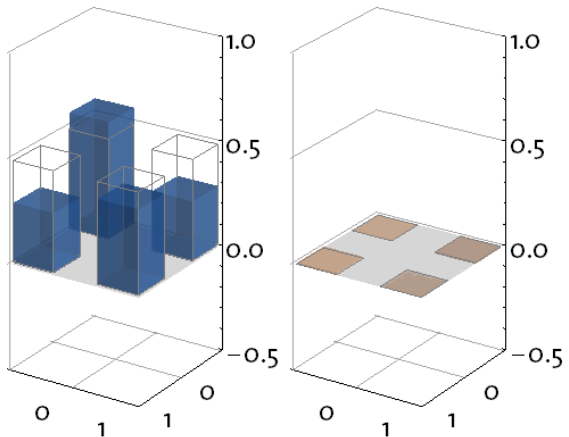
$$\psi_{in} = |0\rangle - i|1\rangle$$

79.4 %



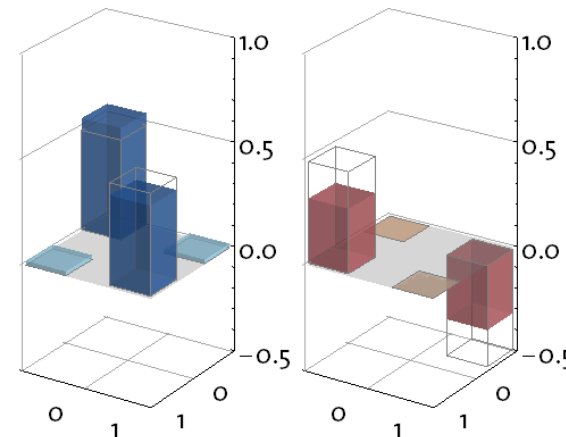
$$\psi_{in} = |0\rangle + |1\rangle$$

84.2 %



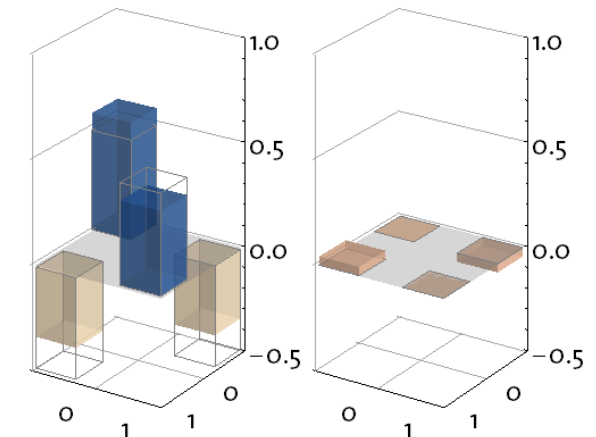
$$\psi_{in} = |0\rangle + i|1\rangle$$

79.5 %



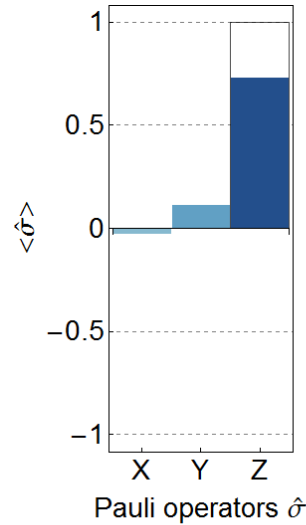
$$\psi_{in} = |0\rangle - |1\rangle$$

83.6 %

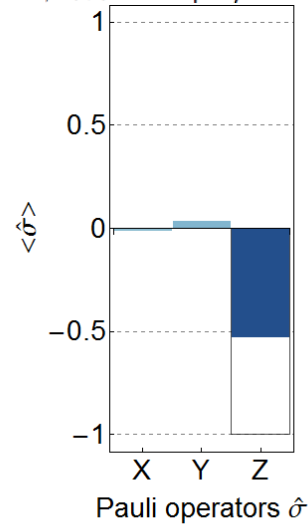


Post-Selection: Pauli-Sets of Teleported States

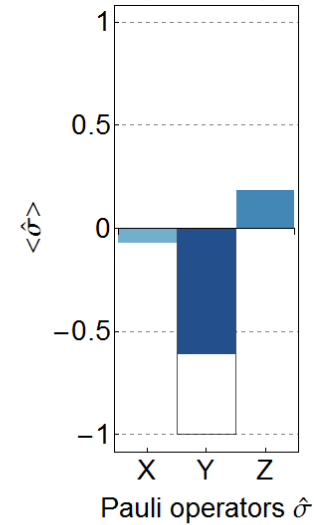
$$\psi_{in} = |0\rangle$$



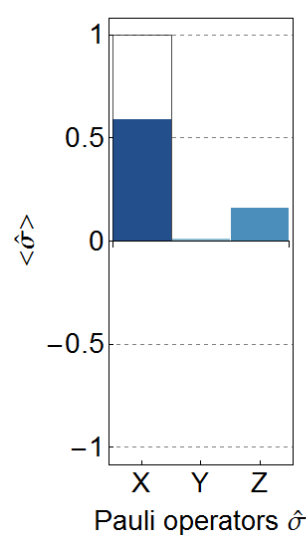
$$\psi_{in} = |1\rangle$$



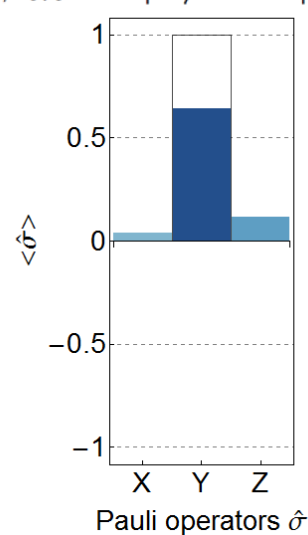
$$\psi_{in} = |0\rangle - i|1\rangle$$



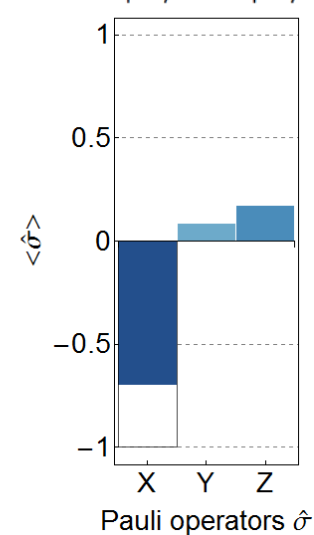
$$\psi_{in} = |0\rangle + |1\rangle$$



$$\psi_{in} = |0\rangle + i|1\rangle$$



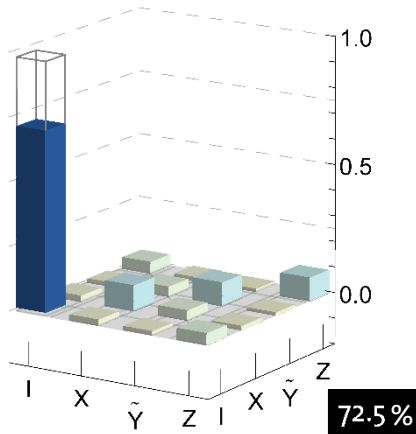
$$\psi_{in} = |0\rangle - |1\rangle$$



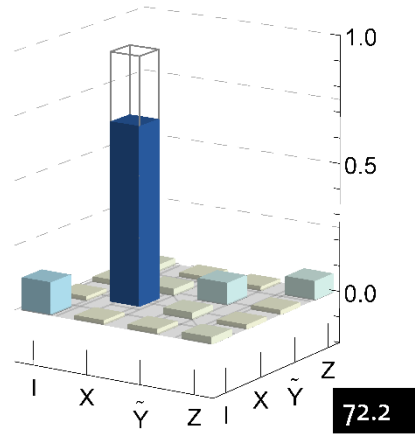
Process Tomography: Teleportation with Post-Selection

absolute value of process matrices $|\chi|$ for state transfer from qubit 1 to qubit 3:

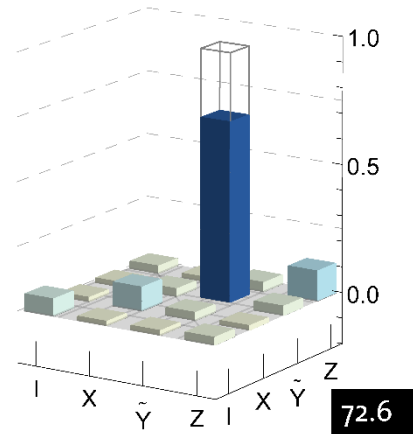
$|00\rangle \hat{=} |\Phi^-\rangle$



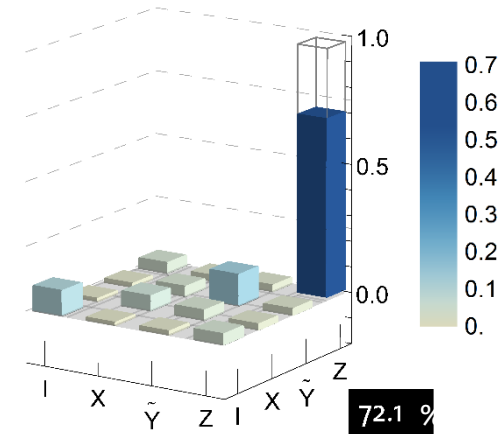
$|01\rangle \hat{=} |\Psi^-\rangle$



$|11\rangle \hat{=} |\Psi^+\rangle$



$|10\rangle \hat{=} |\Phi^+\rangle$



$$|\psi_{\text{out}}\rangle = |\psi_{\text{in}}\rangle$$

$$|\psi_{\text{out}}\rangle = X |\psi_{\text{in}}\rangle$$

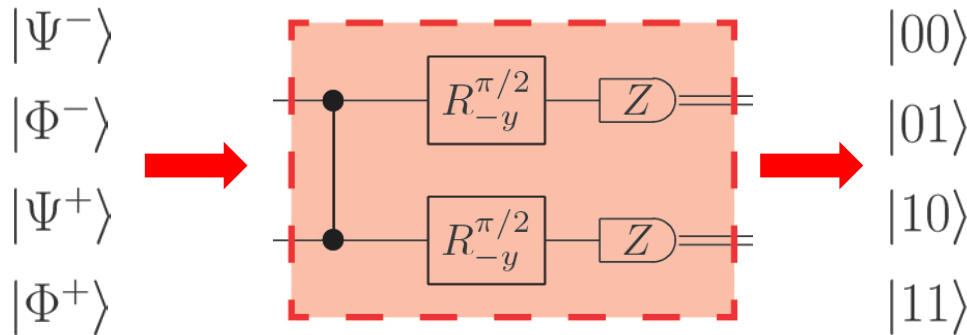
$$|\psi_{\text{out}}\rangle = \tilde{Y} |\psi_{\text{in}}\rangle$$

$$|\psi_{\text{out}}\rangle = Z |\psi_{\text{in}}\rangle$$

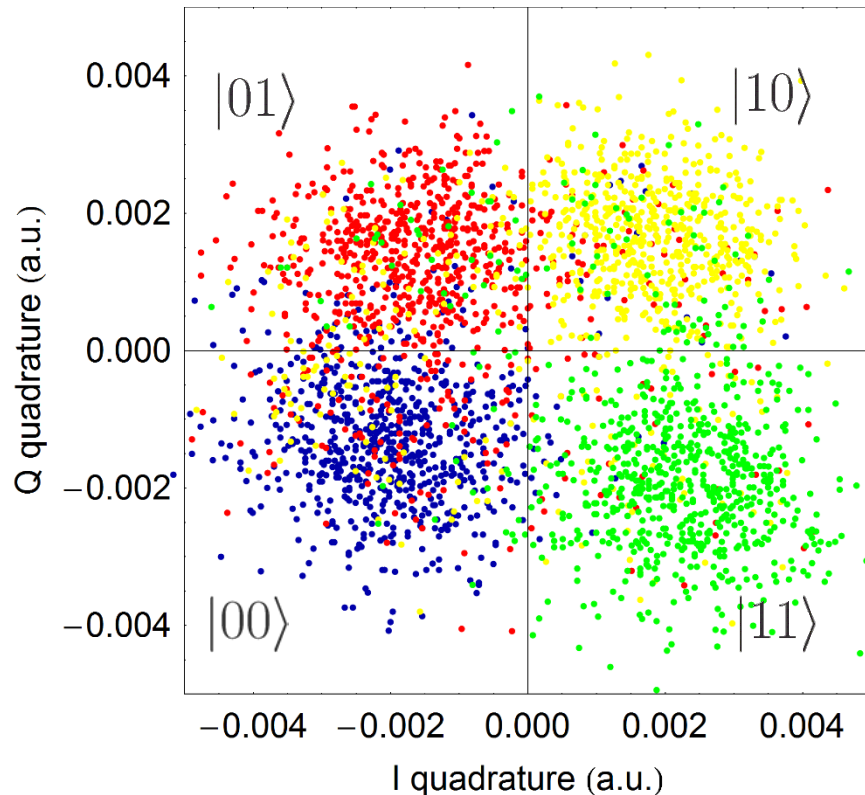
$$X = \hat{\sigma}_x, \tilde{Y} = i\hat{\sigma}_y, Z = \hat{\sigma}_z$$

Average process fidelity **72.3 ± 0.7 %**

Deterministic Bell-Measurement of all 4 States



- map Bell states on basis states
- perform joint two-qubit read-out
- paramp operated in the phase preserving mode to amplify both quadratures



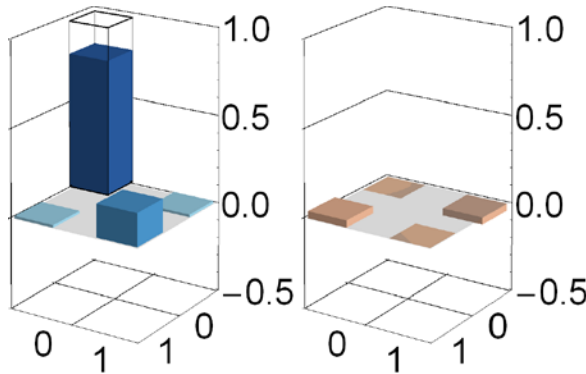
States are identified correctly with ~80% probability

		identified as			
		$ 00\rangle$	$ 01\rangle$	$ 10\rangle$	$ 11\rangle$
Prepared as	$ 00\rangle$	0.86	0.09	0.02	0.02
	$ 01\rangle$	0.14	0.73	0.04	0.09
	$ 10\rangle$	0.03	0.05	0.84	0.09
	$ 11\rangle$	0.08	0.10	0.09	0.73

Deterministic Teleportation – State Tomography

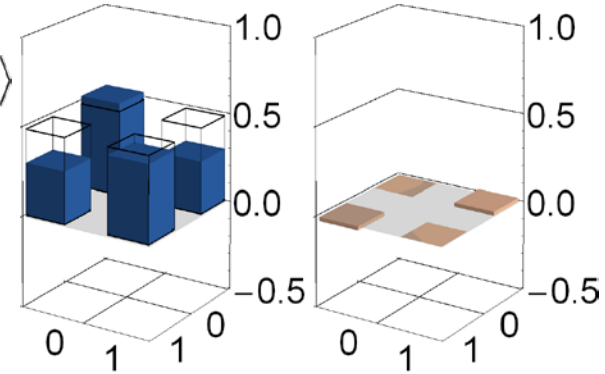
$$|\psi_{\text{in}}\rangle = |0\rangle$$

78.8 %



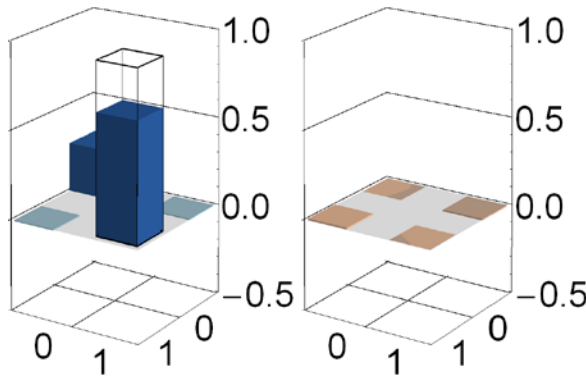
$$|\psi_{\text{in}}\rangle = |0\rangle + |1\rangle$$

82.0 %



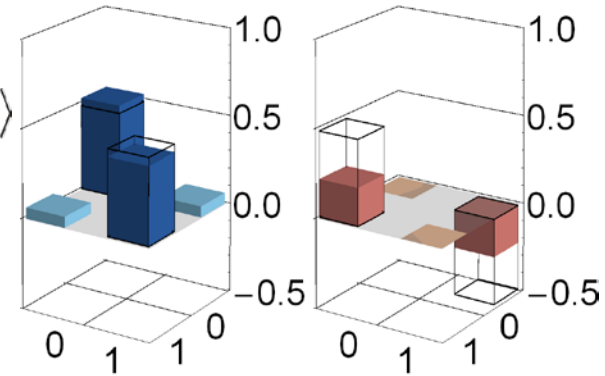
$$|\psi_{\text{in}}\rangle = |1\rangle$$

75.2 %



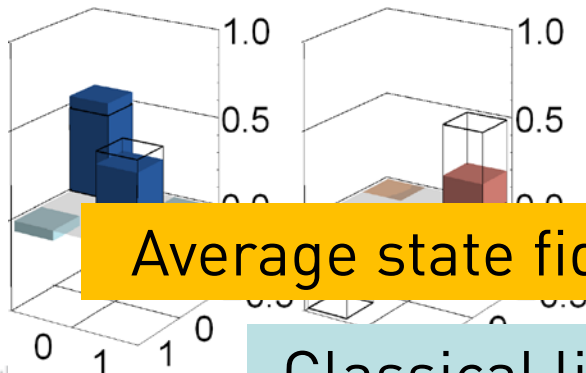
$$|\psi_{\text{in}}\rangle = |0\rangle + i|1\rangle$$

75.6 %



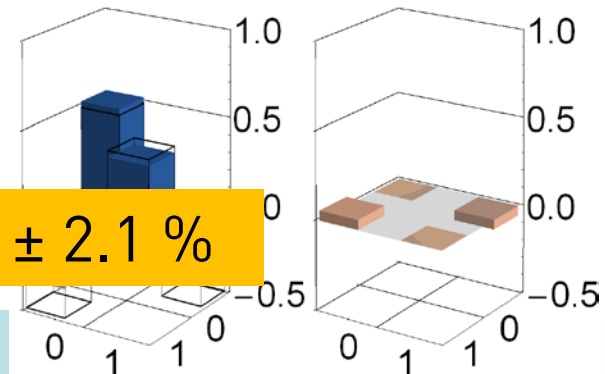
$$|\psi_{\text{in}}\rangle = |0\rangle - i|1\rangle$$

76.0 %



$$|\psi_{\text{in}}\rangle = |0\rangle - |1\rangle$$

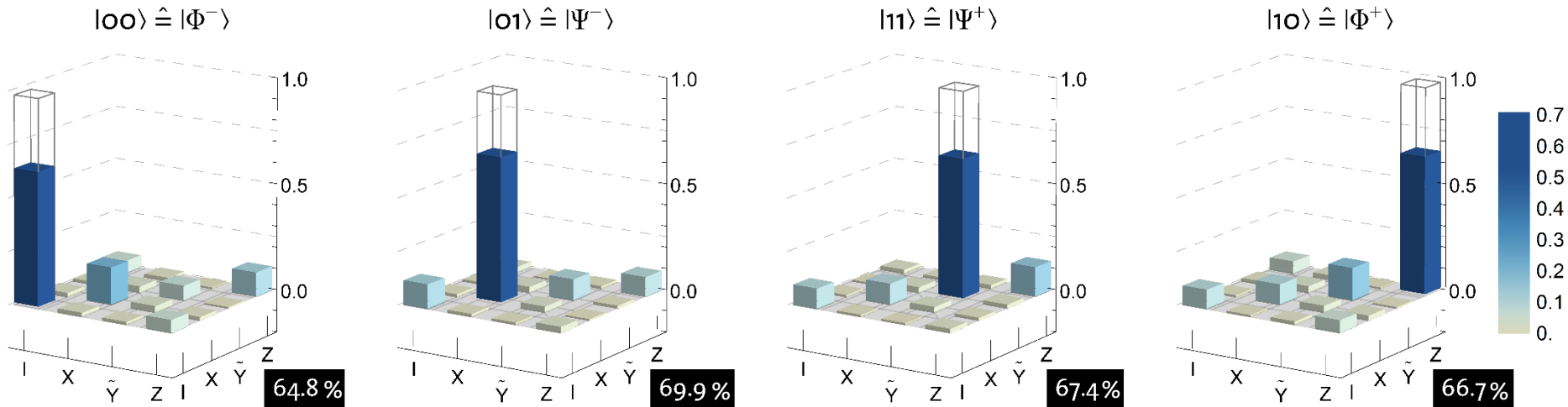
81.2 %



Average state fidelity of 78.1 ± 2.1 %

Classical limit: 66.7 %

Deterministic Teleportation – Process Tomography



$$|\psi_{\text{out}}\rangle = |\psi_{\text{in}}\rangle$$

$$|\psi_{\text{out}}\rangle = X |\psi_{\text{in}}\rangle$$

$$|\psi_{\text{out}}\rangle = \tilde{Y} |\psi_{\text{in}}\rangle$$

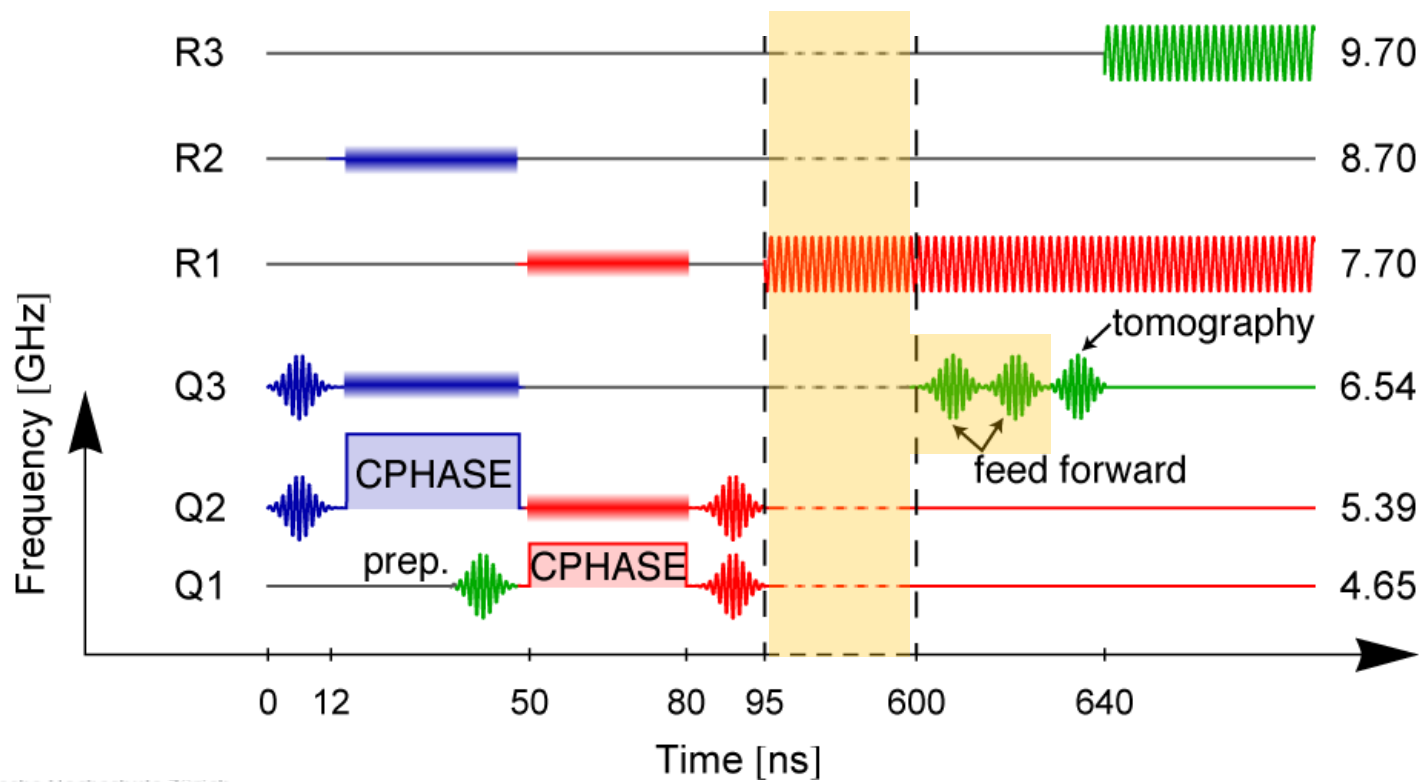
$$|\psi_{\text{out}}\rangle = Z |\psi_{\text{in}}\rangle$$

Average process fidelity **67.1 ± 0.5 %**

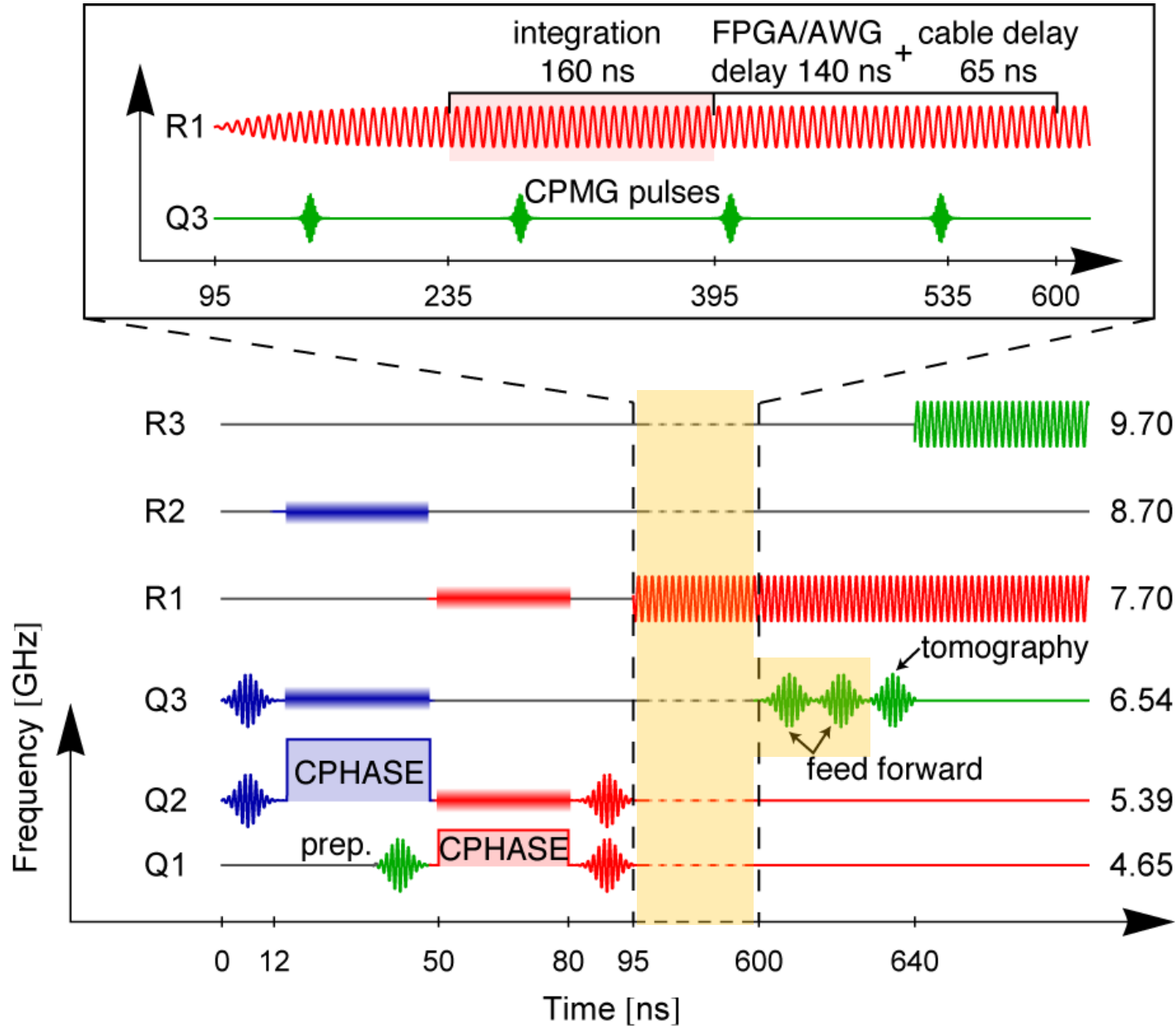
Average state fidelity **78.1 ± 0.9 %**

$$\mathcal{F}_p = (\mathcal{F}_s(d + 1) - 1)/d$$

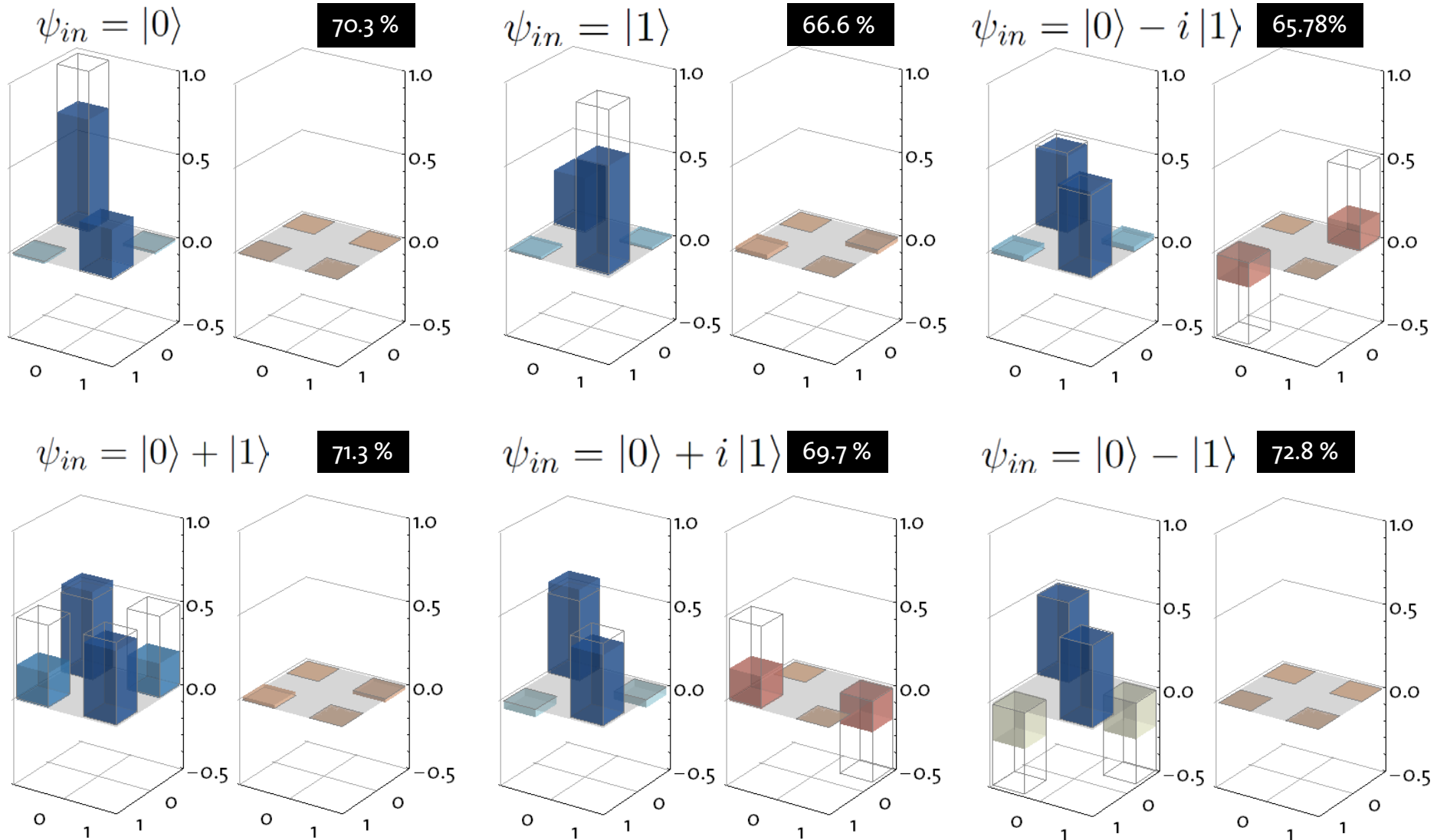
Pulse Scheme



Pulse Scheme



Tomography of Teleported States with Feed-Forward



Average state fidelity of **69.5±0.1 %**

classical limit: 66.7 %

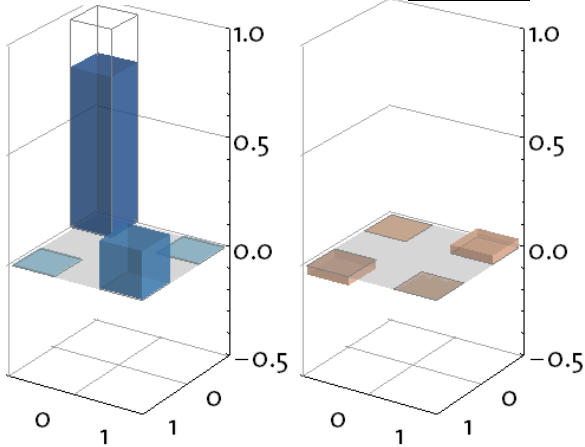
Steffen *et al.*, *Nature* 500, 319 (2013)

Tomography of Teleported States with Feed-Forward

averaged readout of qubit 3

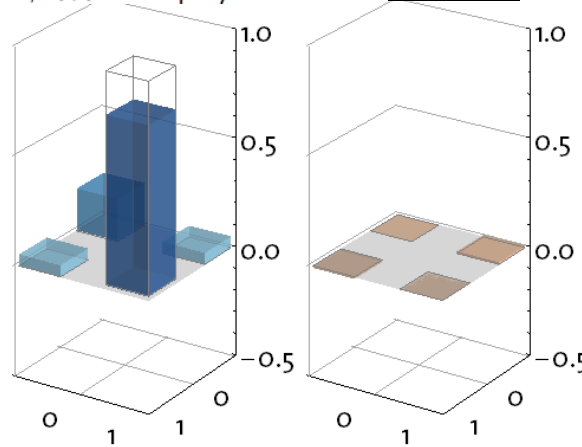
$$\psi_{in} = |0\rangle$$

77.5 %



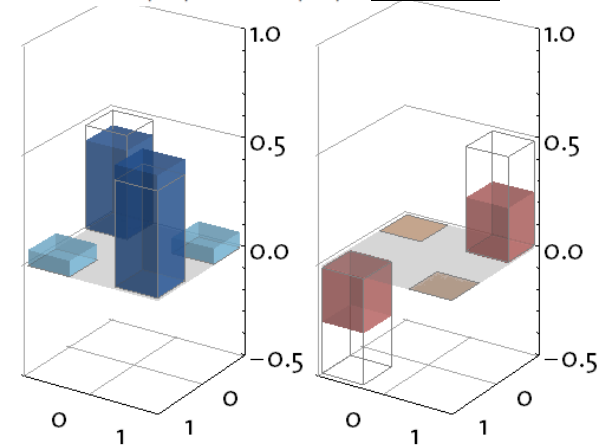
$$\psi_{in} = |1\rangle$$

79.9 %



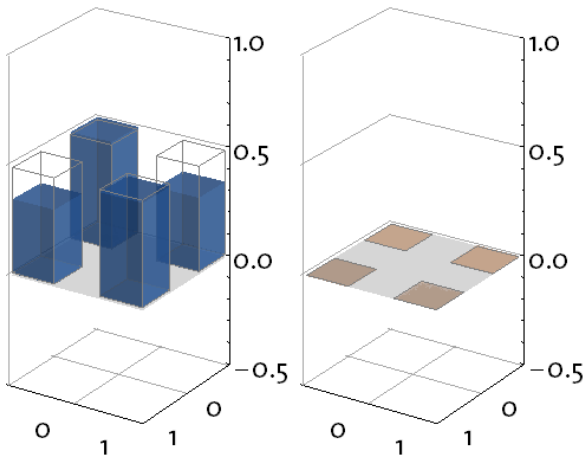
$$\psi_{in} = |0\rangle - i|1\rangle$$

76.2 %



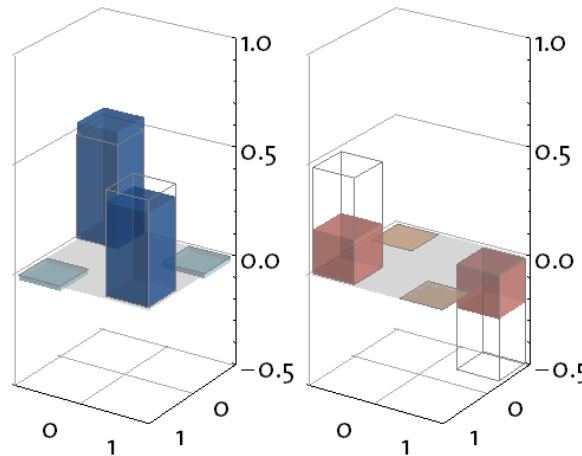
$$\psi_{in} = |0\rangle + |1\rangle$$

85.3 %



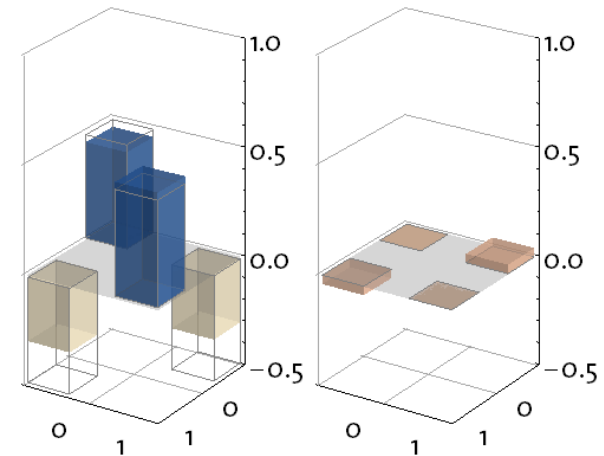
$$\psi_{in} = |0\rangle + i|1\rangle$$

71.2 %



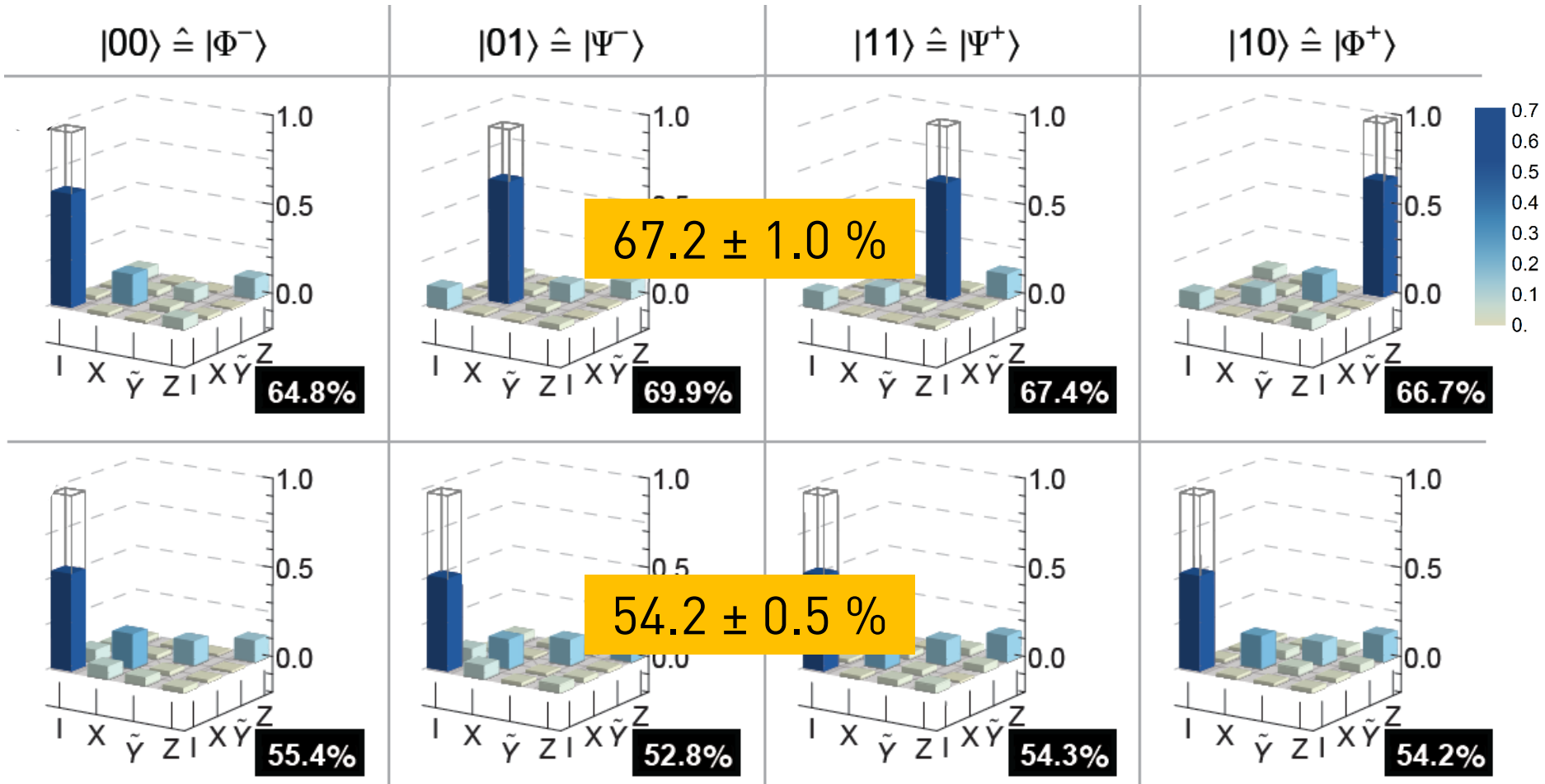
$$\psi_{in} = |0\rangle - |1\rangle$$

80.7 %



Average state fidelity of **78.5 ± 0.9%**

Process Tomography – w/o and with Feed-Forward



$$|\psi_{out}\rangle = |\psi_{in}\rangle$$

$$|\psi_{out}\rangle = |\psi_{in}\rangle$$

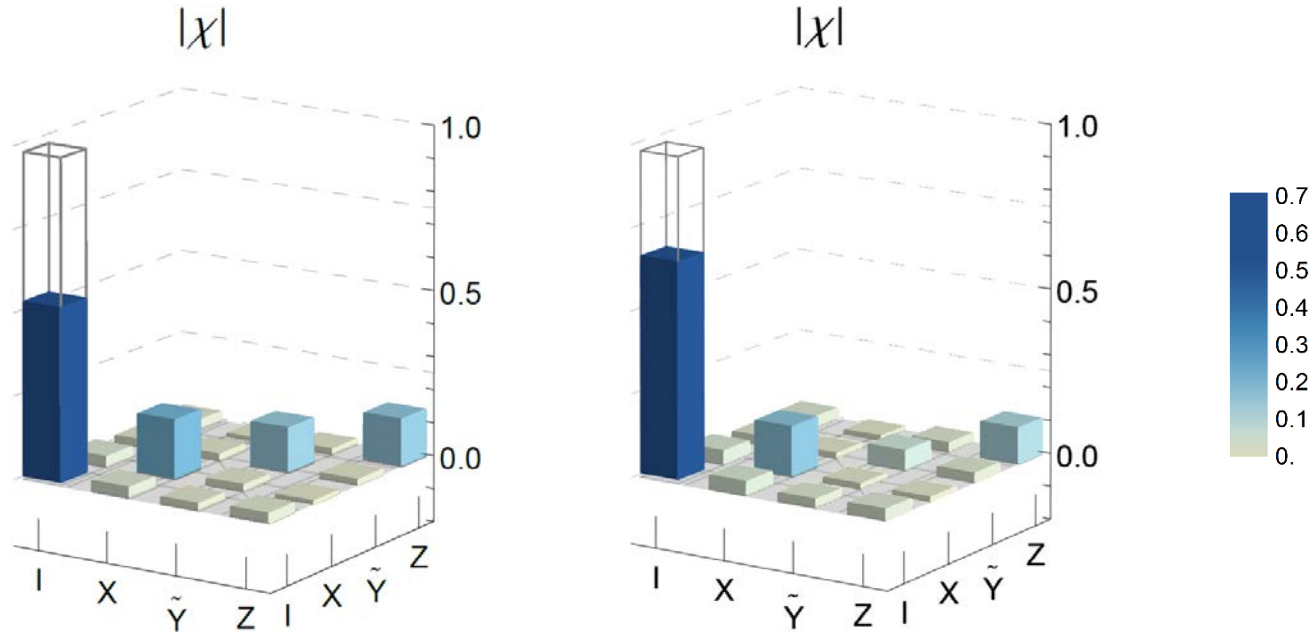
$$|\psi_{out}\rangle = |\psi_{in}\rangle$$

$$|\psi_{out}\rangle = |\psi_{in}\rangle$$

Classical limit: 50 %

$$\mathcal{F}_p = (\mathcal{F}_s(d+1) - 1)/d$$

Teleportation Process with Feed-Forward



Average process fidelity with single shot readout: **$54.2 \pm 0.1 \%$**

Average process fidelity with averaged readout: **$67.7 \pm 1.1 \%$**

Teleportation State and Process Fidelities Summary

	States:	Processes:
Thresholds	66.7 %	50 %
Post-selection	81.6 ± 3.0 %	72.3 ± 1.4 %
Deterministic	78.1 ± 2.1 %	67.2 ± 1.0 %
Feed-forward	69.5 ± 0.5 %	54.2 ± 0.5 %
Averaged	78.5 ± 1.0 %	67.7 ± 0.6 %

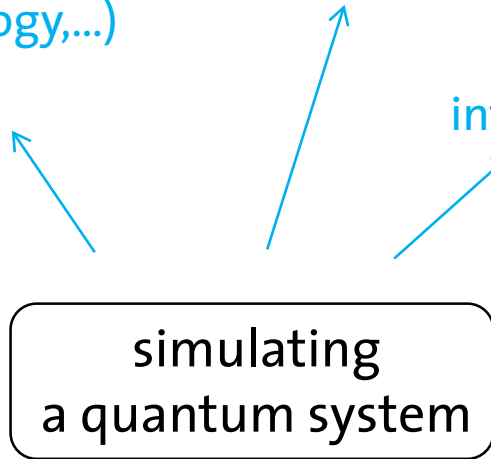
- First experimental implementation of full quantum teleportation in a solid state system
- High qubit teleportation rate (~10.000/s)

Quantum Simulation

describes physical system of interest (physics, chemistry, biology,...)

encodes hard classical problem

interesting toy model



... difficult on classical computer!

universal quantum computer

... still to be realized!

... OR ...

quantum simulator

... sufficient controllability, flexibility!

use

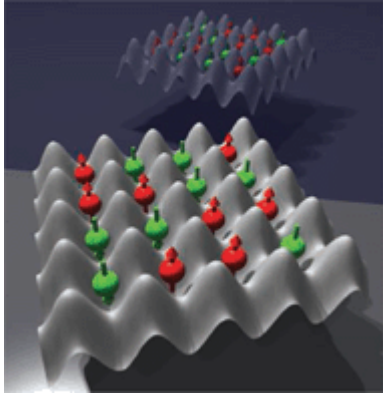
map!

Feynman, *Int. Journal of Th. Phys.* 21, 467 (1982)

Lloyd, *Science* 273, 5278 (1996)

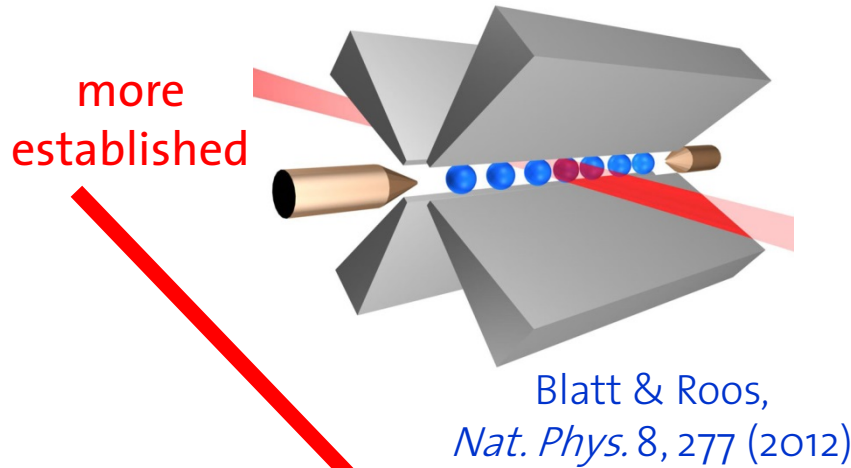
Systems for Quantum Simulation

Ultracold gases

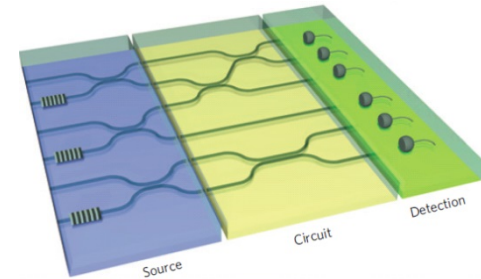


Bloch et al.,
Nat. Phys. 8, 267 (2012)

Trapped ions

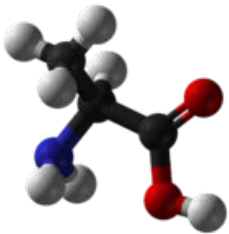


Optical photons



Aspuru-Guzik & Walther,
Nat. Phys. 8, 258 (2012)

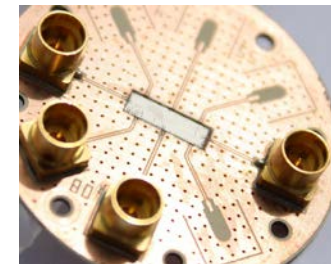
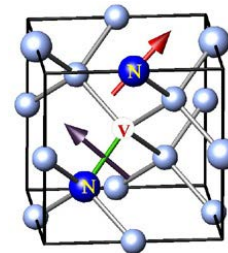
Nuclear magnetic resonance



Vandersypen & Chuang,
RMP 76, 1037 (2004)

under
development

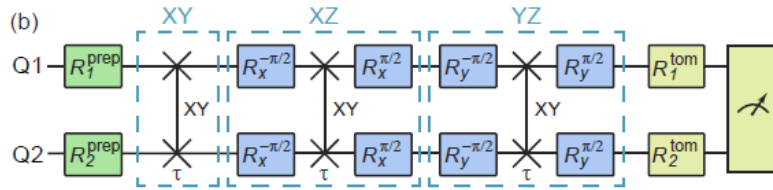
Solid state quantum
devices



Georgescu et al., *RMP* 86, 153 (2014)

Quantum Simulation with Superconducting Circuits

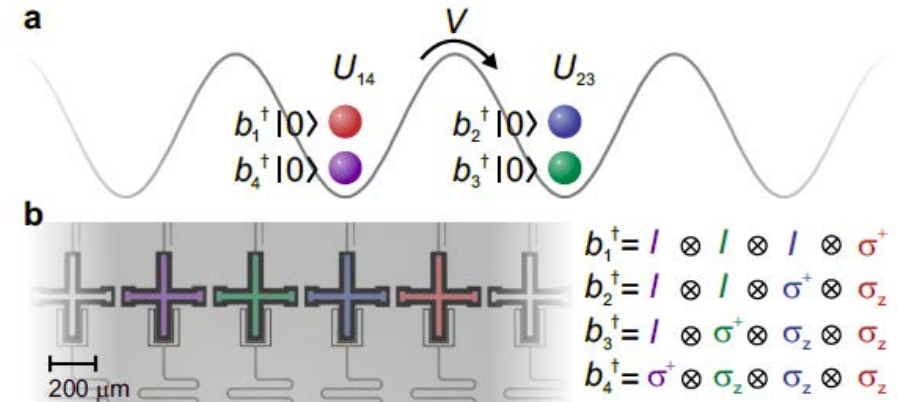
Digital simulation of exchange, Heisenberg, Ising spin models



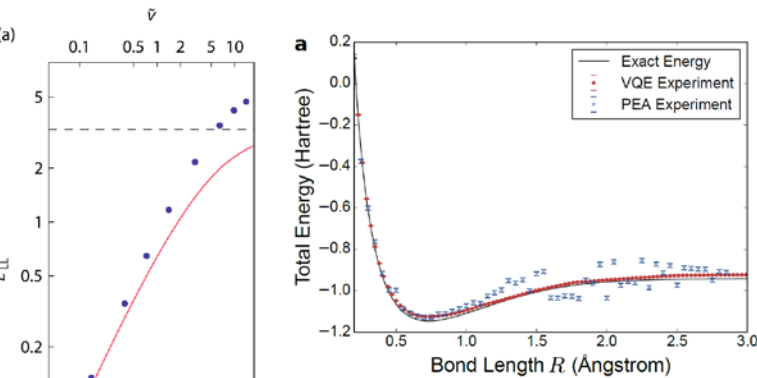
Salathe *et al.*, *PRX* 5, 021027 (2015)

Quantum simulation of correlated systems with variational Ansatz based on MPS

... two-mode fermionic Hubbard models



Barends *et al.*, *Nat. Com.* 6, 7654 (2015)

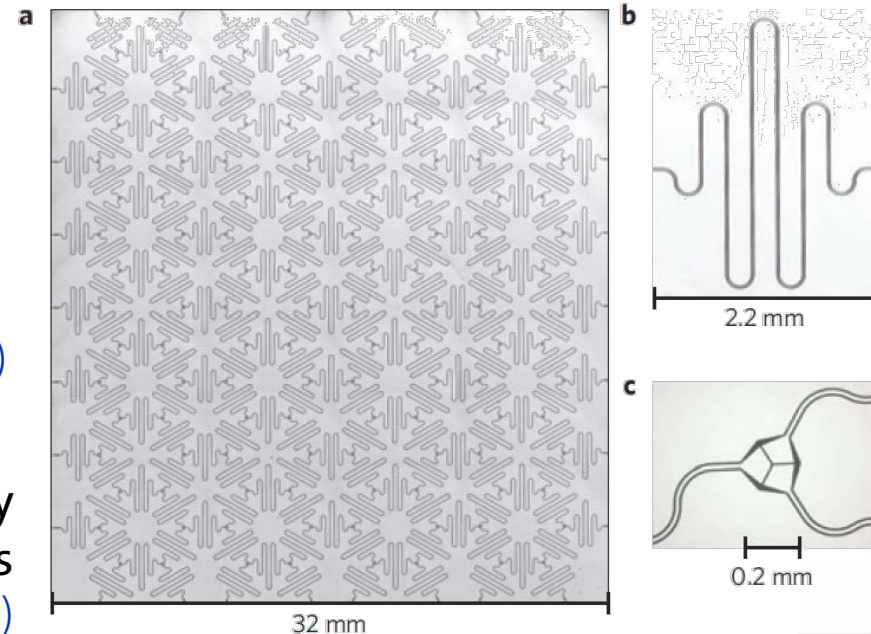


Eichler *et al.*, *Phys. Rev. X* 5, 041044 (2015)

O'Malley *et al.*, *arXiv:1512.06860* (2015)

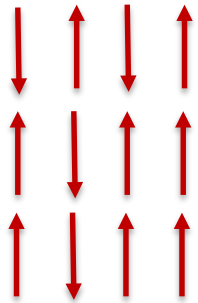
Analog simulations with cavity and/or qubit arrays

Houck *et al.*, *Nat Phys.* 8, 292 (2012)

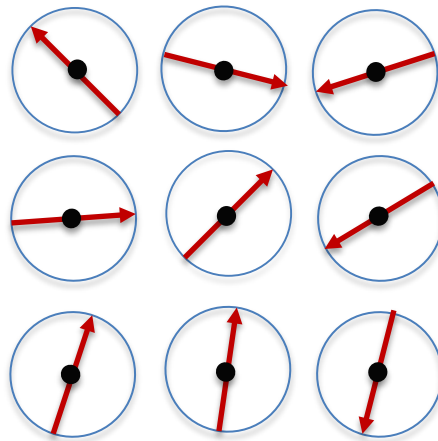


Digital Quantum Simulation of Spin Models

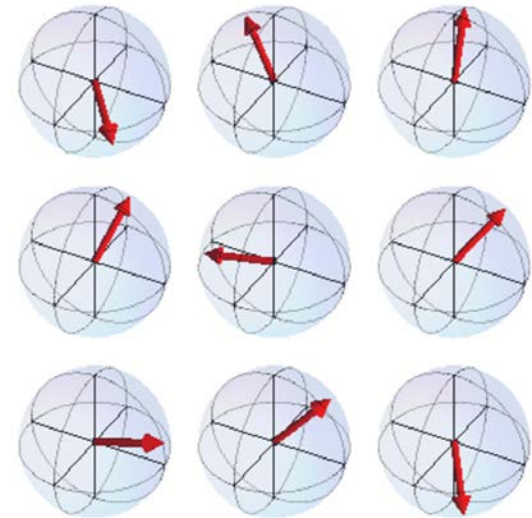
Ising (up-down)



XY (Compass)



XYZ (Heisenberg)



interaction term

$$J \sum_{(i,j)} \sigma_i^x \sigma_j^x$$

$$\sum_{(i,j)} J_x \sigma_i^x \sigma_j^x + J_y \sigma_i^y \sigma_j^y$$

$$\sum_{(i,j)} J_x \sigma_i^x \sigma_j^x + J_y \sigma_i^y \sigma_j^y + J_z \sigma_i^z \sigma_j^z$$

+ magnetic field term, e.g. $B \sum_{i=1}^n \sigma_i^z$

Potential applications in physics, chemistry and material science.

Number of spins > 100  classical simulations intractable

Digital Quantum Simulation

model Hamiltonian
to be realized in 'software'

$$H = \sum_{k=1}^N H_k$$

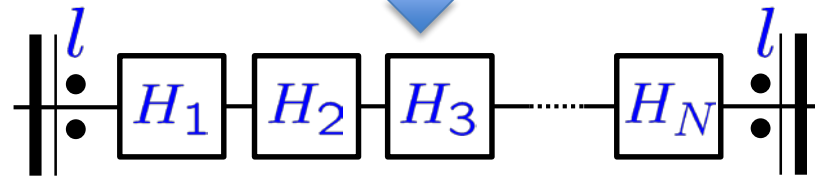
realizable in 'hardware'

time evolution:

Suzuki-Lie-Trotter approximation

$$\exp(-iHt) = \left(\prod_{k=1}^N \exp(-iH_k t/l) \right)^l + E(t, l)$$

higher-order terms



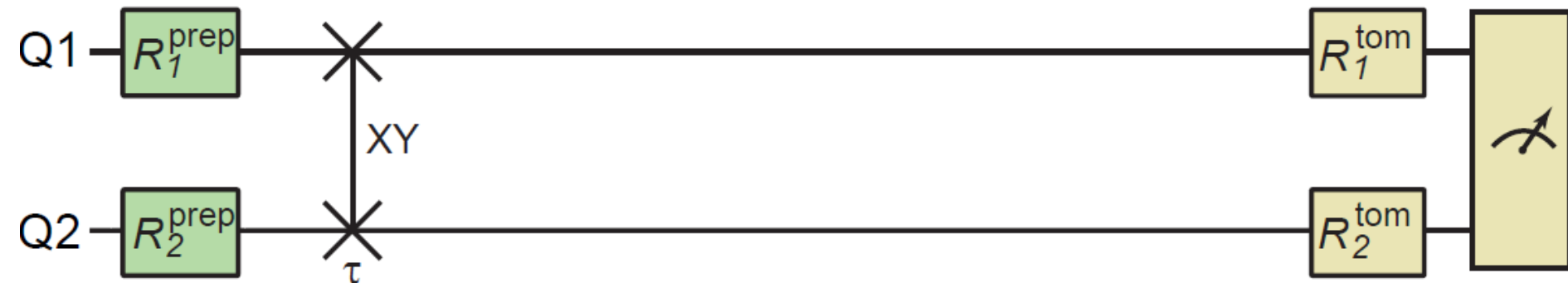
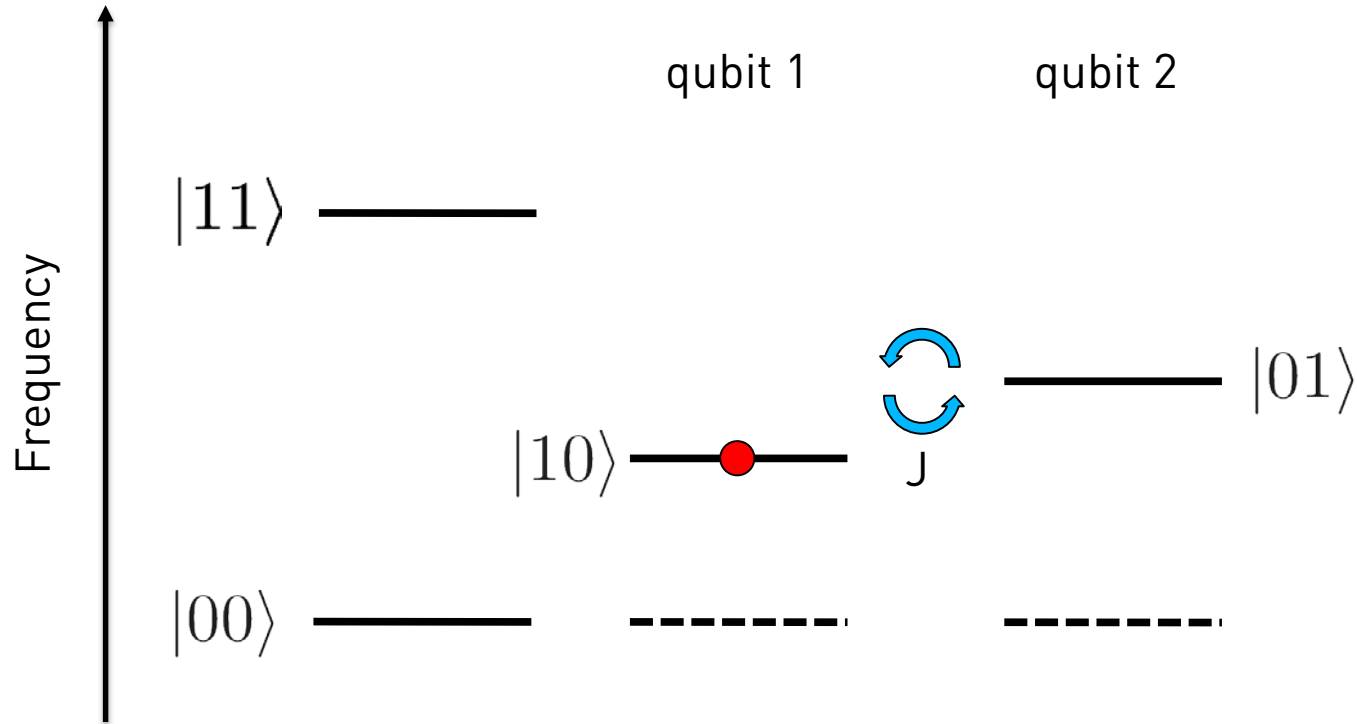
"digital quantum simulation"

$$E(t, l) \rightarrow 0 \text{ if } l \rightarrow \infty$$

works for all Hamiltonians with local interactions (universal)

XY Interaction (Hardware) Controlled by Detuning

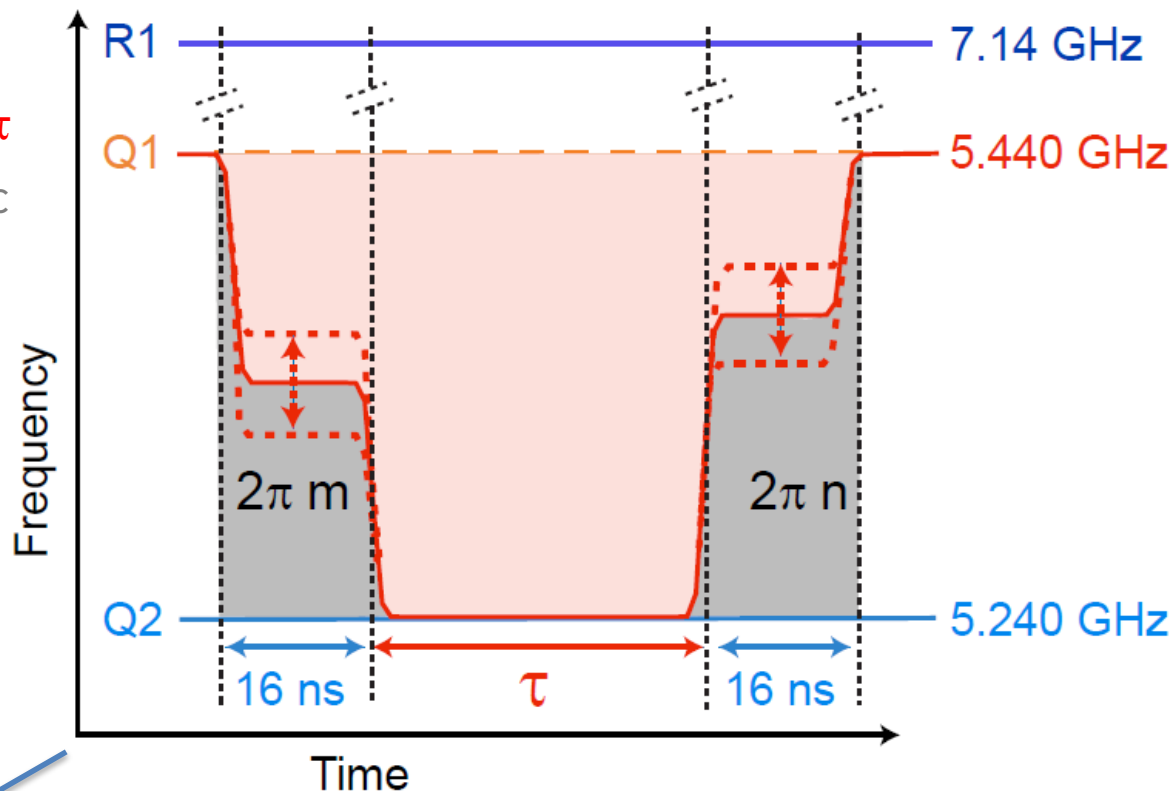
$$H_{2Q} = -\frac{1}{2}\omega_1\hat{\sigma}_1^z - \frac{1}{2}\omega_2\hat{\sigma}_2^z + J(\hat{\sigma}_1^+\hat{\sigma}_2^- + \hat{\sigma}_1^-\hat{\sigma}_2^+)$$



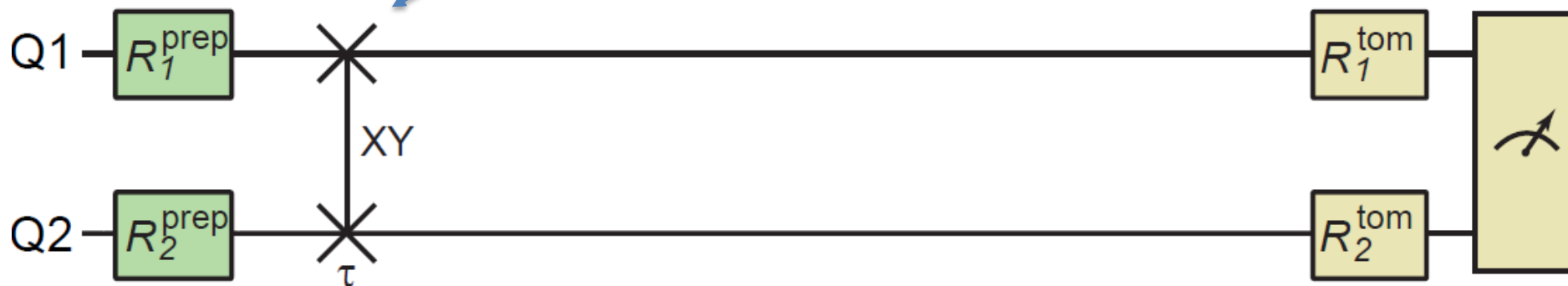
Frequency Control through Flux Bias

Profile of flux pulse:

- tunable **interaction time τ**
- compensation of dynamic phase

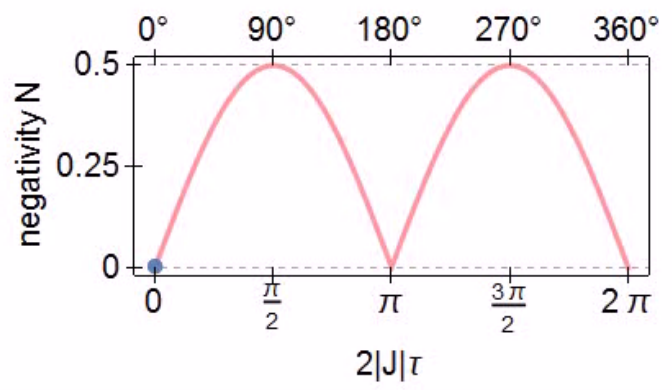
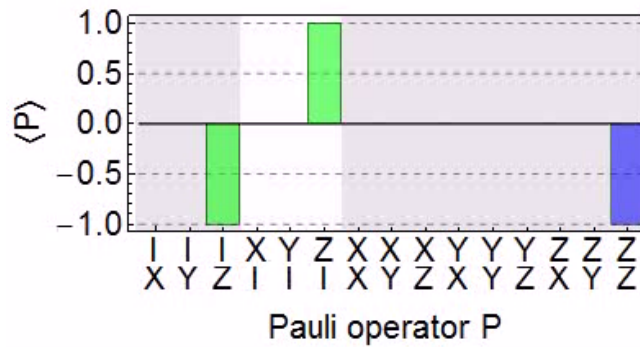
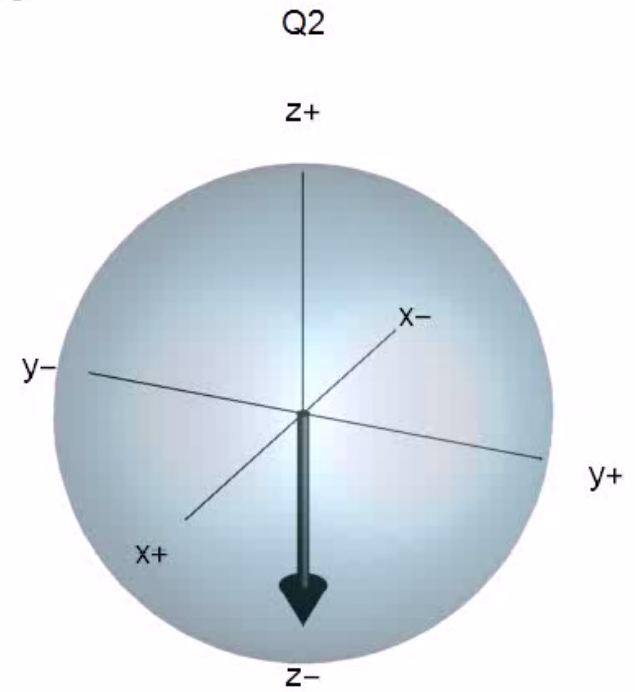
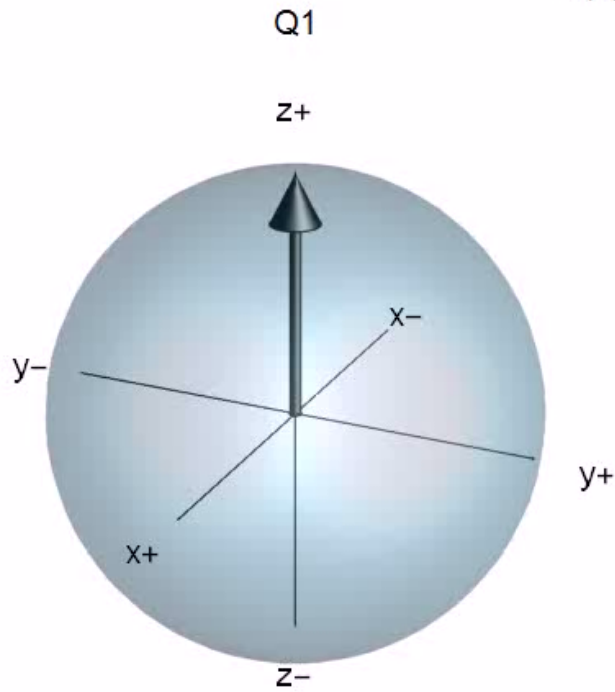


Circuit representation of gate sequence:



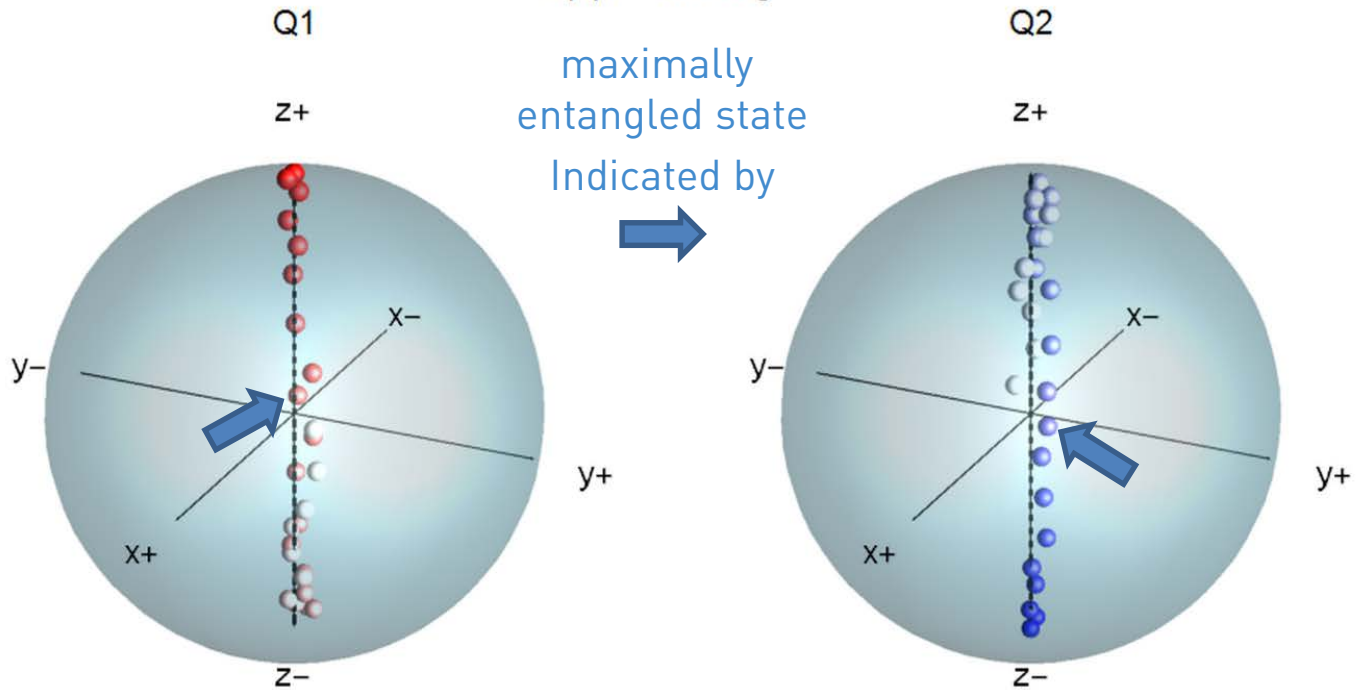
XY Interaction: Calculation

initial state: $|0\rangle \otimes |1\rangle$
 $2|J|\tau = 0 \text{ deg}$

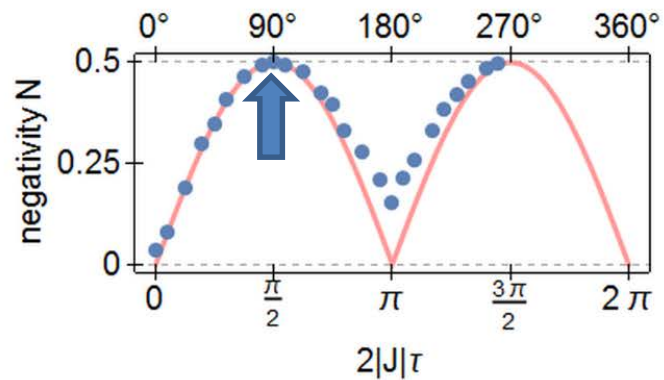
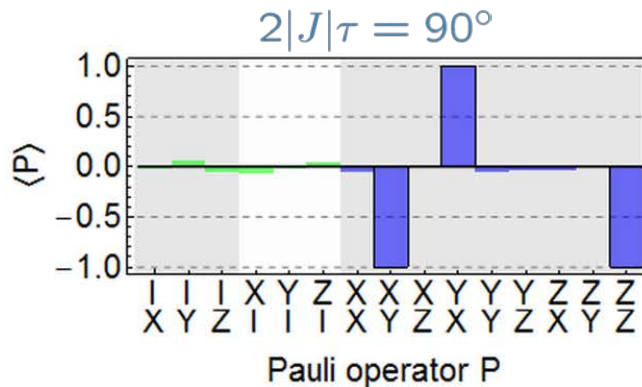


XY Interaction: Experimental Data

initial state: $|0\rangle \otimes |1\rangle$
 $2|J|\tau = 270 \text{ deg}$



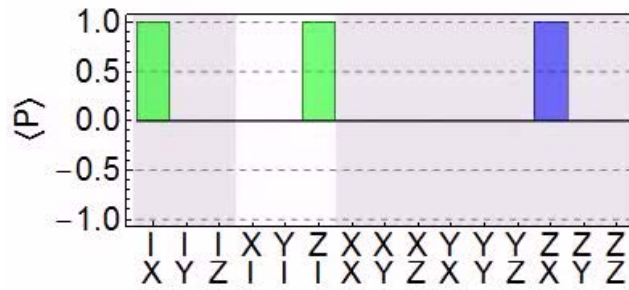
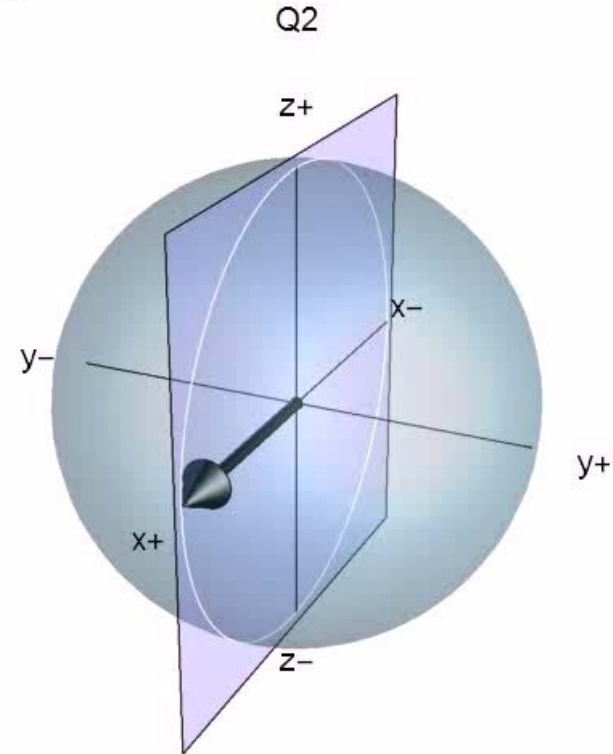
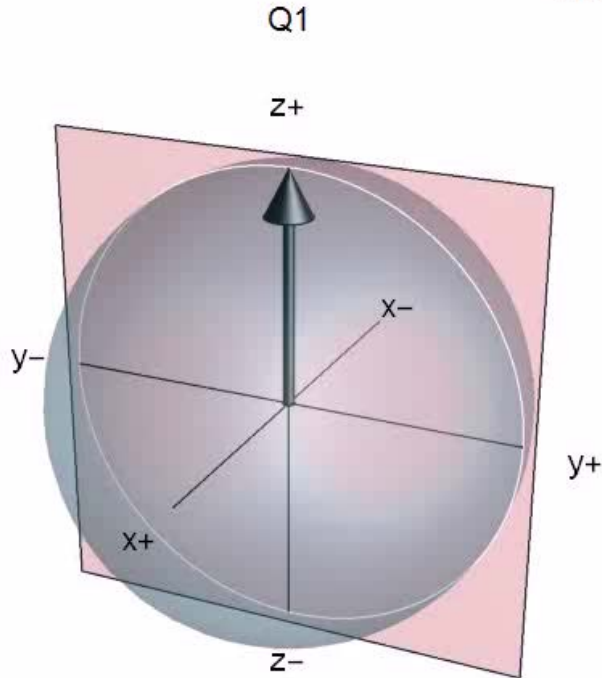
state fidelity: 99.7 %



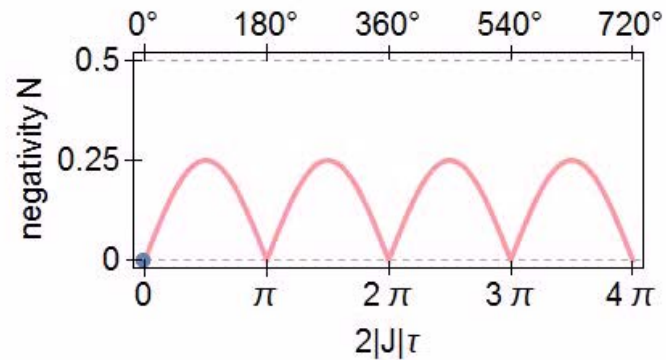
XY Interaction: Calculation

initial state: $\frac{1}{\sqrt{2}}|0\rangle \otimes (|0\rangle + |1\rangle)$

$2|J|\tau = 0 \text{ deg}$

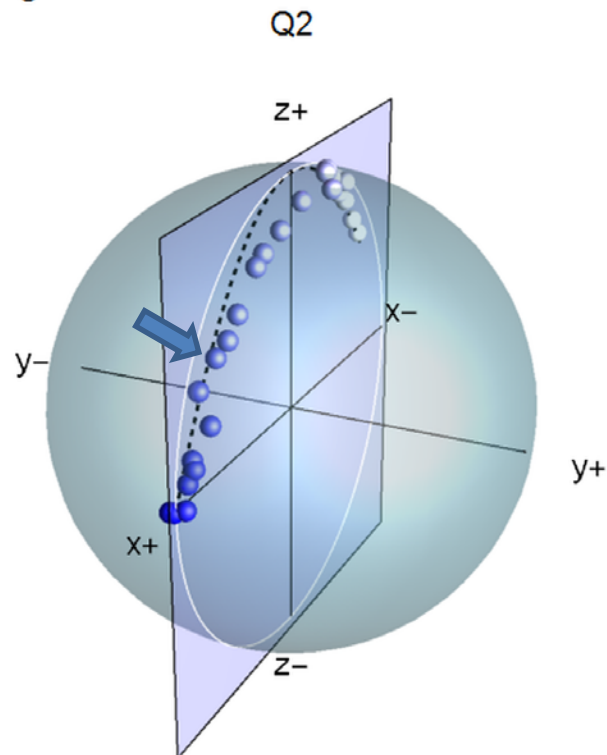
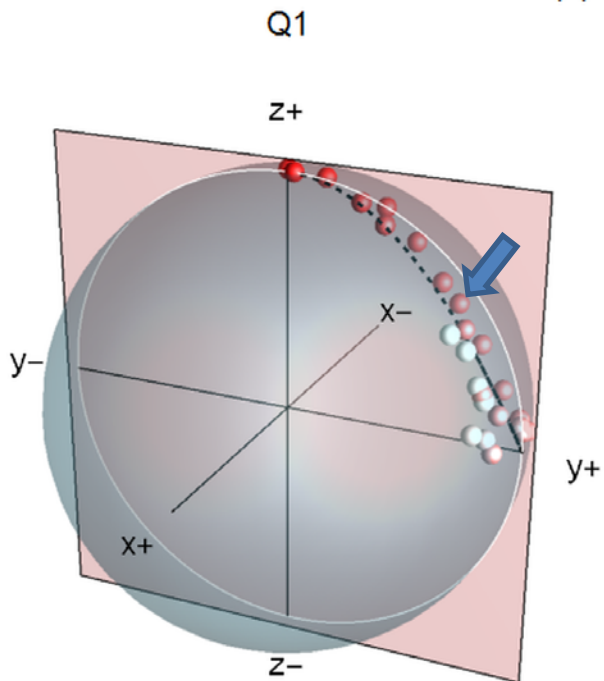


Pauli operator P

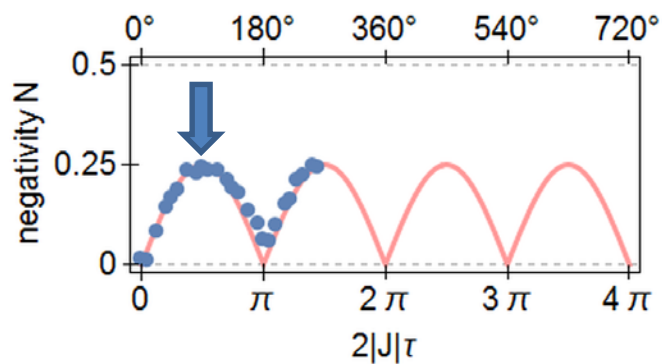
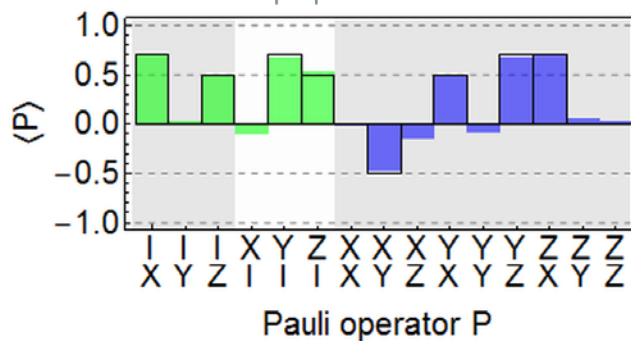


XY Interaction: Experimental Data

initial state: $\frac{1}{\sqrt{2}}|0\rangle\otimes(|0\rangle+|1\rangle)$
 $2|J|\tau = 270 \text{ deg}$



state fidelity: $F = 99.4 \%$
 $2|J|\tau = 90^\circ$



$$F \equiv \left(\text{Tr} \sqrt{\sqrt{\rho_{\text{theo}}} \rho_{\text{exp}} \sqrt{\rho_{\text{theo}}}} \right)^2$$

Digital Simulation of Heisenberg XYZ Interaction

Hamiltonian to be simulated: $H_{xyz} = J(\sigma_1^x \sigma_2^x + \sigma_1^y \sigma_2^y + \sigma_1^z \sigma_2^z)$

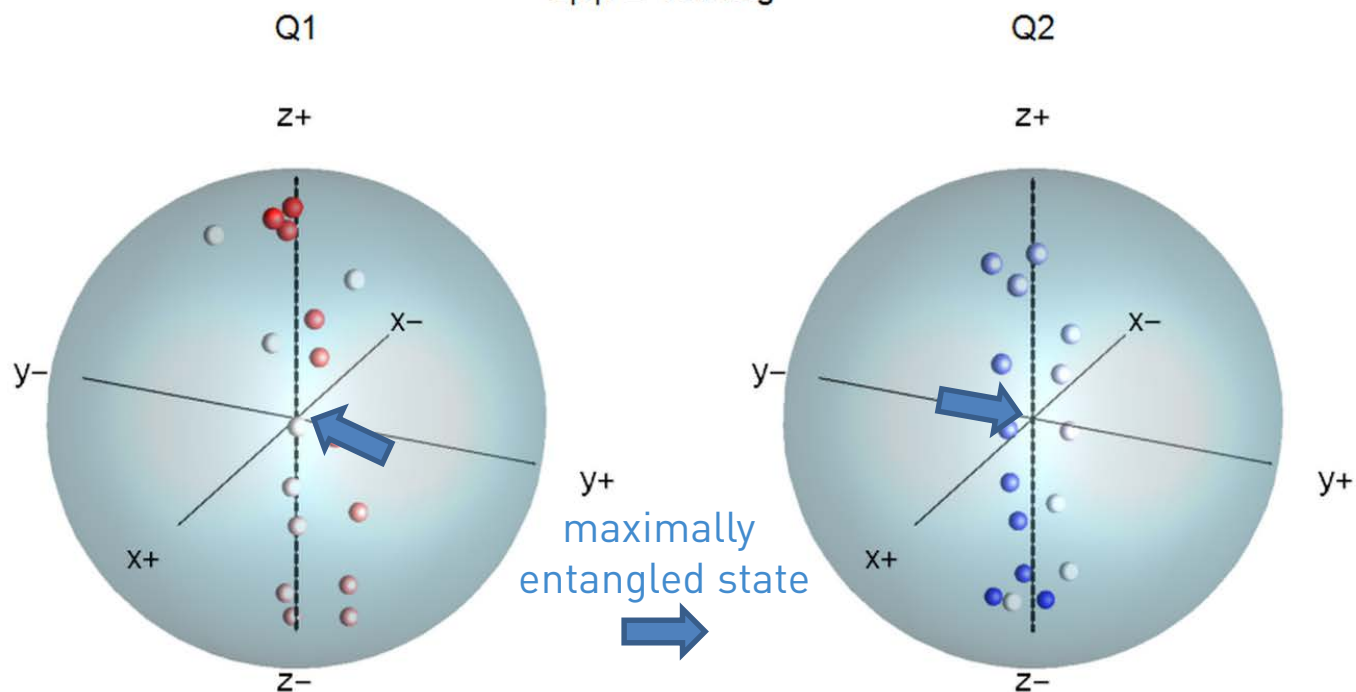
operators commute

-> no Trotter decomposition required, exact result in one step

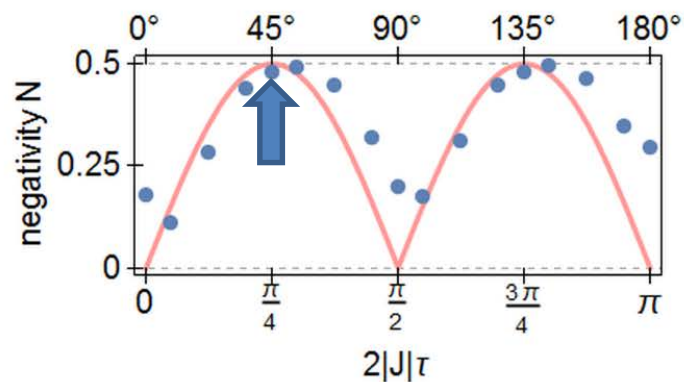
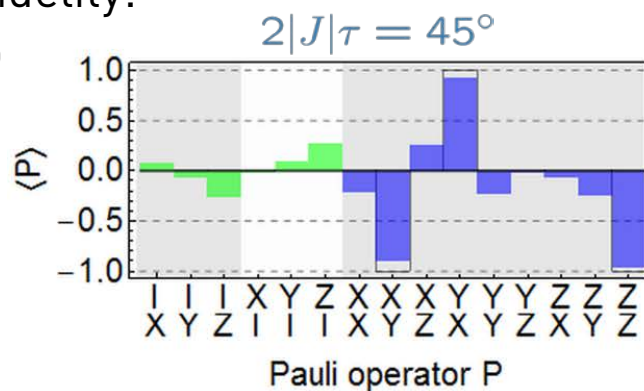
$$e^{-iH\tau} = e^{-iH_{xz}\tau}$$

Heisenberg XYZ Interaction: Experimental Data

initial state: $|0\rangle \otimes |1\rangle$
 $2|J|\tau = 180 \text{ deg}$



state fidelity:
 94.7 %



Heisenberg XYZ Interaction: Calculation

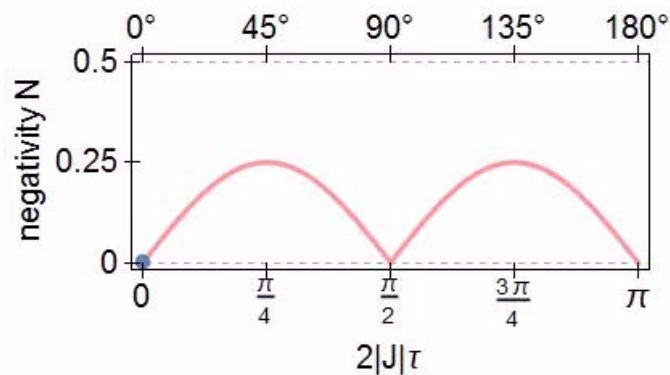
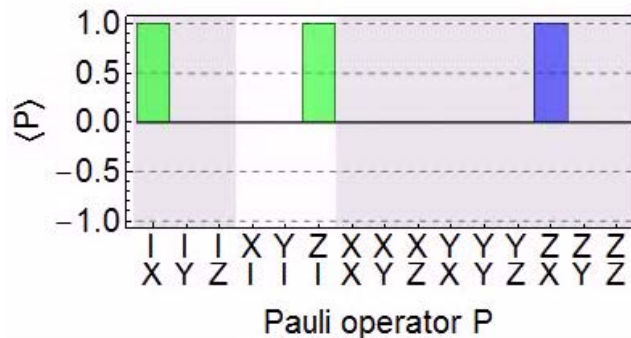
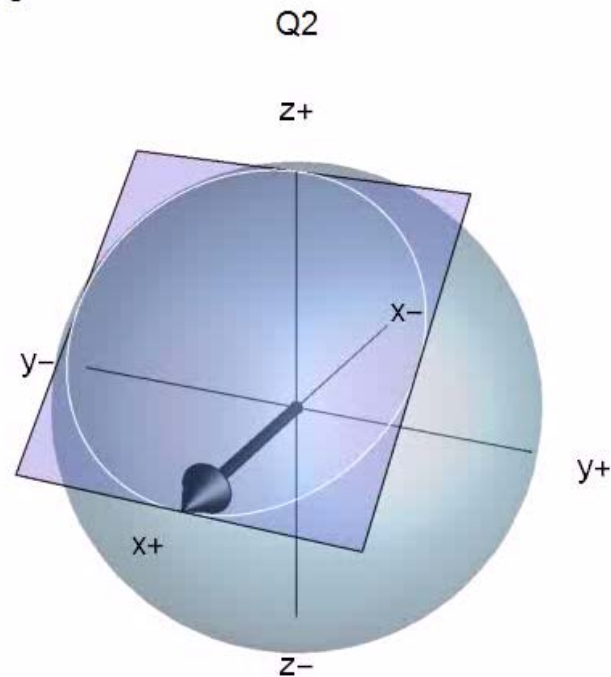
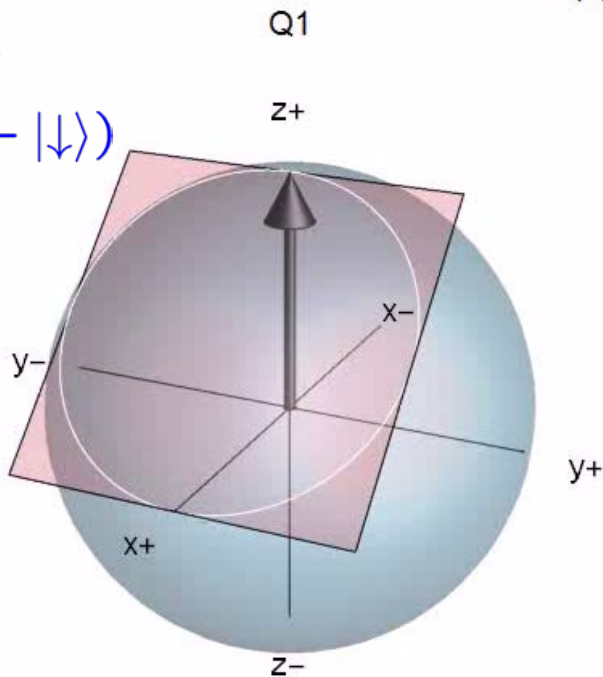
$$H_{xyz} = J (\sigma_1^x \sigma_2^x + \sigma_1^y \sigma_2^y + \sigma_1^z \sigma_2^z)$$

$2|J|\tau = 0 \text{ deg}$

initial state

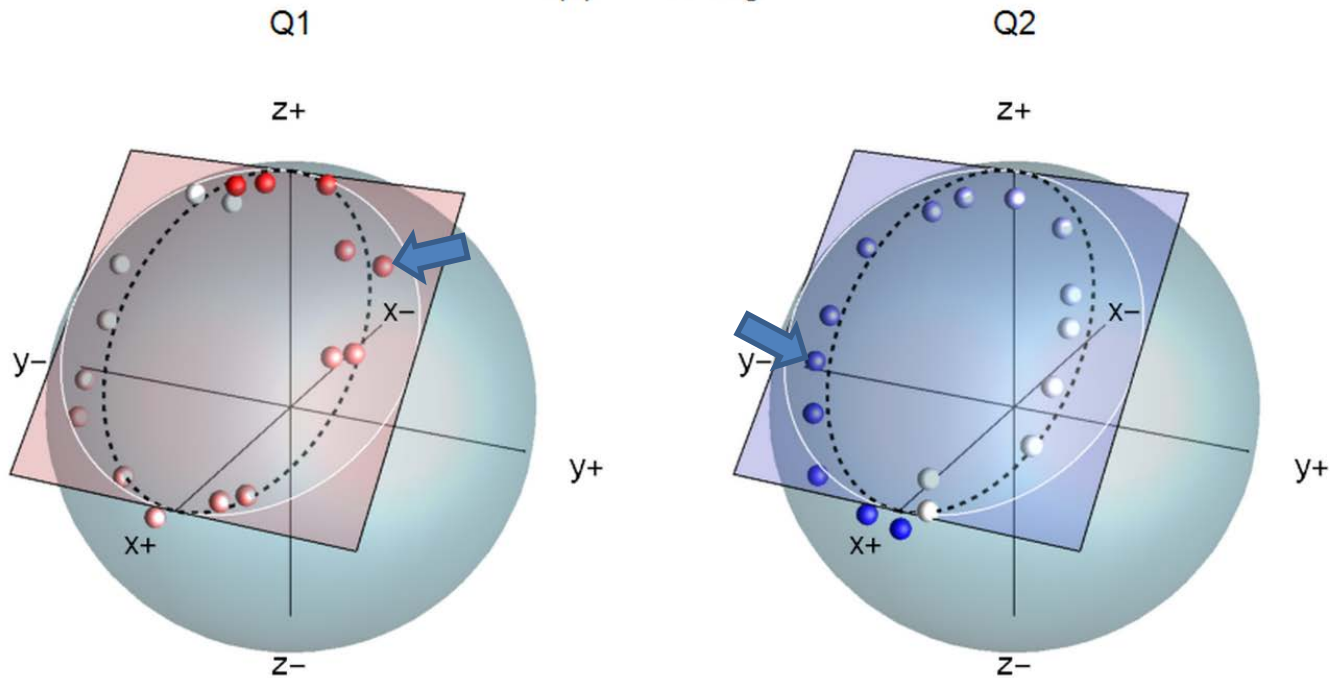
$$\frac{1}{\sqrt{2}} (|\uparrow\rangle (|\uparrow\rangle + |\downarrow\rangle))$$

red: Q1
blue: Q2



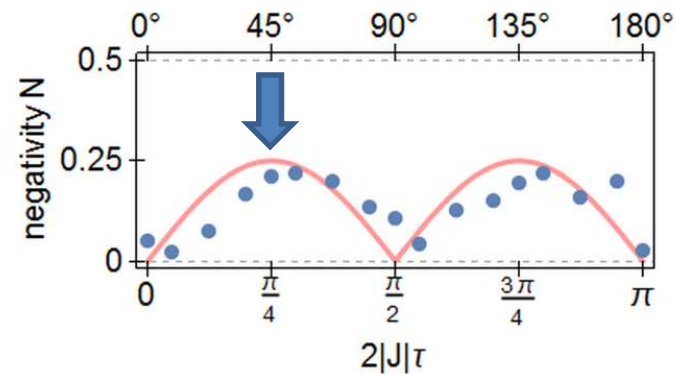
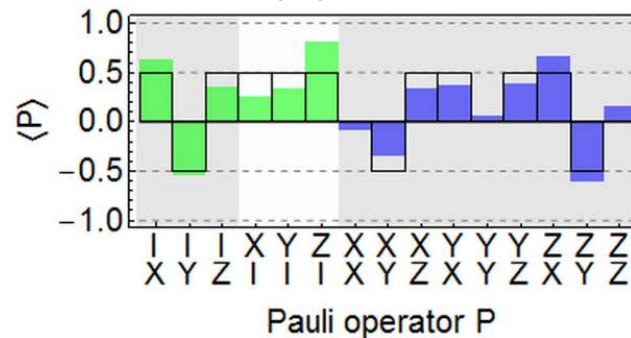
Heisenberg XYZ Interaction: Experimental Data

$2|J|\tau = 180 \text{ deg}$



state fidelity: 95.3 %

$2|J|\tau = 45^\circ$

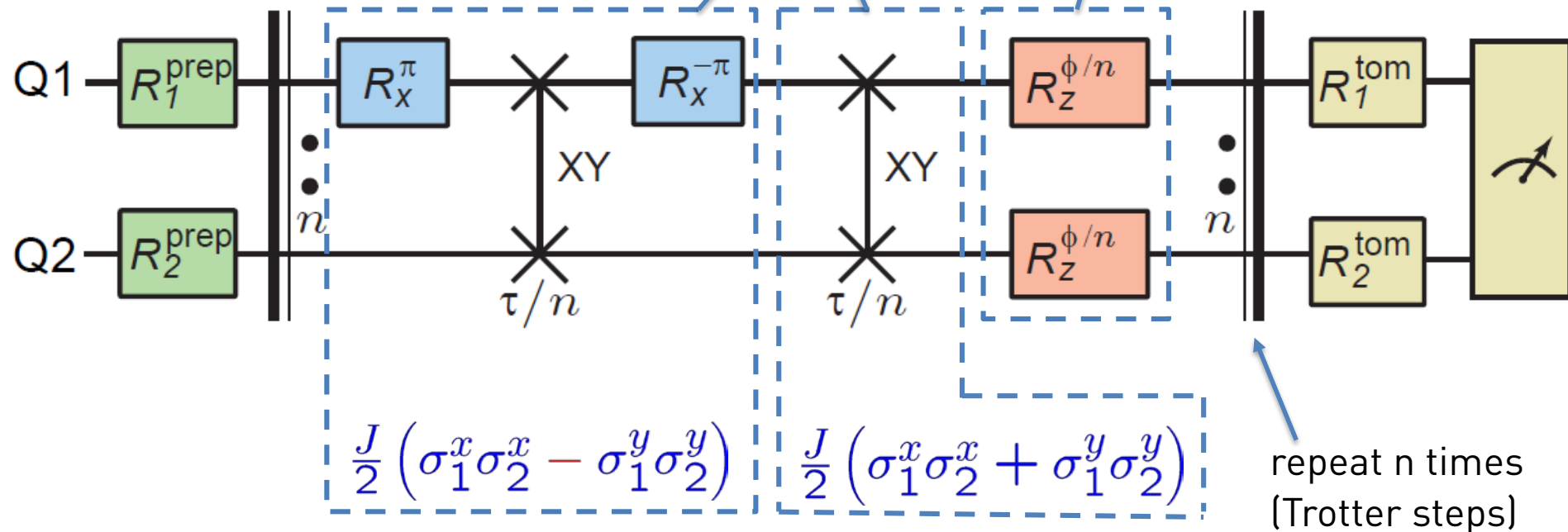


Ising Model with External Field

Hamiltonian to be simulated:

$$H_I = J \sum_{(i,j)} \sigma_i^x \sigma_j^x + \frac{B}{2} \sum_i \sigma_i^z$$

Gate sequence:

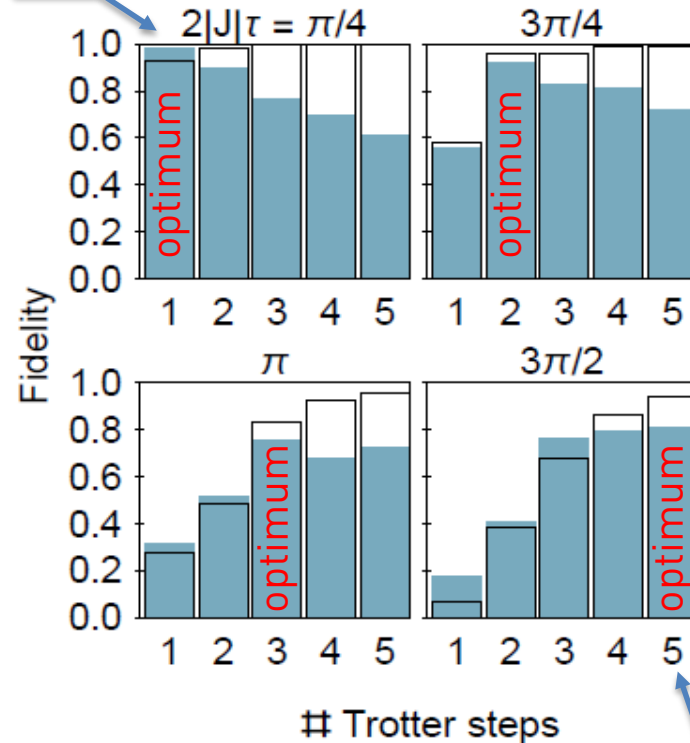


Fidelity of Simulation vs. Trotter Step and Int. Angle

Fidelity of simulated state:
$$F = \left(\text{Tr} \sqrt{\sqrt{\hat{\rho}_{\text{theo}}} \hat{\rho}_{\text{meas}} \sqrt{\hat{\rho}_{\text{theo}}}} \right)^2$$

fewer Trotter steps needed for **small** phase angles (higher-order terms are less important)

Optimal final state fidelity reached at **finite** number of Trotter steps due to limited fidelity individual gates



wire frame:
ideal Trotter fidelity

colored bars:
experimental fidelity

more Trotter steps needed for **large** interaction phase angles (higher order terms are important)

Summary and Outlook

Conclusions:

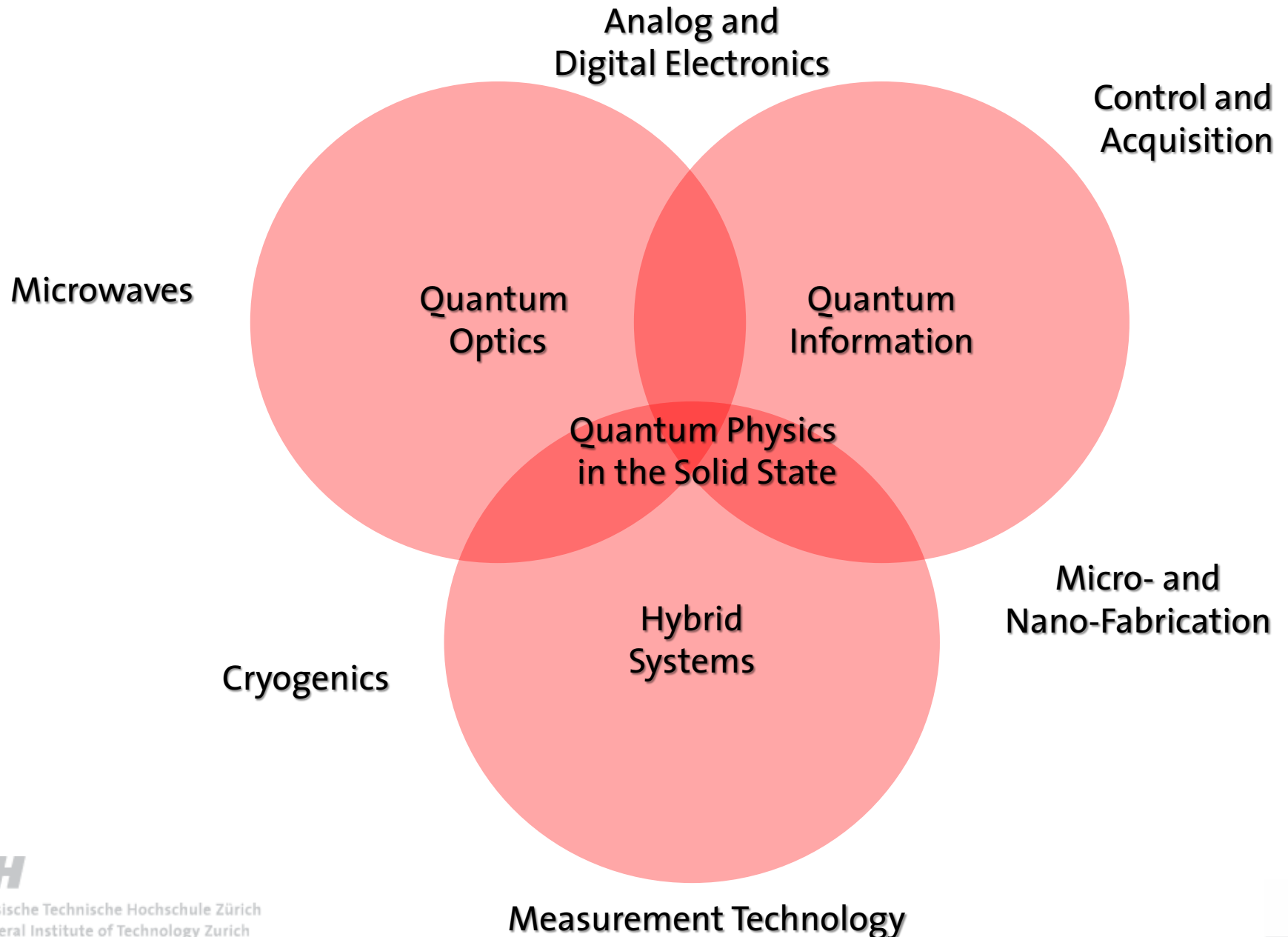
- one of the first digital quantum simulation with superconducting qubits
- simulated time-evolution of paradigmatic spin models
- typical state fidelities above 80% for XYZ and Ising
- used up to 10 two-qubit gates with a continuous interaction parameter combined with 25 single-qubit gates
- challenge: gate calibration with continuous control parameter

Interesting perspectives:

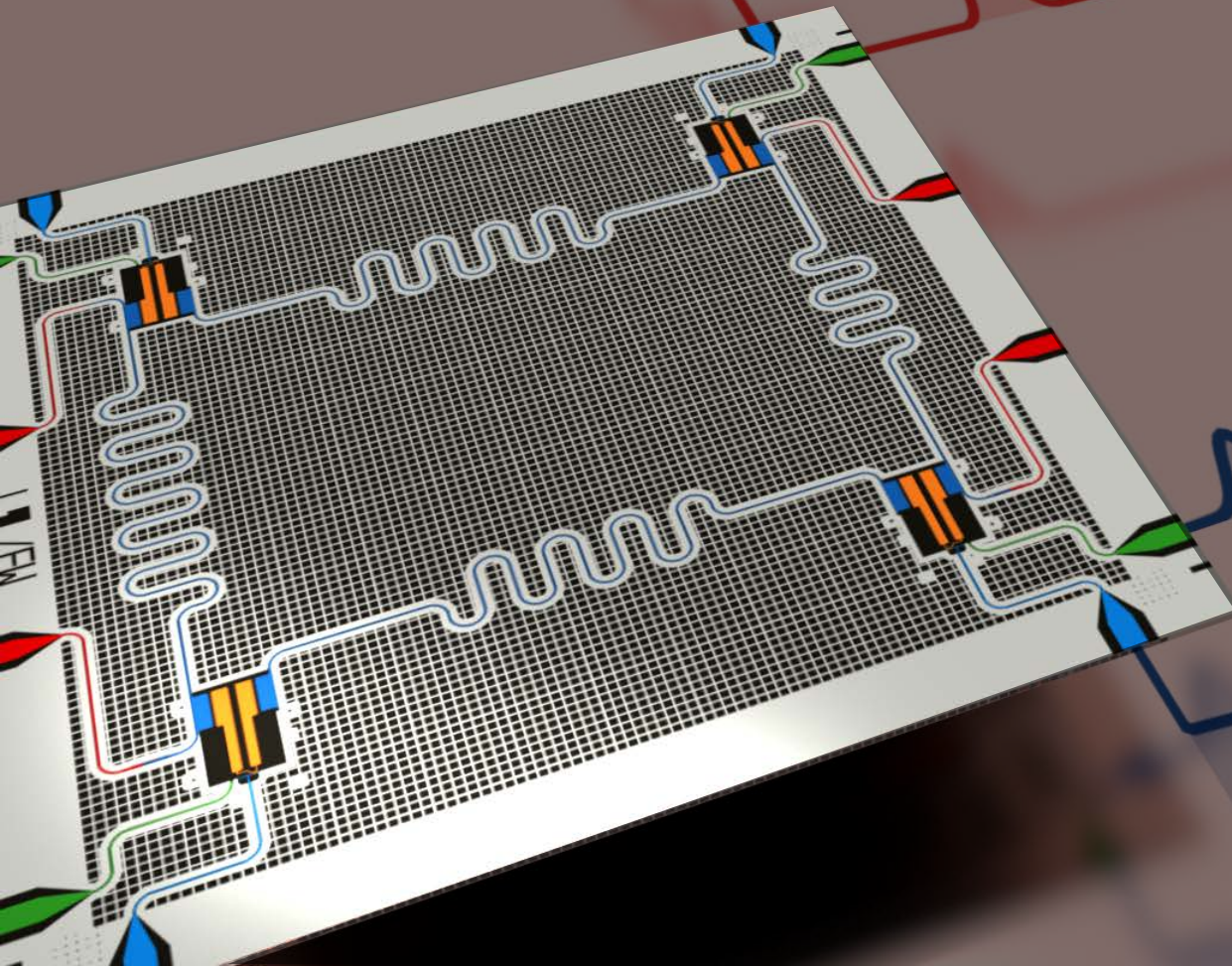
- simulation of Hamiltonians with local interactions using Trotter decomposition
- combine analog and digital methods
- explore time-dependent Hamiltonians
- use bosonic building blocks explicitly
- expand to larger number of spins



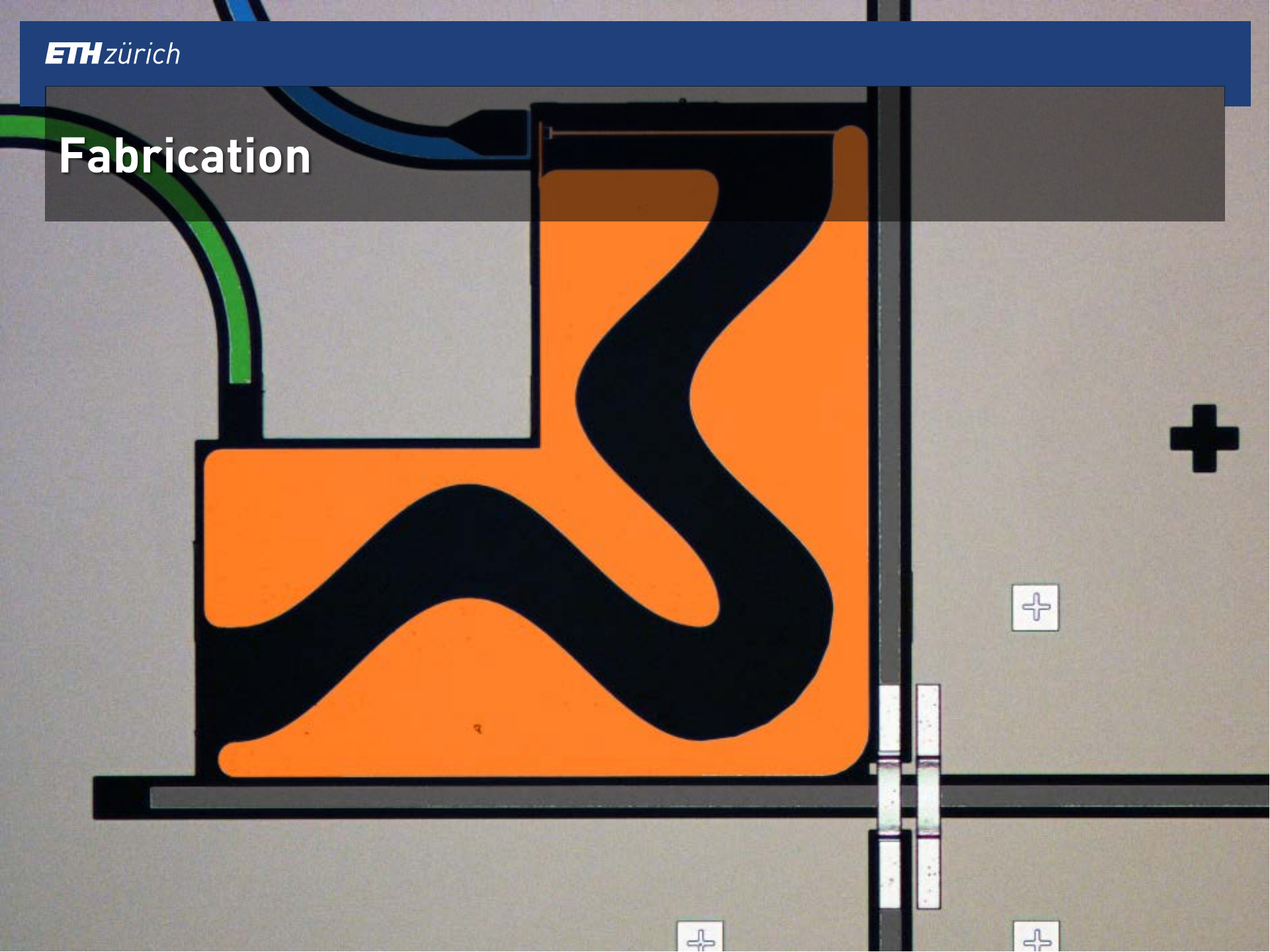
Superconducting Circuits for Quantum Technologies



Design



Fabrication



Control



Automation

Cryogenics



Quantum Engineering



The ETH Zurich Quantum Device Lab

incl. undergrad and summer students



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